

# Functional inequalities for perturbed measures with applications to log-concave measures and to some Bayesian problems.

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We study functional inequalities (Poincaré, Cheeger, log-Sobolev) for probability measures obtained as perturbations. Several explicit results for general measures as well as log-concave distributions are given. The initial goal of this work was to obtain explicit bounds on the constants in view of statistical applications. These results are then applied to the Langevin Monte-Carlo method used in statistics in order to compute Bayesian estimators.

*Keywords:* logconcave measure; Poincaré inequality; Cheeger inequality; logarithmic Sobolev inequality; perturbation; bayesian statistic; sparse learning

## 1. Introduction.

Let  $\mu(dx) = Z_V^{-1} e^{-V(x)} dx$  be a probability measure defined on  $\mathbb{R}^n$ . A priori, we do not require regularity for  $V$  and allow it to take values in  $\mathbb{R} \cup \{-\infty, +\infty\}$ . We denote by  $\mu(f)$  the integral of  $f$  w.r.t.  $\mu$ .

We define the Poincaré constant  $C_P(\mu)$  as the smallest constant  $C$  satisfying

$$\text{Var}_\mu(f) := \mu(f^2) - \mu^2(f) \leq C \mu(|\nabla f|^2), \quad (1.1)$$

for all smooth  $f$ , for instance all bounded  $f$  with bounded derivatives. Similarly the logarithmic-Sobolev constant  $C_{LS}(\mu)$  is defined as the smallest constant such that, for all smooth  $f$  as before,

$$\text{Ent}_\mu(f^2) := \mu(f^2 \ln(f^2)) - \mu(f^2) \ln(\mu(f^2)) \leq C \mu(|\nabla f|^2). \quad (1.2)$$

As it is well known, the Poincaré and the log-Sobolev constants are linked to the exponential stabilization of some markovian dynamics, like the Langevin diffusion i.e. the diffusion semi-group with generator

$$\mathbf{A} = \Delta - \nabla V \cdot \nabla.$$

We shall give more explanations later. For simplicity we will say that  $\mu$  satisfies a Poincaré or a log-Sobolev inequality provided  $C_P(\mu)$  or  $C_{LS}(\mu)$  are finite.

One can also introduce the  $\mathbb{L}^1$  Poincaré constant  $C_C(\mu)$  of  $\mu$ , as the smallest constant such that, for all smooth  $f$ ,

$$\mu(|f - \mu(f)|) \leq C_C(\mu) \mu(|\nabla f|).$$

Replacing  $\mu(f)$  by  $m_\mu(f)$  any  $\mu$  median of  $f$  defines another constant, the Cheeger constant  $C'_C(\mu)$  which satisfies  $\frac{1}{2} C_C \leq C'_C \leq C_C$ .

It is well known [38] that  $\mathbb{L}^1$  and  $\mathbb{L}^2$  Poincaré constants are related by the following

$$C_P(\mu) \leq 4(C'_C)^2(\mu) \leq 4C_C^2(\mu). \quad (1.3)$$

The Cheeger constant is connected to the isoperimetric profile of  $\mu$ , see Ledoux [37].

The initial goal of this work is to study the transference of these inequalities to perturbed measures. Namely, let

$$\mu_F = \frac{e^{-F}}{\mu(e^{-F})} \mu \quad (1.4)$$

be a new probability measure. The question is: what can be said for the Poincaré or the log-Sobolev constant of  $\mu_F$  in terms of the one of  $\mu$  and the properties of  $F$ ? The question includes explicit controls, not only the finiteness of the related constants.

This question is of course not new and some results have been obtained for a long time. We shall only recall results with explicit controls on the constants. The most famous is certainly the following general result of Holley and Stroock (see for example [5])

**Theorem 1.1.** *If  $F$  is bounded, then*

$$C_P(\mu_F) \leq e^{Osc F} C_P(\mu), \quad \text{and} \quad C_{LS}(\mu_F) \leq e^{Osc F} C_{LS}(\mu).$$

Other results, where the constants are not easy to trace, have been obtained in [1] (also see [14] section 7). The result reads as follows: if  $\mu$  satisfies a log-Sobolev inequality, and  $e^{|\nabla F|^2}$  belongs to all the  $\mathbb{L}^p(\mu)$  for  $p < +\infty$ , then  $\mu_F$  also satisfies a log-Sobolev inequality. In particular the result holds true if  $F$  is Lipschitz. We shall revisit this result in subsection 2.2. For the Poincaré inequality, some general results have been obtained in [32] (see e.g. [14] proposition 4.4 for a simplified formulation). Most of the other known results assume some convexity property.

**Definition 1.2.** We shall say that  $\mu$  is *log-concave* if  $V$  is a convex function defined on some convex subset  $U$ . Since  $V$  can be infinite, this definition contains in particular the uniform measure on a convex body.

If  $V$  is strongly convex, i.e. for all  $x \in \mathbb{R}^n$ ,  $\langle u, \nabla^2 V(x) u \rangle \geq \rho |u|^2$ , where  $\langle \cdot, \cdot \rangle$  denotes the Euclidean scalar product and  $\nabla^2 V(x)$  the Hessian of  $V$  computed at point  $x$ , a consequence of Brascamp-Lieb inequality ([13]) is the inequality

$$C_P(\mu) \leq 1/\rho. \quad (1.5)$$

This relation was extended to more general situations and is often called the Bakry-Emery criterion or the curvature-dimension criterion  $CD(\rho, \infty)$  (see [5] for a complete description of curvature-dimension criteria). Some improvements for variable curvature bounds are contained in [16]. Recall

that the Bakry-Emery approach allows to show that in the strongly convex situation

$$C_{LS}(\mu) \leq 2/\rho. \quad (1.6)$$

Combining these results with the Holley-Stroock perturbation result, one can relax the strong convexity assumption in a bounded subset (i.e. assume strong convexity at infinity only).

That  $C_P(\mu) < +\infty$  for general log-concave measures was first shown in 1999 by S. Bobkov in [11]. Another proof was given in [3] using Lyapunov functions, as introduced in [4]. Once one knows that  $C_P(\mu)$  is finite, a particularly important problem is to get some explicit estimates. A celebrated conjecture due to Kannan, Lovász and Simonovits (KLS for short) is that there exists a universal constant  $C$  such that

$$\sigma^2(\mu) \leq C_P(\mu) \leq C \sigma^2(\mu) \quad (1.7)$$

where  $\sigma^2(\mu)$  denotes the largest eigenvalue of the Covariance matrix  $Cov_{i,j}(\mu) = Cov_\mu(x_i, x_j)$  and  $\mu$  is log-concave. The left hand side is immediate. Since, a lot of work has been devoted to this conjecture, satisfied in some special cases. We refer to the book [2] for references before 2015, and to [19] for more information on the Poincaré constant of log-concave measures. Up to very recently the best general known result was

$$C_P(\mu) \leq C n^{\frac{1}{2}} \sigma^2(\mu) \quad (1.8)$$

as a consequence of the results by Lee and Vempala ([39]). It has been very recently announced (see [23] Theorem 1) a drastically better bound namely the existence of an universal constant  $C$  such that

$$C_P(\mu) \leq e^{C\sqrt{\ln(n) \ln(1+\ln(n))}} \sigma^2(\mu). \quad (1.9)$$

The dimension dependence of such results is what is important. Recall that for  $n = 1$  one knows that  $C \leq 12$  according to Bobkov's result (see [11] corollary 4.3).

In this framework, complementary perturbation results have been shown

**Theorem 1.3.** (1) (Miclo, see lemma 2.1 in Bardet et al [7]) *If  $\nabla^2 V \geq \rho Id$  for some  $\rho > 0$  and  $F$  is  $L$ -Lipschitz then*

$$C_P(\mu_F) \leq \frac{2}{\rho} e^{4\sqrt{2n/\pi}} \frac{L^2}{\rho^2}.$$

(2) (see [17] example (3) section 7.1) *With the same assumptions as in (1),*

$$C_P(\mu_F) \leq \frac{1}{2} \left( \frac{2L}{\rho} + \sqrt{\frac{8}{\rho}} \right)^2 e^{L^2/2\rho}.$$

(3) (Barthe-Milman [9] Theorem 2.8) *If  $\mu_F$  is log-concave and  $\mu_F(e^{-F} > K \mu(e^{-F})) \leq \frac{1}{8}$  then*

$$C_P(\mu_F) \leq C^2(1 + \ln K)^2 C_P(\mu).$$

*Here  $C$  is a universal constant.*

The final result (3) is the most general one obtained by transference in the log-concave situation. [9] contains a lot of other results in this direction, [19] contains alternative, simpler but worse results.

Notice that (3) is wrongly recalled in [19] where a square is missing. The remarkable property of (3) or (2) is that the bound is dimension free. On generic examples, like  $V(x) = |x|^2$  and  $F(x) = -|x|$ , the dimension is however present in the Lipschitz constant or in the choice of  $K$ . One can find various other perturbation results in [6] relying on growth conditions, and usually stronger inequalities to get weaker ones.

In [19] we have studied several properties of the Poincaré constant for log-concave measures in particular the transference of these inequalities using absolute continuity, distance, mollification. This study was based on the fact that weak forms of the Poincaré inequality imply the usual form. A similar result was first stated by E. Milman ([41]) and the section 9.2 in [19] is devoted to extensions of E. Milman's results. Actually, in subsection 9.3.1 of [19] devoted to the transference via absolute continuity we missed the point. As explained in [9] the right way to obtain good results is to use the concentration properties of the initial measure (see the proof of Theorem 2.7 in [9]). The main (only) drawback of [41, 9] is that these papers obtain results up to universal constants that are difficult to trace. Thanks to the weak Poincaré inequalities used in [19], it is possible to obtain explicit bounds for these universal constants, sometimes up to a small loss. Why are we so interested in numerical bounds ?

The motivation of this note came from a statistical question. Indeed, log-concave distributions recently deserve attention in Statistics, see e.g. the survey [45]. Our starting point was a question asked to us by S. Gadat on the work [27] by Dalalyan and Tsybakov on sparse regression learning, we shall recall in more details in section 4. The question should be formulated as follows. Let  $\mu_F = \frac{e^{-F}}{\mu(e^{-F})} \mu$  as in (1.4) be a new probability measure. Assume that  $\mu$  is log-concave and that  $F$  is convex (log-concave perturbation of a log-concave measure). Is it possible to control  $C_P(\mu_F)$  by  $C_P(\mu)$  at least up to an universal multiplicative constant ? Since in a sense  $\mu_F$  is "more" log-concave than  $\mu$ , such a statement seems plausible, at least when the "arg-infimum" of  $F$  coincides with the one of  $\mu$ . A first partial answer was obtained by F. Barthe and B. Klartag in [8] Theorem 1:

**Theorem 1.4.** *For  $n \geq 2$ , choose  $V(x) = \sum_{i=1}^n |x_i|^p$  for some  $1 \leq p \leq 2$ , and assume that  $F$  is an even convex function then*

$$C_P(\mu_F) \leq C (\ln(n))^{\frac{2-p}{p}} C_P(\mu),$$

where  $C$  is some universal constant.

Of course here  $C_P(\mu)$  does not depend on the dimension since  $\mu$  is a product measure and

$$C_P(\mu_1 \otimes \dots \otimes \mu_k) \leq \max_{j=1, \dots, k} C_P(\mu_j).$$

Unfortunately this result does not apply to sparse regression as in [27], where  $F$  is not even. We will thus first study perturbation for log-concave measures in section 3 and then see how it can be applied to the aforementioned statistical question. For the latter explicit numerical bounds are of key interest. We will now describe the contents of the present paper.

In the next section we describe general perturbation results both for the Poincaré and the log-Sobolev constants. The naive method we are using does not seem to have been explored with the exception of some results for the log-Sobolev constant contained in [1, 32, 15]. The results can be summarized as follows: for a not too big Lipschitz perturbation  $F$ ,  $\mu_F$  one may explicitly compare  $C_P(\mu_F)$  and

$C_P(\mu)$ , the same for the log-Sobolev constants. If  $F$  is  $C^2$  one can replace the Lipschitz norm by a bound for  $\mathbf{A}F - \frac{1}{2}|\nabla F|^2$ .

In subsection 2.3 we give a surprisingly simple application to mollified measures, extending part of the results in [7].

In section 3 we show how to improve these results when  $\mu_F$  is log-concave. We first study how to get explicit controls in the results obtained by Barthe and Milman [9] using concentration, by plugging our explicit results of [19]. We then transpose the results of the first section. Essentially, if in general one needs to control the uniform norm of  $\nabla F$ , in the log-concave case it is enough to control its  $\mathbb{L}^2$  norm. As an immediate consequence, we obtain in subsection 3.3, explicit controls for the pre-constants in results by Bobkov [11] and Klartag ([35]). The idea of proof is to consider  $\mu(dx) = e^{-V(x)} dx$  as a perturbation  $\nu_F$  of  $\nu(dx) = e^{-V(x) - \lambda^p \sum |x_i|^p} dx$ . For  $p = 2$ ,  $1 < p < 2$  and  $p > 2$  respectively we may use Bakry-Emery, Theorem 1.4 and another result obtained in [8] for unconditional measures, in order to control  $C_P(\nu)$ . We then use our results to come back to  $C_P(\mu)$ . The advantage of our method is that it furnishes an explicit bound for the constant  $C$ .

The final section is devoted to the application to two statistical models: linear regression, following [27] and identification as in [31]. In both cases the authors propose to use a Langevin Monte-Carlo algorithm in order to compute some bayesian estimator. To get some explicit rate of convergence for this algorithm is thus of key importance. This method already used in [47] becomes very popular these last years. One can look at [24, 25, 30] among the huge recent literature. We shall focus on the linear regression case and give explicit controls for the rate of convergence of the Langevin Monte-Carlo algorithm proposed in [27], that are much better than the ones suggested therein. The same comparison can be made with several papers using the Foster-Lyapunov function approach popularized by Meyn and Tweedie. Notice that the kinetic Langevin algorithm recently proposed in [26] does not enter the framework of the present paper, since in this case the required Poincaré inequality involving the square gradient operator of Bakry-Emery (and no more the square of the Euclidean gradient) is not satisfied. Ideas from [20] involving weighted Poincaré inequalities should be an interesting substitute.

## 2. Smooth perturbations in general.

### 2.1. Poincaré inequalities.

We shall follow a direct perturbation approach, as the one of [15] section 4 used for the log-Sobolev constant. Since we shall use a similar but slightly different approach in the next sections we first isolate the starting point of the proof.

For a smooth  $f$  and a constant  $a$  we may write

$$\text{Var}_{\mu_F}(f) \leq \mu_F((f - a)^2) = \frac{1}{\mu(e^{-F})} \mu\left(\left((f - a)e^{-\frac{1}{2}F}\right)^2\right).$$

We choose

$$a = \frac{\mu\left(f e^{-\frac{1}{2}F}\right)}{\mu\left(e^{-\frac{1}{2}F}\right)} \tag{2.1}$$

so that the function  $(f - a)e^{-\frac{1}{2}F}$  is  $\mu$  centered. One can thus use the Poincaré inequality for  $\mu$  in order to get, for all  $a$  and all  $\varepsilon > 0$ ,

$$\mu_F((f - a)^2) \leq C_P(\mu) \int |\nabla f - \frac{1}{2}(f - a)\nabla F|^2 d\mu_F. \quad (2.2)$$

We can state a first elementary result

**Theorem 2.1.** *If  $F$  is  $L$ -Lipschitz on the support of  $\mu$  and if there exists  $\varepsilon > 0$  such that*

$$s := \frac{1}{4}(1 + \varepsilon)C_P(\mu)L^2 < 1,$$

then

$$C_P(\mu_F) \leq \frac{(1 + \varepsilon^{-1})C_P(\mu)}{1 - s}.$$

**Proof.** Starting with (2.2), one deduces, for all  $\varepsilon > 0$ ,

$$\mu_F((f - a)^2) \leq C_P(\mu) \left( (1 + \varepsilon^{-1})\mu_F(|\nabla f|^2) + \frac{1 + \varepsilon}{4}\mu_F((f - a)^2|\nabla F|^2) \right). \quad (2.3)$$

If  $F$  is  $L$ -Lipschitz we deduce from (2.3),

$$\mu_F((f - a)^2) \leq (1 + \varepsilon^{-1})C_P(\mu)\mu_F(|\nabla f|^2) + \frac{1 + \varepsilon}{4}C_P(\mu)L^2\mu_F((f - a)^2),$$

so that

$$\text{Var}_{\mu_F}(f) \leq \mu_F((f - a)^2) \leq \frac{(1 + \varepsilon^{-1})C_P(\mu)}{1 - \frac{1}{4}(1 + \varepsilon)C_P(\mu)L^2}\mu_F(|\nabla f|^2)$$

as soon as  $\exists \varepsilon > 0$  such that

$$\frac{1}{4}(1 + \varepsilon)C_P(\mu)L^2 < 1.$$

□

This theorem is quite sharp. Indeed recall that (see e.g. [5, Prop. 4.4.2]) for every 1-Lipschitz function  $F$  and  $s < \sqrt{4/C_P(\mu)}$ , one has  $\mu(e^{sF}) < \infty$ . The bound on  $s$  in the theorem is thus close to the optimal one ensuring that  $\mu_F$  is a Probability measure.

Similarly for the Cheeger constant we may state

**Proposition 2.2.** *If  $F$  is  $L$ -Lipschitz on the support of  $\mu$  and  $C_C(\mu)L < 1$ ,*

$$C'_C(\mu_F) \leq \frac{C_C(\mu)}{1 - C_C(\mu)L}.$$

**Proof.** First, for all  $a$ ,

$$\mu_F(|f - m_{\mu_F}(f)|) \leq \mu_F(|f - a|) = \frac{1}{\mu(e^{-F})}\mu\left(|(f - a)e^{-F}|\right)$$

so that choosing this time

$$a = \frac{\mu(f e^{-F})}{\mu(e^{-F})}$$

we get the result for the Cheeger constant by using

$$\begin{aligned} \mu_F(|f - a|) &\leq C_C(\mu) (\mu_F(|\nabla f|) + \mu_F(|f - a| |\nabla F|)) \\ &\leq C_C(\mu) \mu_F(|\nabla f|) + C_C(\mu) L \mu_F(|f - a|). \end{aligned}$$

□

The proof of these two statements is so simple that we cannot believe that the result is not known. In fact, the only comparable result we found, using seemingly more intricate techniques, is in [1, Th. 2.7] but with roughly a factor 16 in our favor and an explicit Poincaré constant.

What is remarkable is that we may even allow more general condition than  $F$  to be  $L$ -Lipschitz however assuming more regularity for  $F$ . It requires to be careful when  $\mu$  is compactly supported.

**Theorem 2.3.** a) If  $V$  is of  $C^1$  class and  $F$  is of  $C^2$  class and satisfies for some  $\varepsilon > 0$

$$C_P(\mu) \left\| \left( \mathbf{A}F - \frac{1}{2} |\nabla F|^2 \right)^+ \right\|_{\mathbb{L}^\infty} \leq 2(1 - \varepsilon),$$

where  $u^+ = \max(u, 0)$ , then

$$C_P(\mu_F) \leq \frac{C_P(\mu)}{\varepsilon}.$$

b) Assume that  $U := \{V < +\infty\}$  is an open subset with a smooth boundary  $\partial U$  and that  $V$  is of  $C^1$  class in  $U$ . Let  $F$  be of  $C^2$  class and such that  $\partial_n F \geq 0$  on  $\partial U$  where  $\partial_n$  denotes the normal derivative pointing outward. If  $F$  satisfies for some  $\varepsilon > 0$

$$C_P(\mu) \left\| \left( \mathbf{A}F - \frac{1}{2} |\nabla F|^2 \right)^+ \right\|_{\mathbb{L}^\infty(U)} \leq 2(1 - \varepsilon),$$

where  $u^+ = \max(u, 0)$ , then

$$C_P(\mu_F) \leq \frac{C_P(\mu)}{\varepsilon}.$$

**Proof.** a) If one allows  $F$  to be more regular, one can replace (2.3) by another inequality. Indeed starting with (2.2) and using the classical integration by parts formula

$$\mu_F(\langle \nabla g, \nabla h \rangle) = -\mu_F(g \mathbf{A}_F h)$$

where  $\mathbf{A}_F = \mathbf{A} - \nabla F \cdot \nabla = \Delta - \nabla V \cdot \nabla - \nabla F \cdot \nabla$ , we have

$$\mu_F((f - a)^2) \leq C_P(\mu) \int |\nabla f - \frac{1}{2}(f - a)\nabla F|^2 d\mu_F$$

$$\begin{aligned}
&\leq C_P(\mu) \left( \mu_F(|\nabla f|^2) - \frac{1}{2} \mu_F(\langle \nabla(f-a)^2, \nabla F \rangle) + \frac{1}{4} \mu_F((f-a)^2 |\nabla F|^2) \right) \\
&\leq C_P(\mu) \mu_F(|\nabla f|^2) + \frac{1}{2} C_P(\mu) \mu_F \left( (f-a)^2 \left[ \mathbf{A}_F F + \frac{1}{2} |\nabla F|^2 \right] \right). \quad (2.4)
\end{aligned}$$

Finally

$$\mu_F((f-a)^2) \leq C_P(\mu) \mu_F(|\nabla f|^2) + \frac{1}{2} C_P(\mu) \mu_F \left( (f-a)^2 \left[ \mathbf{A}_F F - \frac{1}{2} |\nabla F|^2 \right] \right). \quad (2.5)$$

It remains to bound

$$\mu_F \left( (f-a)^2 \left[ \mathbf{A}_F F - \frac{1}{2} |\nabla F|^2 \right] \right) \leq \left\| \left( \mathbf{A}_F F - \frac{1}{2} |\nabla F|^2 \right)^+ \right\|_{\mathbb{L}^\infty} \mu_F((f-a)^2)$$

to conclude.

b) We have to start again with (2.2) where integration holds in  $U$ . This yields

$$\begin{aligned}
\mu_F(f-a)^2 &\leq C_P(\mu) \int_U |\nabla f - \frac{1}{2}(f-a)\nabla F|^2 d\mu_F \quad (2.6) \\
&\leq C_P(\mu) \left( \mu_F(|\nabla f|^2) - \frac{1}{2} \mu_F(\langle \nabla(f-a)^2, \nabla F \rangle) + \frac{1}{4} \mu_F((f-a)^2 |\nabla F|^2) \right).
\end{aligned}$$

To control the second term we have to use Green's formula to integrate by parts

$$\mu_F(\langle \nabla(f-a)^2, \nabla F \rangle) = -\mu_F((f-a)^2 \mathbf{A}_F F) + \mu_F^\partial((f-a)^2 \partial_n F) \quad (2.7)$$

where  $\mu_F^\partial$  denotes the surface measure on  $\partial U$  and  $\partial_n$  denotes the normal derivative pointing outward. The end of the proof is then similar.  $\square$

**Example 1.** Let us describe a very simple case which illustrates the difference between Th.2.1 and Th.2.3. For  $n = 1$ , let  $d\mu = \frac{1}{2}e^{-|x|}dx$  for which  $C_P(\mu) = 4$ , and consider  $F(x) = \rho|x|$  which is  $|\rho|$ -Lipschitz. An application of Th.2.1 shows that  $\mu_F$  still satisfies a Poincaré inequality if  $|\rho| < 1$  whereas Th.2.3 implies that  $\mu_F$  satisfies a Poincaré inequality as soon as  $\rho > -1$  which is optimal in this case.  $\diamond$

**Example 2.** For  $\rho \in \mathbb{R}^+$  consider

$$\mu^\rho(dx) = Z_\rho^{-1} e^{-V(x) - \frac{\rho}{2}|x|^2} dx,$$

i.e.

$$F(x) = \frac{\rho}{2}|x|^2$$

which is not Lipschitz continuous. We thus have

$$(\mathbf{A}_F F - \frac{1}{2} |\nabla F|^2)(x) = \rho \left( n - \langle x, \nabla V(x) \rangle - \frac{1}{2} \rho |x|^2 \right).$$



Hence if  $\langle x, \nabla V \rangle \geq -K - K'|x|^2$ ,  $\mathbf{A}F - \frac{1}{2}|\nabla F|^2 \leq \rho(n+K) + (K' - \rho/2)|x|^2$  so that  $C_P(\mu^\rho) \leq \frac{C_P(\mu)}{1-\varepsilon}$  as soon as  $\rho \leq \frac{2(1-\varepsilon)}{C_P(\mu)(n+K)}$  and  $\rho > K'$ .

Thus a (very) small Gaussian perturbation of a measure satisfying some Poincaré inequality is still satisfying some Poincaré inequality. Though natural, we do not know any other way to prove such a result.  $\diamond$

**Example 3.** Assume now that  $\mu$  is supported by some convex body  $\Omega$ . Denoting by  $\nu$  the uniform measure on  $\Omega$  we have  $\mu = \nu_V$ . Hence if  $V$  is  $L$ -Lipschitz with  $L^2 C_P(\nu) \leq 1$  we have  $C_P(\mu) \leq 4C_P(\nu)$ . Can we compare this result with the Holley-Stroock perturbation result? Of course we have  $\text{Osc}(V) \leq LD$  where  $D$  denotes the diameter of  $\Omega$ , so that Holley-Stroock furnishes  $C_P(\mu) \leq e^{LD} C_P(\nu)$ . We may always use the bound  $C_P(\nu) \leq D^2/\pi^2$  going back to Poincaré himself (actually to Payne and Weinberger [44] in the optimal form here, optimal since the bound is exact in dimension 1) yielding  $L \leq \pi/D$  hence  $LD \leq \pi$ , hence  $C_P(\mu) \leq e^\pi D^2/\pi^2$  which is slightly worse than the one we described above since  $e^\pi \geq 4$  but clearly more or less equivalent. It is however well known that in many situations the diameter bound is far to be the good one. For instance if  $\Omega = B(0, R)$  the euclidean ball of radius  $R$ ,  $C_P(\nu) \leq R^2/n$  as soon as  $n \geq 4$ . Hence we may use theorem 2.1 as soon as  $L \leq \sqrt{n}/R$  and obtain a bound for  $C_P(\mu)$  of order  $R^2/n$ , while Holley-Stroock gives a pre-factor of order  $e^{c\sqrt{n}}$ .  $\diamond$

The previous results can be compared with the one in [32] where  $\mu$  is assumed to satisfy a log-Sobolev inequality. Actually, one can almost recover Gong-Wu result. Indeed, in (2.3) and (2.5), the final step requires to control a term in the form

$$\mu_F((f-a)^2 G),$$

with  $G = |\nabla F|^2$  in (2.3) or  $G = [\mathbf{A}F - \frac{1}{2}|\nabla F|^2]^+$  in (2.5). According to the variational definition of relative entropy we have, for  $\alpha > 0$ ,

$$\mu_F((f-a)^2 G) \leq \frac{1}{\alpha} \text{Ent}_{\mu_F}((f-a)^2) + \frac{1}{\alpha} \mu_F((f-a)^2) \ln \mu_F(e^{\alpha G}). \quad (2.8)$$

The relative entropy in the right hand side of the previous inequality can be controlled by the energy provided  $\mu$  satisfies a log-Sobolev inequality. We thus have

**Proposition 2.4.** We suppose here that  $\mu$  satisfies a logarithmic Sobolev inequality and thus that  $C_{LS}(\mu)$  is finite.

(i) Assume that for some positive  $s, t$  and  $\theta$ ,

$$\frac{1}{s} + \frac{1+\theta}{4t} C_{LS}(\mu) := D_1 \leq 1 \quad (2.9)$$

Assume in addition that there exist  $\alpha > 0$  and  $\varepsilon > 0$  such that

$$T'_1 := \frac{(1+\varepsilon)C_P(\mu)}{4\alpha} T_1 < 1$$

where

$$T_1 := \ln \mu_F(e^{\alpha|\nabla F|^2}) + \frac{1}{1-D_1} \left[ \frac{1}{s} \ln \mu_F(e^{sF}) + \ln(\mu(e^{-F})) + \frac{(1+\theta)C_{LS}(\mu)}{4t} \ln \mu_F(e^{t|\nabla F|^2}) \right].$$

Then,

$$C_P(\mu_F) \leq \frac{1}{1-T'_1} C_P(\mu) \left( (1+\varepsilon^{-1}) + \frac{(1+\theta^{-1})(1+\varepsilon)}{4\alpha} C_{LS}(\mu) \right).$$

(ii) Assume for some positive  $s$  and  $t$ ,

$$\frac{1}{s} + \frac{1}{t} C_{LS}(\mu) := D_2 \leq 1 \quad (2.10)$$

Assume in addition that there exists  $\alpha > 0$  such that

$$T'_2 := \frac{C_P(\mu)}{2\alpha} T_2 < 1$$

where

$$T_2 := \ln \mu_F(e^{\alpha[\mathbf{A}F - \frac{1}{2}|\nabla F|^2]}) + \frac{1}{1-D_2} \left[ \frac{1}{s} \ln \mu_F(e^{sF}) + \ln(\mu(e^{-F})) + \frac{C_{LS}(\mu)}{t} \ln \mu_F(e^{t[\mathbf{A}F - \frac{1}{2}|\nabla F|^2]}) \right].$$

Then

$$C_P(\mu_F) \leq \frac{1}{1-T'_2} C_P(\mu) \left( 1 + \frac{C_{LS}(\mu)}{\alpha} \right).$$

**Proof.** Replacing  $F$  by  $F + \ln(\mu(e^{-F}))$  we may assume for simplicity that  $\mu(e^{-F}) = 1$ . Defining  $g = e^{-F/2}(f-a)$  we have  $\mu(g^2) = \mu_F((f-a)^2)$ . Hence for all  $\theta > 0$ ,

$$\begin{aligned} \text{Ent}_{\mu_F}((f-a)^2) &= \text{Ent}_{\mu}(g^2) + \mu(g^2 F) \\ &\leq C_{LS}(\mu) \mu(|\nabla g|^2) + \mu(g^2 F) \\ &\leq C_{LS}(\mu) (1+\theta^{-1}) \mu_F(|\nabla f|^2) + C_{LS}(\mu) \frac{1+\theta}{4} \mu_F((f-a)^2 |\nabla F|^2) \\ &\quad + \mu_F((f-a)^2 F), \end{aligned} \quad (2.11)$$

if we follow the proof of Theorem 2.1, or

$$\text{Ent}_{\mu_F}((f-a)^2) \leq C_{LS}(\mu) \mu_F(|\nabla f|^2) + C_{LS}(\mu) \mu_F \left( (f-a)^2 [\mathbf{A}F - \frac{1}{2}|\nabla F|^2] \right) + \mu_F((f-a)^2 F) \quad (2.12)$$

if we follow the proof of Theorem 2.3. Using (2.8) we get

$$\begin{aligned} \text{Ent}_{\mu_F}((f-a)^2) &\leq C_{LS}(\mu) \mu_F(|\nabla f|^2) + \frac{C_{LS}(\mu)}{t} \text{Ent}_{\mu_F}((f-a)^2) \ln \mu_F(e^{t|\nabla F|^2}) + \\ &\quad + \frac{1}{s} \text{Ent}_{\mu_F}((f-a)^2) \ln \mu_F(e^{sF}), \end{aligned}$$

i.e.

$$\begin{aligned} \text{Ent}_{\mu_F}((f-a)^2) &\leq \frac{1}{1-D_1} \frac{C_{LS}(\mu)}{1+\theta} \mu_F(|\nabla f|^2) + \\ &\quad + \frac{\mu_F((f-a)^2)}{1-D_1} \left[ \frac{1}{s} \ln \mu_F(e^{sF}) + \frac{C_{LS}(\mu)}{t} \ln \mu_F(e^{t|\nabla F|^2}) \right]. \end{aligned}$$

Coming back to (2.3) we have

$$\begin{aligned} \mu_F((f-a)^2) &\leq C_P(\mu)(1+\varepsilon^{-1})\mu_F(|\nabla f|^2) + \\ &\quad + \frac{(1+\varepsilon)C_P(\mu)}{4} \left( \frac{1}{\alpha} \text{Ent}_{\mu_F}((f-a)^2) + \frac{1}{\alpha} \mu_F((f-a)^2) \ln \mu_F(e^{\alpha|\nabla F|^2}) \right) \end{aligned}$$

yielding (i). The proof of (ii) is similar and left to the reader.  $\square$

Of course such a result where we can replace uniform bounds by exponential integrability, is difficult to apply, but the method will be useful to get a perturbation result for the log-Sobolev constant. Part of the result has been described in [15]. Note however that as  $\mu$  satisfies a logarithmic Sobolev inequality then one has Gaussian integrability properties, so that at least for every  $s < 1/C_{LS}(\mu)$

$$\int e^{s|x|^2} d\mu < \infty$$

and thus if for some positive  $a, b$  sufficiently small, one has  $|F|, |\nabla F|^2 < a + b|x|^2$ ,  $T_1$  is then finite and can be made explicit.

## 2.2. Log-Sobolev inequality.

We can similarly look at the log-Sobolev constant

**Theorem 2.5.** *Assume that  $F$  is  $L$ -Lipschitz on the support of  $\mu$  and that  $\mu$  satisfies a log-Sobolev inequality with constant  $C_{LS}(\mu)$ . Also assume for simplicity that  $\mu(e^{-F}) = 1$ .*

(1) *If  $\sup_{x \in \text{supp}(\mu)} F(x) = M$ , then for all  $\theta > 0$ ,*

$$C_{LS}(\mu_F) \leq (1 + \theta^{-1}) C_{LS}(\mu) + C_P(\mu_F) \left( \frac{1 + \theta}{4} L^2 C_{LS}(\mu) + M + 2 \right).$$

(2) *For all  $\beta > 0$ , for all  $\theta > 0$ ,*

$$\begin{aligned} C_{LS}(\mu_F) &\leq \frac{(\beta + 1)(1 + \theta^{-1})}{\beta} C_{LS}(\mu) + C_P(\mu_F) (2 + \mu(F)) \\ &\quad + L^2 C_P(\mu_F) C_{LS}(\mu) \left( \frac{(1 + \theta)(1 + \beta)}{4\beta} + \frac{\beta^2}{2} \right). \end{aligned}$$

*If in addition the condition in Theorem 2.1 is satisfied for some  $\varepsilon > 0$  we may replace  $C_P(\mu_F)$  by the bound obtained in Theorem 2.1.*

Remark that the first statement does not enter the framework of Holley-Stroock's theorem as only a one sided bound is assumed on  $F$ .

**Proof.** Let  $f$  be smooth and such that  $\mu_F(f^2) = 1$ . Defining  $g = e^{-F/2} f$  we have  $\mu(g^2) = 1$ . It follows for all  $\theta > 0$ ,

$$\text{Ent}_{\mu_F}(f^2) = \mu_F(f^2 \ln(f^2)) = \mu(g^2 \ln(g^2)) + \mu(g^2 F)$$

$$\begin{aligned}
&\leq C_{LS}(\mu) \mu(|\nabla g|^2) + \mu(g^2 F) & (2.13) \\
&\leq C_{LS}(\mu) (1 + \theta^{-1}) \mu_F(|\nabla f|^2) + C_{LS}(\mu) \frac{1 + \theta}{4} \mu_F(f^2 |\nabla F|^2) + \mu_F(f^2 F).
\end{aligned}$$

In the first case, we can bound the sum of the last two terms by

$$\left( C_{LS}(\mu) \frac{1 + \theta}{4} L^2 + M \right) \mu_F(f^2)$$

and apply the Poincaré inequality for  $\mu_F$  provided  $\mu_F(f) = 0$ . To conclude it is then enough to recall Rothaus lemma (Lemma 5.14 in [5]),

$$\text{Ent}_\nu(f^2) \leq \text{Ent}_\nu((f - \nu(f))^2) + 2 \text{Var}_\nu(f).$$

If  $F$  is not bounded above we can use the variational definition of relative entropy as before:

$$\mu_F(f^2 F) \leq \frac{1}{\alpha} \text{Ent}_{\mu_F}(f^2) + \frac{1}{\alpha} \mu_F(f^2) \ln(\mu_F(e^{\alpha F})).$$

Gathering all the previous bounds we have obtained, provided  $\alpha > 1$ ,

$$\text{Ent}_{\mu_F}(f^2) \leq \frac{\alpha}{\alpha - 1} \left( C_{LS}(\mu) (1 + \theta^{-1}) \mu_F(|\nabla f|^2) + C \mu_F(f^2) \right) \quad (2.14)$$

with

$$C = C_{LS}(\mu) \frac{1 + \theta}{4} L^2 + \frac{1}{\alpha} \ln(\mu_F(e^{\alpha F})).$$

We may then argue as before using Rothaus lemma again. The bound

$$\mu_F(e^{\alpha F}) = \mu(e^{(\alpha-1)F}) \leq e^{(\alpha-1)\mu(F) + (C_{LS}(\mu) L^2 (\alpha-1)^2 / 2)}$$

is known as the Herbst argument (see e.g. [5] Proposition 5.4.1). □

As before we may replace the Lipschitz assumption by an integrability condition yielding the next result whose proof, similar to the previous one, is omitted

**Theorem 2.6.** *Suppose that  $\mu$  satisfies a logarithmic Sobolev inequality with constant  $C_{LS}(\mu)$  and that  $\mu(e^{-F}) = 1$ . Assume that there exist  $\alpha > 1$  and  $\beta, \theta > 0$  such that*

$$\mu_F(e^{\alpha F}) < \infty, \quad \mu_F(e^{\beta |\nabla F|^2}) < \infty$$

and

$$C_{LS}(\mu) \frac{1 + \theta}{4\beta} + \frac{1}{\alpha} := \delta < 1$$

then  $\mu_F$  also satisfies a logarithmic Sobolev inequality with constant  $C_{LS}(\mu_F)$  equal to

$$\frac{1}{1 - \delta} \left[ C_{LS}(\mu) (1 + \theta^{-1}) + C_P(\mu_F) \left( 2 + C_{LS}(\mu) \frac{1 + \theta}{4\beta} \log \mu_F(e^{\beta |\nabla F|^2}) + \frac{1}{\alpha} \log \mu_F(e^{\alpha F}) \right) \right].$$

The previous Theorem 2.6 is a version of the one obtained in [1] as recalled in the introduction.

Finally, if  $F$  is more regular we may replace (2.13) by the following

$$\text{Ent}_{\mu_F}(f^2) \leq C_{LS}(\mu)\mu(|\nabla f|^2) + \frac{1}{2} C_{LS}(\mu) \mu_F \left( f^2 [\mathbf{A}F - \frac{1}{2}|\nabla F|^2] \right) + \mu_F(f^2 F).$$

Arguing as before we thus obtain

**Theorem 2.7.** *Suppose that  $\mu$  satisfies a logarithmic Sobolev inequality with constant  $C_{LS}(\mu)$  and that  $V$  is  $C^1$ . Assume that there exist  $\alpha > 1$  and  $\beta > 0$  such that*

$$\mu_F(e^{\alpha F}) < \infty, \quad \mu_F(e^{\beta[\mathbf{A}F - \frac{1}{2}|\nabla F|^2]}) < \infty$$

and

$$C_{LS}(\mu) \frac{1}{2\beta} + \frac{1}{\alpha} := \delta < 1$$

then  $\mu_F$  also satisfies a logarithmic Sobolev inequality with constant  $C_{LS}(\mu_F)$  equal to

$$\frac{1}{1-\delta} \left[ C_{LS}(\mu) + C_P(\mu_F) \left( 2 + C_{LS}(\mu) \frac{1}{2\beta} \log \mu_F(e^{\beta[\mathbf{A}F - \frac{1}{2}|\nabla F|^2]}) + \frac{1}{\alpha} \log \mu_F(e^{\alpha F}) \right) \right].$$

If in addition the condition in Theorem 2.3 is satisfied for some  $\varepsilon > 0$  we may replace  $C_P(\mu_F)$  by the bound obtained in Theorem 2.3.

**Remark 2.8.** Notice that if  $\mathbf{A}F - \frac{1}{2}|\nabla F|^2$  is non-positive at infinity (which is often the case in concrete examples), and  $F$  is  $C^2$ ,  $e^{\beta[\mathbf{A}F - \frac{1}{2}|\nabla F|^2]}$  is bounded for all  $\beta > 0$ , so that the condition in the previous theorem reduces to the integrability of  $e^{\alpha F}$  for some  $\alpha > 1$ . The most stringent condition is thus the one in theorem 2.3 ensuring the finiteness of the Poincaré constant.  $\diamond$

### 2.3. Application to mollified measures.

Let  $\nu$  be a given probability measure (non necessarily absolutely continuous) and define  $\nu^\sigma$  as the convolution  $\nu^\sigma = \nu * \gamma_\sigma$  where  $\gamma_\sigma$  denotes the centered Gaussian distribution with covariance matrix  $\sigma^2 Id$ . In other words  $\nu^\sigma$  is the probability distribution of  $X + \sigma G$  where  $X$  is a random variable with distribution  $\nu$  and  $G$  is a standard Gaussian variable. A natural question is to know when  $\nu^\sigma$  satisfies a Poincaré or a log-Sobolev inequality and to get some controls on the corresponding constants. Notice that  $\nu$  is not assumed to satisfy itself such an inequality.

When  $\nu$  has compact support, included in the Euclidean ball  $B(0, R)$ , this question has been partly studied in [48], and the results therein extended in [7]. [48] is using the Lyapunov function method of [21], while [7] is partly using the Bakry-Emery criterion. Since for  $\sigma > 0$ ,  $\nu^\sigma(dx) = e^{-V^\sigma(x)} dx$  for some smooth  $V^\sigma$ , with

$$V^\sigma(x) = - \ln \left( \int e^{-\frac{|x-y|^2}{2\sigma^2}} (2\pi\sigma^2)^{-n/2} \nu(dy) \right),$$

a simple calculation (see [7] p.438) shows that

$$\nabla^2 V^\sigma \geq \left( \frac{1}{\sigma^2} - \frac{R^2}{\sigma^4} \right) Id, \tag{2.15}$$

so that  $C_{LS}(\nu^\sigma) \leq \frac{2\sigma^4}{\sigma^2 - R^2}$  according to (1.6), as soon as  $\sigma > R$  (it seems that the factor 2 is lacking in [7]). The small variance case is more delicate and imposes to use other arguments. Nevertheless it is not difficult using a variance decomposition to prove that the following is always true:

$$C_P(\nu^\sigma) \leq \sigma^2 e^{4R^2/\sigma^2}.$$

If  $V^\sigma$  is not necessarily strongly convex, the Hessian remains bounded from below. Using deep results by E. Milman ([42]) in the spirit of the ones we will recall in the next section (see Theorem 3.2), it is shown in Theorem 4.3 of [7] that  $\nu^\sigma$  is still satisfying a log-Sobolev inequality that does not depend on the dimension  $n$  provided  $\sigma > R/\sqrt{2}$ , but this time the log-Sobolev constant is not explicit.

We will improve the latter result and furnish an explicit constant by directly using our perturbation results. To this end we simply write

$$\nu^\sigma(dx) = Z^{-1} e^{-F(x)} \gamma_\sigma(dx) \quad \text{with} \quad F(x) = V^\sigma(x) - \frac{|x|^2}{2\sigma^2}. \quad (2.16)$$

We have

$$\nabla F(x) = \int \frac{x-y}{\sigma^2} h(x,y) \nu(dy) - \frac{x}{\sigma^2} = - \int \frac{y}{\sigma^2} h(x,y) \nu(dy) \quad (2.17)$$

where

$$h(x,y) = \frac{e^{-\frac{|x-y|^2}{2\sigma^2}}}{\int e^{-\frac{|x-z|^2}{2\sigma^2}} \nu(dz)}.$$

Hence

$$|\nabla F(x)| \leq \frac{1}{\sigma^2} \int |y| h(x,y) \nu(dy) \leq \frac{R}{\sigma^2}. \quad (2.18)$$

It remains to apply Theorem 2.1, part (2) of Theorem 2.5 and the bounds  $C_P(\gamma_\sigma) \leq \sigma^2$  and  $C_{LS}(\gamma_\sigma) \leq 2\sigma^2$  in order to get

**Theorem 2.9.** *Let  $\nu$  be any probability measure whose support is included in  $B(0, R)$ . Define  $\nu^\sigma = \nu * \gamma_\sigma$  where  $\gamma_\sigma$  denotes the centered Gaussian distribution with covariance matrix  $\sigma^2 \text{Id}$ . Then if*

$$s := \frac{1 + \varepsilon}{4} \frac{R^2}{\sigma^2} < 1$$

it holds

$$C_P(\nu^\sigma) \leq \frac{1 + \varepsilon^{-1}}{1 - s} \sigma^2.$$

Similarly for all  $\theta$  and  $\beta$  positive,

$$\begin{aligned} C_{LS}(\nu^\sigma) &\leq \left( \frac{2(\beta+1)(1+\theta^{-1})}{\beta} + 5 \frac{1+\varepsilon^{-1}}{1-s} \right) \sigma^2 \\ &\quad + 2 \frac{1+\varepsilon^{-1}}{1-s} \left( \frac{(1+\theta)(1+\beta)}{4\beta} + \frac{\beta^2}{2} \right) R^2. \end{aligned}$$

Notice that this result covers the range  $\sigma > R/2$  which is larger than the one in [7]. Here we have used  $\mu(F) \leq F(0) + L\mu(|x|)$  in order to simplify the (already intricate) bound for the log-Sobolev constant. Using the elementary general  $C_P(\mu * \nu) \leq C_P(\mu) + C_P(\nu)$ , one has  $C_P(\nu^\sigma) \leq C_P(\nu^{\sigma_0}) + (\sigma - \sigma_0)^2$ , yielding the correct asymptotic behaviour.

The previous proof can easily be extended to more general situations replacing  $\gamma_\sigma$  by some more general  $\mu(dx) = e^{-H(x)} dx$ , provided  $\nabla^2 H$  is bounded, yielding

**Theorem 2.10.** *Let  $\nu$  be any probability measure whose support is included in  $B(0, R)$ . Define  $\nu^H = \nu * \mu$  where  $\mu(dx) = e^{-H(x)} dx$  is a probability measure such that*

$$\sup_x |\nabla^2 H(x)| = K < +\infty.$$

Then if

$$s := \frac{1 + \varepsilon}{4} K^2 R^2 C_P(\mu) < 1$$

it holds

$$C_P(\nu^H) \leq \frac{1 + \varepsilon^{-1}}{1 - s} C_P(\mu).$$

**Proof.** Following the notations of the previous proof we have

$$\nabla F(x) = \int \nabla H(x - y) h(x, y) \nu(dy) - \nabla H(x) = \int (\nabla H(x - y) - \nabla H(x)) h(x, y) \nu(dy)$$

with

$$h(x, y) = \frac{e^{-H(x-y)}}{\int e^{-H(x-z)} \nu(dz)}.$$

It remains to use

$$|\nabla H(x - y) - \nabla H(x)| \leq K |y|,$$

and to use Theorem 2.1. □

**Corollary 2.11.** *Let  $X$  be a random variable supported by  $B(0, R)$  and  $Y$  a random variable with distribution  $\mu(dx) = e^{-H(x)} dx$  such that  $\sup_x |\nabla^2 H(x)| = K < +\infty$ . For  $\sigma \in \mathbb{R}^+$  define  $X^\sigma = X + \sigma Y$  and denote by  $\nu^{\sigma H}$  the distribution of  $X^\sigma$ . Then if*

$$s := \left( \frac{1 + \varepsilon}{4} K^2 R^2 C_P(\mu) \right) / \sigma^2 < 1$$

it holds

$$C_P(\nu^H) \leq \frac{1 + \varepsilon^{-1}}{1 - s} C_P(\mu) \sigma^2.$$

**Proof.** It is enough to remark that the probability density of  $\sigma Z$  is proportional to  $e^{-H(x/\sigma)}$  so that  $|\nabla F| \leq \frac{KR}{\sigma^2}$  and to remember that  $C_P(\sigma Y) = \sigma^2 C_P(Y)$ . □

**Remark 2.12.** One can ask about what happens when  $\nabla H$  is bounded, for instance if  $\mu(dx) = Z^{-1} e^{-\sigma|x|} dx$  in  $\mathbb{R}$ . The proofs above furnish  $|\nabla F| \leq 2/\sigma$  for all  $R$  and all  $\sigma$  so that the condition on  $s$  reads  $s := (1 + \varepsilon) C_P(\mu) < 1$  which is impossible since  $C_P(\mu) = 4$ . This is another argument showing that our perturbation result is close to be optimal.  $\diamond$

### 3. Perturbation with log-concavity.

We will now give some new results relating the Poincaré constant of both measures  $\mu$  and  $\mu_F$  when at least one of them is log-concave.

For log-concave distributions it is often more convenient to use the Cheeger constant instead of the Poincaré constant. Recall the following

**Proposition 3.1.** Recall that in all cases  $C_P(\mu) \leq 4(C'_C)^2(\mu) \leq 4C_C^2(\mu)$ . If in addition  $\mu$  is log-concave the following converse inequality is satisfied :

$$C'_C(\mu) \leq C_C(\mu) \leq \frac{16}{\pi} \sqrt{C_P(\mu)}.$$

The first inequality is contained in [11], while the second one is shown in [19] proposition 9.2.11. With the slightly worse constant 6 instead of  $\frac{16}{\pi}$  the result is due to Ledoux in [37].

A remarkable property of log-concave measures, we shall intensively use in the sequel, is that a very weak form of the Poincaré (or Cheeger) inequality is enough to imply the true one. For simplicity we recall here the two main results we obtained in [19] (Theorem 9.2.7 and Theorem 9.2.14), improving on the beautiful seminal result by E. Milman ([41])

**Theorem 3.2.** *Let  $\nu$  be a log-concave probability measure.*

- (1) *Assume that there exists some  $0 \leq s < 1/2$  and some  $\beta(s)$  such that for any Lipschitz function  $f$  it holds*

$$\nu(|f - m_\nu(f)|) \leq \beta(s) \|\nabla f\|_\infty + s \text{Osc}(f).$$

*Then*

$$C'_C(\nu) \leq \frac{4\beta(s)}{\pi(\frac{1}{2} - s)^2}.$$

- (2) *Assume that there exists some  $0 \leq s < 1/6$  and some  $\beta(s)$  such that for any Lipschitz function  $f$  it holds*

$$\text{Var}_\nu(f) \leq \beta(s) \nu(|\nabla f|^2) + s \text{Osc}^2(f).$$

*Then*

$$C'_C(\nu) \leq \frac{4\sqrt{\beta(s) \ln 2}}{1 - 6s}.$$

*In both cases recall that  $C_P(\nu) \leq 4(C'_C(\nu))^2$ .*



### 3.1. From Holley-Stroock to Barthe-Milman.

We start by mimicking the proof of Holley-Stroock perturbation result replacing the uniform bound on  $F$  by exponential moments

**Proposition 3.3.** If  $\mu_F$  is log-concave then

$$C'_C(\mu_F) \leq \frac{32}{\pi} \frac{\mu^{1/2}(e^{-2F})}{\mu(e^{-F})} C_P^{1/2}(\mu)$$

so that

$$C_P(\mu_F) \leq \frac{4 \times 32^2}{\pi^2} \frac{\mu(e^{-2F})}{\mu^2(e^{-F})} C_P(\mu).$$

**Proof.** Let  $f$  be a bounded Lipschitz function. Of course in the definition of  $\mu_F$  we may always replace  $e^{-F}$  by  $e^{-(F-\min F)}$  provided  $F$  is bounded from below. Hence, for simplicity we may first assume that  $F \geq 0$  so that  $e^{-F} \leq 1$  is in all the  $\mathbb{L}^p(\mu)$ . Then :

$$\begin{aligned} \mu_F(|f - \mu_F(f)|) &\leq 2\mu_F(|f - m_{\mu_F}(f)|) \leq 2\mu_F(|f - \mu(f)|) \\ &\leq 2 \frac{\mu(|f - \mu(f)| e^{-F})}{\mu(e^{-F})} \\ &\leq \frac{2}{\mu(e^{-F})} \mu^{1/2}(|f - \mu(f)|^2) \mu^{1/2}(e^{-2F}) \\ &\leq \frac{2\mu^{1/2}(e^{-2F})}{\mu(e^{-F})} C_P^{1/2}(\mu) \mu^{1/2}(|\nabla f|^2) \leq \frac{2\mu^{1/2}(e^{-2F})}{\mu(e^{-F})} C_P^{1/2}(\mu) \|\nabla f\|_\infty. \end{aligned}$$

If  $F$  is not bounded from below, just using a cut-off (for instance reducing the support of  $\mu_F$  to a large Euclidean ball) we obtain the same result, with a possibly infinite right hand side. Using Theorem 3.2 (1), we can thus conclude  $\square$

Compared to Holley-Stroock theorem, the previous result is interesting but far from optimal. As shown in Theorem 2.7 of [9], there exists an universal constant  $c$  such that

$$C_P(\mu_F) \leq c \left( 1 + \ln \left( \frac{\mu^{1/2}(e^{-2F})}{\mu(e^{-F})} \right) \right)^2 C_P(\mu).$$

We shall recover this result and furnish a numerical bound for  $c$ . To this end first recall

**Definition 3.4.** The concentration profile of a probability measure  $\nu$  denoted by  $\alpha_\nu$ , is defined as

$$\alpha_\nu(r) := \sup \left\{ 1 - \nu(A + B(0, r)); \nu(A) \geq \frac{1}{2} \right\}, \forall r > 0,$$

where  $B(y, r)$  denotes the Euclidean ball centered at  $y$  with radius  $r$ .

The following is shown in [19] Corollary 9.2.10

**Proposition 3.5.** For any log-concave probability measure  $\nu$ ,

$$C'_C(\nu) \leq \inf_{0 < s < \frac{1}{4}} \frac{16 \alpha_\nu^{-1}(s)}{\pi(1-4s)^2} \quad \text{and} \quad C_P(\nu) \leq \inf_{0 < s < \frac{1}{4}} \left( \frac{32 \alpha_\nu^{-1}(s)}{\pi(1-4s)^2} \right)^2.$$

Actually a better result, namely

$$C'_C(\nu) \leq \frac{\alpha_\nu^{-1}(s)}{1-2s}$$

that holds for all  $s < \frac{1}{2}$  was shown by E. Milman in Theorem 2.1 of [43], when  $\nu$  is the uniform measure on a convex body. The result extends presumably to any log-concave measure, but the proof of this result lies on deep geometric results (like the Heintze-Karcher theorem) while ours is elementary.

Now recall the statement of Proposition 2.2 in [9], in a simplified form: if  $M = \frac{\mu^{\frac{1}{2}}(e^{-2F})}{\mu(e^{-F})}$ , then

$$\alpha_{\mu_F} \leq 2M \alpha_\mu^{1/2}(r/2).$$

We may use this result together with proposition 3.5 to deduce corollary 9.3.2 in [19]

**Corollary 3.6.** If  $\mu_F$  is log-concave, denoting  $M = \frac{\mu^{\frac{1}{2}}(e^{-2F})}{\mu(e^{-F})}$ ,

$$C'_C(\mu_F) \leq \inf_{0 < s < \frac{1}{4}} \frac{32 \alpha_\mu^{-1}((s/2M)^2)}{\pi(1-4s)^2}.$$

Finally, as shown by Gromov and V. Milman, the concentration profile of a measure with a finite Poincaré constant is exponentially decaying. One more time it is not as easy to find an explicit version of Gromov-Milman's result. We found two of them in the literature: the first one in [46] Théorème 25 (in french) and Proposition 11

$$\alpha_\mu(r) \leq 16 e^{-\frac{r}{\sqrt{2C_P(\mu)}}},$$

the second one in [10] Theorem 2:

$$\alpha_\mu(r) \leq e^{-\frac{r}{3\sqrt{C_P(\mu)}}}.$$

We shall use the first one due to the lower constant and obtain

**Theorem 3.7.** If  $\mu_F$  is log-concave, denoting  $M = \frac{\mu^{\frac{1}{2}}(e^{-2F})}{\mu(e^{-F})}$ , for all  $0 < s < \frac{1}{4}$ ,

$$C_P(\mu_F) \leq \frac{C}{(1-4s)^4} (3 \ln 2 + \ln(1/s) + \ln M)^2 C_P(\mu)$$

where the constant  $C$  satisfies  $C \leq \left( \frac{64\sqrt{2}}{\pi} \right)^2$ .

**Remark 3.8.** Of course  $M \geq 1$ . In order to get a presumably more tractable bound, first remark that adding a constant to  $F$  does not change  $M$  so that, provided  $F$  is bounded from below (otherwise we

may argue as in the proof of proposition 3.3), we may always assume that  $\min F = 0$ . It thus follows  $\mu(e^{-2F}) \leq \mu(e^{-F})$ . According to Jensen's inequality we also have  $\mu(e^{-F}) \geq e^{-\mu(F)}$  so that finally

$$M \leq e^{\frac{1}{2}\mu(F)}.$$

Finally

$$C_P(\mu_F) \leq (C_1 + C_2 \mu^2(F)) C_P(\mu)$$

for some explicit constants  $C_1$  and  $C_2$ . Choosing for instance  $s = 1/8$  in theorem 3.7, we get

$$C_P(\mu_F) \leq \frac{2^{17}}{\pi^2} (4 \ln(2) + \frac{1}{2} \mu(F))^2. \quad (3.1)$$

◇

**Example 4. Gaussian perturbation.**

For  $\rho \in \mathbb{R}^+$  consider

$$\mu^\rho(dx) = Z_\rho^{-1} e^{-V(x) - \frac{1}{2} \rho |x|^2} dx,$$

i.e.

$$F(x) = \frac{\rho}{2} |x|^2.$$

If  $V$  is convex (hence  $\mu^\rho$  log-concave), we deduce from the previous theorem and Bakry-Emery criterion that,

$$C_P(\mu^\rho) \leq \min \left( \frac{1}{\rho}; \left( C_1 + C_2 \ln^2 \left( \frac{\mu^{\frac{1}{2}}(e^{-\rho|x|^2})}{\mu(e^{-\rho|x|^2/2})} \right) \right) C_P(\mu) \right),$$

for some explicit universal constants  $C_1$  and  $C_2$ . This indicates that we can find a bound for the Poincaré constant of  $\mu^\rho$  that does not depend on  $\rho$ . We shall try to get some explicit result.

First according to remark 3.8, the  $\ln^2$  can be bounded up to some universal constant by  $\mu^2(\rho|x|^2)$  so that if  $\mu$  is isotropic i.e. is centered with a covariance matrix equal to identity, we obtain a bound in  $\rho^2 n^2$ . Optimizing in  $\rho$ , we see that the worst case is for  $\rho \sim n^{-\frac{2}{3}} C_P^{-\frac{1}{3}}(\mu)$  yielding

**Proposition 3.9.** There exists som universal constant  $C$  such that, if  $\mu(dx) = Z^{-1} e^{-V(x)} dx$  is log-concave and isotropic, then for all  $\rho \geq 0$ , and  $\mu^\rho(dx) = Z_\rho^{-1} e^{-V(x) - \frac{1}{2} \rho |x|^2} dx$ ,

$$C_P(\mu^\rho) \leq C n^{2/3} C_P^{\frac{1}{3}}(\mu).$$

This result looks disappointing. A direct approach via the KLS inequality obtained by Chen will presumably give a better dimensional bound provided we are able to get a good bound for the covariance matrix of the perturbed measure. When  $V$  is even we can directly estimate the covariance matrix of  $\mu^\rho$  as in Theorem 18 in [8]. ◇

### 3.2. Smooth perturbations.

The previous subsection was devoted to recover with explicit constants the results by Barthe and Milman. Using the specific properties of log-concave measures we will now use the ideas of section 2. In particular we will replace uniform bounds (on  $|\nabla F|$  for example) by moments of similar quantities. Let us state a first result in this direction

**Theorem 3.10.** *If  $\mu_F$  is log-concave we have for  $\varepsilon > 0$ ,*

$$C_P(\mu_F) \leq \frac{64 \ln(2) (1 + \varepsilon^{-1}) C_P(\mu)}{(1 - 6s)^2} \quad \text{provided} \quad \frac{1 + \varepsilon}{4} C_P(\mu) \mu_F(|\nabla F|^2) := s < \frac{1}{6}.$$

In particular this bound is available as soon as  $F$  is  $L$ -Lipschitz with  $L^2 < 2/(3(1 + \varepsilon)C_P(\mu))$ , but this condition is worse than the one in theorem 2.1. The gain of the previous theorem is thus to replace the uniform bound on  $\nabla F$  by a bound on its second moment.

*Proof.* We deduce from (2.3)

$$\text{Var}_{\mu_F}(f) \leq (1 + \varepsilon^{-1}) C_P(\mu) \mu_F(|\nabla f|^2) + \frac{1 + \varepsilon}{4} C_P(\mu) \text{Osc}^2(f) \mu_F(|\nabla F|^2). \quad (3.2)$$

So that, if  $\mu_F$  is log-concave, using Theorem 3.2 (2) we get the result.  $\square$

We may similarly modify the proof of proposition 2.2 to similarly get a Cheeger inequality.

**Theorem 3.11.** *If  $\mu_F$  is log-concave we have*

$$C'_C(\mu_F) \leq \frac{16 C_C(\mu)}{\pi(1 - 2s)^2} \quad \text{provided} \quad C_C(\mu) \mu_F(|\nabla F|) := s < \frac{1}{2}.$$

*In particular if  $\mu$  is also log-concave we have*

$$C_P(\mu_F) \leq \frac{256 \times 64}{\pi^4} \frac{C_P(\mu)}{(1 - 2s)^4},$$

*for  $s$  as before.*

At the level of Cheeger inequality, we found no other comparable perturbation result despite, once again, the very simple argument involved here.

Starting with (2.5) we also have

$$\text{Var}_{\mu_F}(f) \leq \mu_F((f - a)^2) \leq C_P(\mu) \mu_F(|\nabla f|^2) + \frac{1}{2} C_P(\mu) \text{Osc}^2(f) \mu_F([\mathbf{A}F - \frac{1}{2} |\nabla F|^2]^+)$$

so that we obtain an improvement of Theorem 2.3

**Theorem 3.12.** *If  $V$  is  $C^1$ ,  $\mu_F$  is log-concave,  $F$  is of  $C^2$  class and satisfies*

$$C_P(\mu) \mu_F([\mathbf{A}F - \frac{1}{2} |\nabla F|^2]^+) := s < \frac{1}{3}$$

then

$$C_P(\mu_F) \leq \frac{64 \ln(2)}{(1-3s)^2} C_P(\mu).$$

**Remark 3.13.** As in the previous section it is interesting to extend the result to compactly supported log-concave measures. We will thus assume that the set  $U = \{V < +\infty\}$  is convex with a smooth boundary and that  $V$  is  $C^1$  in  $U$ . Then the conclusion of the previous Theorem is still available provided in addition  $\partial_n F \geq 0$  on  $\partial U$ .  $\diamond$

**Remark 3.14.** Assume that  $V$  is  $C^2$  on  $\mathbb{R}^n$  and satisfies  $\langle u, \nabla^2 V(x)u \rangle \geq \rho|u|^2$  for all  $u$  and  $x$  in  $\mathbb{R}^n$ . Let  $U$  be an open convex subset given by  $U = \{W < 1\}$  where  $W$  is a smooth (say  $C^2$ ) convex function. Consider

$$\mu(dx) = Z^{-1} e^{-V(x)} \mathbf{1}_{W(x) \leq 1} dx.$$

It turns out that once again

$$C_P(\mu) \leq 1/\rho.$$

To prove this result one can use the  $\Gamma_2$  theory of Bakry-Emery, but one has to carefully define the algebra  $\mathcal{A}$  (see section 1.16 in [5]). The devil is in this definition when looking at reflected semi-groups. We prefer to give an elementary proof of what we claimed.

Define  $H(x) = (W(x) - 1)^4 \mathbf{1}_{W(x) \geq 1}$ .  $H$  is smooth and

$$\partial_{ij}^2 H(x) = 4(W(x) - 1)^2 \left( (W - 1) \partial_{ij}^2 W + 3 \partial_i W \partial_j W \right) \mathbf{1}_{W(x) \geq 1}$$

so that

$$\langle u, \nabla^2 H(x)u \rangle = 4(W(x) - 1)^2 \left( (W - 1) \langle u, \nabla^2 W(x)u \rangle + 3 \langle u, \nabla W \rangle^2 \right) \mathbf{1}_{W(x) \geq 1} \geq 0.$$

If we consider

$$\mu^\varepsilon(dx) = Z_\varepsilon^{-1} e^{-V(x) - \frac{1}{\varepsilon} H(x)} dx,$$

$\mu^\varepsilon$  satisfies the Bakry-Emery criterion so that  $C_P(\mu^\varepsilon) \leq 1/\rho$ . It remains to let  $\varepsilon$  go to 0 and to pass to the limit in (1.1) by using Lebesgue convergence theorem.  $\diamond$

Let us finish by proving the same type of result at the level of Brascamp-Lieb inequality. Let us consider  $d\mu = e^{-V} dx$  with  $\nabla^2 V > 0$  in the sense of positive definite matrix, the celebrated Brascamp-Lieb inequality is then

$$\text{Var}_\mu(f) \leq \int (\nabla f)^t (\nabla^2 V)^{-1} \nabla f d\mu.$$

We will see that we can easily have some perturbation result, leading to some modified Brascamp-Lieb inequality.

**Theorem 3.15.** *Let us consider  $d\mu = e^{-V} dx$  with  $\nabla^2 V > 0$  in the sense of positive definite matrix and suppose that there exists  $\varepsilon$  such that*

$$\frac{1}{4}(1 + \varepsilon) \|(\nabla F)^t (\nabla^2 V)^{-1} \nabla F\|_\infty < 1$$

then

$$\text{Var}_{\mu_F}(f) \leq \frac{(1 + \varepsilon^{-1})}{1 - \frac{1}{4}(1 + \varepsilon) \|(\nabla F)^t (\nabla^2 V)^{-1} \nabla F\|_\infty} \int \nabla f^t (\nabla^2 V)^{-1} \nabla f d\mu.$$

In particular if  $\mu_F$  is log-concave then

$$C_P(\mu_F) \leq \frac{64(1 + \varepsilon^{-1}) \ln(2)}{1 - \frac{1}{4}(1 + \varepsilon) \|(\nabla F)^t (\nabla^2 V)^{-1} \nabla F\|_\infty} \int \|(\nabla^2 V)^{-1}\|_{HS} d\mu$$

where  $\|\cdot\|_{HS}$  denotes the Hilbert-Schmidt norm.

**Proof.** We follow the idea already developed for the Poincaré inequality.

$$\text{Var}_{\mu_F}(f) \leq \mu_F((f - a)^2) = \mu^{-1}(e^{-F}) \mu \left( ((f - a)e^{-\frac{1}{2}F})^2 \right)$$

for which we choose

$$a = \frac{\mu \left( f e^{-\frac{1}{2}F} \right)}{\mu \left( e^{-\frac{1}{2}F} \right)}$$

so that we may apply Brascamp-Lieb inequality for  $\mu$ .

$$\begin{aligned} \mu_F((f - a)^2) &\leq \int (\nabla f - \frac{1}{2}(f - a) \nabla F)^t (\nabla^2 V)^{-1} (\nabla f - \frac{1}{2}(f - a) \nabla F) d\mu_F \\ &\leq (1 + \varepsilon^{-1}) \int (\nabla f)^t (\nabla^2 V)^{-1} \nabla f d\mu_F \\ &\quad + \frac{1}{4}(1 + \varepsilon) \int (f - a)^2 (\nabla F)^t (\nabla^2 V)^{-1} \nabla F d\mu_F. \end{aligned}$$

We then use our growth condition to conclude.

If  $\mu_F$  is log-concave, we have

$$\text{Var}_{\mu_F}(f) \leq \frac{(1 + \varepsilon^{-1})}{1 - \frac{1}{4}(1 + \varepsilon) \|(\nabla F)^t (\nabla^2 V)^{-1} \nabla F\|_\infty} \|\nabla f\|_\infty^2 \int \|(\nabla^2 V)^{-1}\|_{HS} d\mu,$$

and we can conclude by using Theorem 3.2 (2).  $\square$

Of course it is illusory to expect a Brascamp-Lieb inequality for  $\mu_F$  as  $\nabla^2(V + F)$  is not necessarily positive. However it may be useful for concentration inequalities, indeed, reproducing the proof of the exponential integrability for Poincaré inequality due to Bobkov-Ledoux, see [5], under the assumptions of the previous theorem, if  $f$  is such that

$$\|(\nabla f)^t (\nabla^2 V)^{-1} \nabla f\|_\infty \leq 1$$

then

$$\forall s < \sqrt{\frac{4 \left(1 - \frac{1}{4}(1 + \epsilon)\|(\nabla F)^t (\nabla^2 V)^{-1} \nabla F\|_\infty\right)}{1 + \epsilon^{-1}}}, \quad \int e^{sf} d\mu_F < \infty.$$

This implies exponential concentration for  $\mu_F$  for some particular class of functions.

### 3.3. Coming back: from the perturbed measure to the initial one.

Any probability measure  $\mu(dx) = e^{-V(x)} dx$  can be seen as a perturbation of a perturbed measure, namely  $\mu(dx) = Z^{-1} e^{F(x)} \mu_F(dx)$ . In some cases the measure  $\mu_F$  is simpler to study, so that one can expect some results for the initial one using our perturbation method. We already discussed an example 3 comparing a compactly supported  $\mu$  with the uniform distribution on its support. In this section we will compare a log-concave distribution with another one which is “more” log-concave.

For  $p \geq 1$  and  $\lambda > 0$  consider

$$\nu(dx) = Z^{-1} e^{-V(x) - \lambda^p \sum_{i=1}^n |x_i|^p} dx,$$

where we assume that  $V$  is a convex function. If we denote

$$F(x) = -\lambda^p \sum_{i=1}^n |x_i|^p$$

we have

$$\mu(dx) = \nu_F(dx).$$

We thus have

$$|\nabla F|^2(x) \leq \lambda^{2p} p^2 \sum_{i=1}^n |x_i|^{2(p-1)}, \quad (3.3)$$

so that

$$\mu(|\nabla F|^2) = \nu_F(|\nabla F|^2) \leq \lambda^{2p} p^2 \mu \left( \sum_{i=1}^n |x_i|^{2(p-1)} \right). \quad (3.4)$$

Choosing for simplicity  $\epsilon = 1$  and  $s = 1/12$  in Theorem 3.10 we thus have to choose (if this choice is possible)

$$\lambda^{2p} p^2 C_P(\nu) \mu \left( \sum_{i=1}^n |x_i|^{2(p-1)} \right) \leq \frac{1}{6}. \quad (3.5)$$

Notice that we may always use an upper bound for  $C_P(\nu)$  furnishing a lower bound for  $\lambda$  and an upper bound for  $C_P(\mu)$ .

Hence the first thing is to get some explicit bound for  $C_P(\nu)$ . For  $p = 2$  we may use Bakry-Emery criterion, while for  $1 \leq p \leq 2$  we may use Theorem 1.4 and for  $p \geq 2$  we may use Theorem 17 in [8]. In the last two cases some additional assumptions on  $V$  are necessary.

**Theorem 3.16.** *Let  $\mu(dx) = Z^{-1} e^{-V(x)} dx$  be a log-concave distribution. Then*

- (i) it holds  $C_P(\mu) \leq 32 \times 81 \ln(2) \mu(|x - \mu(x)|^2)$ .  
(ii) If in addition  $V$  is even we have for all  $1 < p \leq 2$ ,

$$C_P(\mu) \leq C 6^{\frac{1}{p-1}} p^{\frac{2}{p-1}} \mu^{\frac{1}{p-1}} \left( \sum_{i=1}^n |x_i|^{2(p-1)} \right) (\ln(n))^{\frac{2-p}{p-1}}, \quad (3.6)$$

for some universal  $C$ .

- (iii) If  $V$  is unconditional, i.e.  $V(x_1, \dots, x_n) = V(|x_1|, \dots, |x_n|)$  for all  $x$ , and  $n \geq 2$ , it holds

$$C_P(\mu) \leq C \ln^2(3n) \sigma^2(\mu)$$

with  $C = 512 e^2 \ln(2)$  (recall that  $\sigma^2(\mu)$  has been introduced in (1.7)).

**Proof.** *The case  $p = 2$ .* The situation here is particular since according to the Bakry-Emery criterion  $C_P(\nu) \leq \frac{1}{2\lambda^2}$  for all  $\lambda > 0$ . Hence we may apply Theorem 3.10 as soon as  $s := 2\mu(|x|^2)\lambda^2 \leq 1/6$  yielding

$$C_P(\mu) \leq 128 \ln(2) \frac{1}{2s(1-6s)^2} \mu(|x|^2).$$

The optimal choice of  $s$  is  $1/18$  and of course we may always center  $\mu$  without changing the Poincaré constant. Hence the result (i).

*The case  $p < 2$ .* In this case assuming that  $V$  is even we know that  $C_P(\nu) \leq C \lambda^{-2} (\ln n)^{\frac{2-p}{p}}$  since  $C_P(Z^{-1} e^{-|x_i|^p})$  is bounded above by an universal constant for  $1 \leq p \leq 2$ . (3.5) then furnishes (3.6) for  $p \neq 1$ .

*The case  $p > 2$ .* When  $V$  (or  $\mu$ ) is unconditional, it follows from Theorem 17 in [8] and the dilation property of the Poincaré constant, that

$$C_P(\nu) \leq \lambda^{-2} C_P(S_p(dx_1))$$

for all  $\lambda$ , where

$$S_p(dx_1) = \frac{1}{z_p} e^{-|x_1|^p}.$$

According to Bobkov's one dimensional result ([11] Corollary 4.3),

$$C_P(S_p) \leq 12 \text{Var}_{S_p}(x) = 12 \frac{\Gamma(3/p)}{\Gamma(1/p)}. \quad (3.7)$$

A better result is obtained by combining [12] Theorem 2.1 and the dilation property of Poincaré constants yielding

$$C_P(S_p) \leq \frac{p^{1-2/p}}{2(1+p)^{1-2/p}}.$$

Furthermore, since  $V$  is even,  $\mu(x) = 0$  and for  $2 < p$ ,

$$\mu(|x_i|^{2(p-1)}) \leq \frac{\Gamma(2p+1)}{2^{p-1}} \mu^{p-1}(|x_i|^2) \leq (p-1)^{2(p-1)} \mu^{p-1}(|x_i|^2) \quad (3.8)$$



according to [33] corollary 5.7 and remark 5.8 (also see in [36] the discussion after definition 2). Hence in the unconditional case,

$$\begin{aligned} C_P(\mu) &\leq C 6^{\frac{1}{p-1}} \left(\frac{p^2}{2}\right)^{\frac{1}{p-1}} (p-1)^2 n^{\frac{1}{p-1}} \frac{p^{\frac{p-2}{p-1}}}{(1+p)^{\frac{p-2}{p-1}}} \sigma^2(\mu) \\ &\leq 4C 3^{\frac{1}{p-1}} (p-1)^2 n^{\frac{1}{p-1}} \sigma^2(\mu), \end{aligned} \quad (3.9)$$

with  $C = 512 \ln(2)$ . Here we used  $p^{2/(p-1)} \leq 4$  for  $p \geq 2$ . It remains to optimize in  $p$ , the optimal value being  $p-1 = \ln(3n)/2$  (which is larger than 1 for  $n \geq 2$ ).  $\square$

**Remark 3.17.** The result (i) is well known and according to [2] p.11 is contained in [34] with a much better pre-constant 4 (also see [11] (1.8) with a non explicit constant). Applying Theorem 3.2 it is easily seen (see (9.2.13) in [19]) that

$$C'_C(\mu) \leq \frac{16}{\pi} \mu(|x - m_\mu(x)|) \leq \frac{16}{\pi} \mu(|x - \mu(x)|). \quad (3.10)$$

If we replace Theorem 3.10 by Theorem 3.11 we obtain using our perturbation method the worse

$$C'_C(\mu) \leq \frac{100\sqrt{10}}{\pi^2} \mu(|x - \mu(x)|).$$

$\diamond$

**Remark 3.18.** (3.6) is not really satisfactory. For instance for  $p = 3/2$  we get a worse result than the one deduced from (3.10) and  $C_P \leq 4(C'_C)^2$  since some dimension dependence appears (actually Barthe and Klartag do not know whether the  $\ln(n)$  pre-constant is really necessary). In addition for log-concave measures all moments can be compared so that even for  $p < 3/2$  the result does not seem as interesting.  $\diamond$

**Remark 3.19.** The result (iii) is not (completely) new, and is due to Klartag in [35] with a non explicit constant. Another proof (still with a difficult to trace constant) is contained in [19]. The constant here is explicit (but certainly far to be sharp), but may be more interesting is the fact that this result can be obtained via Subbotin perturbation.

One may ask whether the previous method extends to non unconditional log-concave distributions. Unfortunately, [8] contains an example (see subsection 3.4) where  $C_P^{\frac{p}{p-1}}(\nu)$  behaves like  $n^{\frac{p-2}{p-1}}$ , but for a measure  $\mu$  which is highly non isotropic. In order to make a step in the direction of the KLS conjecture it should be interesting to get some version of the unconditional result for isotropic distributions.  $\diamond$

## 4. Application to some problems in Bayesian statistics.

In this section we shall discuss two classical problems in Bayesian statistics: linear regression and parameter identification. Actually we shall only quickly mention how the results of the present paper can be used in the second example, and we will mainly focus on the first one. More precisely we will compare the rate of convergence of the Langevin Monte-Carlo algorithm suggested in [27] and the one one can obtain via the methods exposed in the present paper. As said in the introduction a lot of works have been devoted during the last ten years to similar problems. In order to avoid the intricacies of the

obtention of a bound for the rate of convergence, a large part of these works assume some uniform convexity in order to use the Bakry-Emery criterion. Only few are dealing with general log-concave measures or perturbations of such measures. [27] is one of the first paper in this direction and we decided to focus on it.

#### 4.1. Sparse linear regression.

In [27] the authors proposed a Bayesian strategy for the liner regression model

$$Y_i = \langle X_i, \lambda^* \rangle + \xi_i \quad (4.1)$$

where  $\lambda^*$  and each  $X_i$  belong to  $\mathbb{R}^k$ ,  $\xi_i$  are i.i.d. scalar noises and  $i = 1, \dots, n$ .  $n$  is thus the size of a sample while  $M$  is the dimension of the predictor. Given a collection of design points  $X_i$  the exponentially weighted aggregate estimator of  $\lambda^*$  is given by

$$\hat{\lambda}_n(X) = \int \lambda \hat{\pi}_{n,\beta}(d\lambda) \quad (4.2)$$

where

$$\hat{\pi}_{n,\beta}(d\lambda) = C \exp\left(-\frac{1}{\beta} \sum_{i=1}^n |Y_i - \langle X_i, \lambda \rangle|^2\right) \pi(d\lambda) \quad (4.3)$$

is the posterior probability distribution associated to the prior  $\pi(d\lambda) = e^{-W(\lambda)} d\lambda$  and the temperature  $\beta$ . Here and in all what follows  $C$  is a normalizing constant that can change from line to line.

In their Theorem 2 they obtain in particular an explicit bound for the  $\mathbb{L}^2$  error, when the prior is (almost) chosen as

$$\pi(d\lambda) = C \prod_{j=1}^k \frac{e^{-\alpha|\lambda_j|}}{(\tau^2 + \lambda_j^2)^2} \mathbf{1}_{\sum_{j=1}^k |\lambda_j| \leq R} d\lambda = e^{-W(\lambda)} d\lambda, \quad (4.4)$$

for some positive  $\alpha$  and  $R$ . This choice is motivated by dimensional reasons when  $k \gg n$  and  $\lambda^*$  is sparse.

In order to compute  $\hat{\lambda}_n$ , they propose to use the ergodic theorem applied to the Langevin diffusion process

$$dL_t = \sqrt{2} dB_t - \nabla W(L_t) dt - \frac{2}{\beta} \sum_{i=1}^n (\langle X_i, L_t \rangle - Y_i) X_i dt \quad (4.5)$$

where  $B_t$  is a standard  $\mathbb{R}^k$  valued Brownian motion, i.e. the  $\mathbb{L}^1$  convergence of  $\frac{1}{t} \int_0^t L_s ds$  to the desired  $\hat{\lambda}_n$  as  $t \rightarrow +\infty$ . The  $\mathbf{1}_{\sum_{j=1}^k |\lambda_j| \leq R}$  is no more considered here, and thus denote

$$\nu_{n,\beta}(d\lambda) = C \exp\left(-\frac{1}{\beta} \sum_{i=1}^n |Y_i - \langle X_i, \lambda \rangle|^2 - \sum_{i=1}^k 2 \ln(\tau^2 + \lambda_i^2)\right) \prod_{j=1}^k e^{-\alpha|\lambda_j|} d\lambda. \quad (4.6)$$

Remark that  $\nu_{n,\beta}$  is not log-concave. To justify some rate of convergence they call upon the Meyn-Tweedie theory (see e.g. [40]) of Foster-Lyapunov functions. Notice that using instead  $\frac{1}{t} \int_t^{2t} L_s ds$  the bound (8) in [27] becomes  $C\theta^t$  for some  $\theta < 1$ .

Let us recall some facts about the part of the Meyn-Tweedie theory the authors are calling upon. For a diffusion process  $X_t$  with infinitesimal generator  $\mathbf{A} = \Delta - \langle \nabla U, \nabla \rangle$  and invariant (symmetric) measure  $d\mu = e^{-U} dx$  with  $U$  smooth enough, we will say that  $G$  is a (Foster-)Lyapunov function if  $G \geq 1$  and there exist  $\kappa > 0$ ,  $b \geq 0$  and  $K$  some connected bounded domain, such that

$$\mathbf{A}G \leq -\kappa G + b\mathbf{1}_K.$$

The existence of a Lyapunov function implies exponential ergodicity in the following sense: denote by  $\mathbb{E}_x$  the law of  $X_t$  conditional to  $X_0 = x$ , then for any measurable  $f$  such that  $f/G \leq 1$ ,

$$|\mathbb{E}_x(f(X_t)) - \mu(f)| \leq CG(x)\theta^t$$

for some constant  $C$  and some  $\theta < 1$  (see [29, 40]). The theory was extended to more general Foster-Lyapunov functions in order to cover non exponential decays (see e.g. [28]).

If the arguments give an exponential rate  $\theta^t$  for some  $\theta < 1$  it is very hard to find an explicit expression for  $\theta$  (actually we do not find any such expression in the literature). The same holds with  $C$ .

During the last ten years (actually a little bit more), we have studied the links between the Meyn-Tweedie theory and functional inequalities like Poincaré and log-Sobolev inequalities. We refer to [4, 3, 22, 18]. The message is that, in the symmetric case we are looking at, a Poincaré inequality is “equivalent” to the existence of a Lyapunov function. In addition one can give some relationship between all constants (see [18]). The one we are interested here is the following: if  $G$  is a Lyapunov function as before, then (see [3, 22])

$$C_P(\mu) \leq \frac{1}{\kappa} (1 + bC_P(\mu_K)) \quad (4.7)$$

where  $\mu_K$  denotes the restriction of  $\mu$  to  $K$  and  $\text{Osc}_K$  denotes the Oscillation on  $K$ . Using Holley-Stroock perturbation argument we have

$$C_P(\mu_K) \leq e^{\text{Osc}_K U} C_P(\mathcal{U}_K)$$

if  $\mathcal{U}_K$  denotes the uniform measure on  $K$ . Recall that if  $K$  is convex, the Payne-Weinberger estimate furnishes  $C_P(\mathcal{U}_K) \leq D^2/\pi^2$  where  $D$  is the diameter of  $K$ . Hence if  $K$  is convex (in [27] the authors are considering  $K = B(0, R)$ ), we get

$$C_P(\mu) \leq \frac{1}{\kappa} \left( 1 + \frac{bD^2 e^{\text{Osc}_K U}}{\pi^2} \right). \quad (4.8)$$

Finally recall that the Poincaré inequality implies

$$\int |\mathbb{E}_x(f(X_t)) - \mu(f)|^2 \mu(dx) \leq e^{-2t/C_P(\mu)} \mu(|f - \mu(f)|^2)$$

for any  $f \in \mathbb{L}^2(\mu)$ . Using the ellipticity of the generator one sees that the distribution of  $X_t$ , starting from  $x$ , has a density  $p(t, x, \cdot)$  w.r.t. Lebesgue measure. A simple argument detailed in [22] (proof of Proposition (3.1)) shows that this density belongs to  $\mathbb{L}^2(\mu)$  and satisfies

$$\|p(1, x, \cdot)\|_{\mathbb{L}^2(\mu)}^2 \leq (2\pi)^{-k/2} e^{C_U} e^{U(x)}$$

where

$$C_U = \sup_y (\Delta U - \frac{1}{2} |\nabla U|^2).$$

We thus have, using

$$\mathbb{E}_x((f - \mu(f))(X_t)) = \int \mathbb{E}_y((f - \mu(f))(X_{t-1})) p(1, x, y) \mu(dy)$$

and Cauchy-Schwarz inequality, that for  $t \geq 1$

$$|\mathbb{E}_x(f(X_t)) - \mu(f)| \leq \frac{e^{(C_U + U(x))/2}}{(2\pi)^{k/4}} e^{-(t-1)/C_P(\mu)} \mu(|f - \mu(f)|^2). \quad (4.9)$$

The authors of [27] propose to use  $G(\lambda) = e^{\gamma|\lambda|}$  as a Lyapunov function ( $|\lambda|$  denotes the Euclidean norm). An immediate computation yields

$$\frac{\mathbf{A}G}{G} = \gamma^2 - \frac{4\gamma}{|\lambda|} \sum_{j=1}^k \frac{\lambda_j^2}{\lambda_j^2 + \tau^2} - \alpha\gamma \frac{\sum_{j=1}^k |\lambda_j|}{|\lambda|} - \frac{2\gamma}{\beta|\lambda|} \sum_{i=1}^n (\langle X_i, \lambda \rangle - Y_i) \langle X_i, \lambda \rangle.$$

It is not clear that  $\gamma = \alpha$  can be chosen (in particular for the  $\lambda$ 's such that  $\langle X_i, \lambda \rangle = 0$  for all  $i$ ). However for  $|\lambda| \geq 1$ , it holds

$$\frac{\mathbf{A}G}{G} \leq \gamma(\gamma - \alpha) + \frac{2}{\beta} \sum_{i=1}^n |Y_i| |X_i|$$

so that, provided  $\alpha > \frac{2}{\beta} \sum_{i=1}^n |Y_i| |X_i|$ , we can choose  $\gamma$  small enough for  $G$  to be a Lyapunov function. It is then not difficult to see that the Oscillation of the potential in the ball  $|\lambda| \leq 1$  exceeds  $\alpha\sqrt{k}$  yielding a large bound in (4.8).

Our goal is thus to get some interesting (and more tractable) bounds for  $C_P(\nu_{n,\beta})$  by using our direct approach. To this end we may use two methods. In what follows  $Z$  is a normalizing constant that may change from line to line.

First, we may write  $\nu_{n,\beta} = \mu_F$  with

$$\mu(d\lambda) = Z^{-1} \exp\left(-\frac{1}{\beta} \sum_{i=1}^n \langle X_i, \lambda \rangle^2\right) \prod_{j=1}^k e^{-\alpha|\lambda_j|} d\lambda$$

and

$$F(\lambda) = -\frac{2}{\beta} \sum_{i=1}^n Y_i \langle X_i, \lambda \rangle - \sum_{i=1}^n \log(\tau^2 + \lambda_i^2) = -\frac{2}{\beta} \sum_{j=1}^k \lambda_j \langle Y, \tilde{X}^j \rangle - \sum_{i=1}^n \ln(\tau^2 + \lambda_i^2)$$

where  $\tilde{X}^j = (X_1^j, \dots, X_n^j)$ . According to Barthe and Klartag result Theorem 1.4, the tensorisation property and the quadratic behaviour of the Poincaré with respect to dilation we have

$$C_P(\mu) \leq C \ln^2(k) \frac{1}{\alpha^2},$$

for some universal constant. Since  $F$  is  $L$ -Lipschitz with

$$L = \frac{2}{\beta} \sup_{j=1, \dots, k} |\langle Y, \tilde{X}^j \rangle| + \frac{1}{\tau}$$

we may apply Theorem 2.1

**Theorem 4.1.** *There exist two universal positive constants  $c$  and  $C$  such that, provided*

$$\frac{\ln(k)}{\beta \alpha} \left( \sup_{j=1, \dots, k} |\langle Y, \tilde{X}^j \rangle| + \beta \tau \right) \leq c$$

then

$$C_P(\nu_{n, \beta}) \leq C \frac{\beta}{\sup_{j=1, \dots, k} |\langle Y, \tilde{X}^j \rangle| + \beta \tau}.$$

Depending on whether the logarithmic factor  $\ln(k)$  is necessary in Theorem 1.4 or not will yield a result that does not depend on the dimension. Notice that the previous result imposes

$$\alpha \beta > C' \ln(k) \sup_{j=1, \dots, k} |\langle Y, \tilde{X}^j \rangle|$$

while the Lyapunov function approach imposes

$$\alpha \beta > C' \sum_i |X_i| |Y_i|,$$

though we did not really try to optimize this bound. In addition if we choose  $\alpha$  big enough the Poincaré constant becomes small, while the bound obtained via the Lyapunov function will explode.

As said before the situation considered here is  $k \gg n$ . Another approach to compute  $C_P(\mu)$  will thus be to use Theorem 3.3 instead of Theorem 1.4. The discussion will thus be very similar to the one in Example 4. Of course because the quadratic form  $\sum_{i=1}^n \langle X_i, \lambda \rangle^2$  is very degenerate on  $\mathbb{R}^k$  it is impossible to use Bakry-Emery criterion. Since our goal is not to rewrite [27] but to see how one can control the rate of convergence of the Langevin dynamics, we will assume for simplicity that the  $X_i$ 's are an orthogonal family. So, after an orthogonal transform, we may write

$$\mu(d\lambda) = Z^{-1} \exp \left( -(1/\beta) \sum_{i=1}^n |X_i|^2 \lambda_i^2 \right) e^{-V(\lambda)} d\lambda$$

where  $e^{-V(\lambda)} d\lambda$  is an orthogonal change of  $\prod_{j=1}^k e^{-\alpha |\lambda_j|} d\lambda$ , hence shares the same Poincaré constant. Since the Poincaré constant of the symmetric exponential isotropic distribution in dimension 1, hence the one of the tensor product of such distributions, is equal to 4, the dilation scaling yields  $C_P(e^{-V(\lambda)} d\lambda) = 4/\alpha^2$ .

Of course we may consider, when  $k \geq n$ ,

$$\eta_n(d\lambda_1, \dots, d\lambda_n) = \left( \int e^{-V(\lambda)} d\lambda_{n+1} \dots d\lambda_k \right) d\lambda_1 \dots d\lambda_n$$

which is a new log-concave distribution according to Prekopa-Leindler theorem. It is still isotropic up to a dilation of scale  $\alpha$  and

$$\int \exp\left(-\frac{1}{\beta} \sum_{i=1}^n |X_i|^2 \lambda_i^2\right) e^{-V(\lambda)} d\lambda = \int \exp\left(-\frac{1}{\beta} \sum_{i=1}^n |X_i|^2 \lambda_i^2\right) \eta_n(d\lambda_1, \dots, d\lambda_n).$$

In addition the Poincaré constant of  $\eta_n$  is less than the one of  $e^{-V} d\lambda$ . We may thus argue as in Example 4 when  $\alpha = 1$ , and then use the dilation property of the Poincaré inequality to conclude that for all  $\beta$  and  $X$ ,

$$C_P(\mu) \leq C \frac{n^{2/3}}{\alpha^2},$$

for some universal constant  $C$ . We can now follow what we did previously to get

**Theorem 4.2.** *There exist two universal positive constants  $c$  and  $C$  such that, provided*

$$\frac{n^{1/3}}{\beta \alpha} \left( \sup_{j=1, \dots, k} |\langle Y, \tilde{X}^j \rangle| + \beta \tau \right) \leq c$$

then

$$C_P(\nu_{n,\beta}) \leq C \frac{\beta}{\sup_{j=1, \dots, k} |\langle Y, \tilde{X}^j \rangle| + \beta \tau}.$$

Hence depending whether  $\ln(k) > C'n^{1/3}$  or not we may use either the bound in Theorem 4.2 or the one in Theorem 4.1

## 4.2. Parameter identification.

We shall briefly indicate another problem, parameter identification via a Bayesian approach as studied in the recent [31]. Given a convex function  $U : \mathbb{R}^q \times \mathbb{R}^d \mapsto \mathbb{R}$  one considers the family of probability densities

$$f_\theta(x) := f(x, \theta) = Z_\theta^{-1} e^{-U(x, \theta)}$$

on  $\mathbb{R}^q$ . Given the observation of a sample  $X = (X_1, \dots, X_n)$  of i.i.d. random vectors with density  $f_{\theta^*}$ , one wants to estimate the unknown parameter  $\theta^*$ , here again using a Bayesian procedure.

For a prior distribution density  $\pi_0(\theta)$ , the posterior density is thus

$$\pi_n(\theta) = \pi_0(\theta) \prod_{i=1}^n f_\theta(X_i), \quad (4.10)$$

and the natural bayesian estimator of  $\theta^*$  is once again given by

$$\hat{\theta}_n = \int \theta \pi_n(\theta) d\theta. \quad (4.11)$$

It is shown in [31] that under mild assumptions this estimator is consistent and a bound for the  $\mathbb{L}^p$  error is given, provided there exists a constant  $C_P^U$  such that

$$C_P(f_\theta(x) dx) \leq C_P^U \quad \text{for all } \theta \in \text{supp}(\pi_0), \quad (4.12)$$

An important case where this assumption is satisfied is the location problem, i.e. when  $U(x, \theta) = V(x - \theta)$  for some convex function  $V$ , and of course  $q = d$ .

Here again the authors propose to use a Langevin Monte Carlo procedure to compute  $\hat{\theta}_n$ , and as in the previous subsection the problem is now to estimate the Poincaré constant of the measure  $\pi_n(\theta) d\theta$ . The situation is of course much simpler here if one chooses  $\pi_0$  as a strictly log-concave log concave distribution since we obtain a bound that only depends on the curvature of  $\pi_0$ . The case of a general log-concave measure  $\pi_0$  is studied in [31].

### Acknowledgements

This work has been (partially) supported by the Project EFI ANR-17-CE40-0030 of the French National Research Agency.

The authors want to acknowledge an anonymous reviewer for an accurate reading and constructive comments.

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