

In honour of the 19th :-) birthday of  
Patrick and Christian

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# *Nonparametric estimation for stochastic damping Hamiltonian systems under partial observation*

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# Outline of the talk

- 1 Introduction
- 2 Assumptions, probabilistic properties of the model
- 3 Adaptive estimation
- 4 Numerical results

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# Model

We consider a damped hamiltonian system with stochastic noise:  
( $Z_t := (X_t, Y_t) \in \mathbb{R}^{2d}$ ,  $t \geq 0$ ) given by the s.d.e.

$$\begin{cases} dX_t = Y_t dt \\ dY_t = \sigma(X_t, Y_t) dB_t - (c(X_t, Y_t) Y_t + \nabla V(X_t)) dt \end{cases}$$

where  $B$  is a standard Brownian motion.

# Particle in contact with a heat bath

## Particle in a contact with a heat thermal reservoir:

Consider a particle in a potential  $V$ . Its dynamics is described through the Hamiltonian  $H(p, q) = \frac{1}{2} p^2 + V(q)$ :  $\dot{q} = \partial_p H$ ,  $\dot{p} = -\partial_q H$ .

This particle is in contact with some heat bath with temperature  $T > 0$ , modeled by an Ornstein-Uhlenbeck process acting as a noise on the momentum  $p$  only:

$$\begin{aligned}dq_t &= p_t dt \\ dp_t &= (-\gamma p_t - \nabla V(q_t))dt + \sqrt{2\gamma T} dB_t.\end{aligned}$$

# Chain of oscillators in contact with two heat baths (1/2)

## Chain of oscillators in contact with two heat bath reservoirs:

The dynamics is given by the Hamiltonian

$$H(p, q) = \sum_{1 \leq i \leq d} \frac{p_i^2}{2} + V(q)$$

with  $V : \mathbb{R} \rightarrow \mathbb{R}$  in the form

$$\sum_{1 \leq i \leq d} U^{(1)}(q_i) + \sum_{1 \leq i \leq d-1} U^{(2)}(q_i - q_{i+1}).$$

$U^{(1)}$  is a trapping potential,  $U^{(2)}$  an interaction potential.



# Chain of oscillators in contact with two heat baths (2/2)

The two heat baths with temperature  $T_1$  and  $T_d$  act on the momenta of particles 1 and  $d$  like Ornstein-Uhlenbeck processes:

$$dq_j(t) = p_j(t)dt \quad 1 \leq j \leq d$$

$$dp_1(t) = (-\gamma p_1(t) - \partial_{q_1} V(q_t))dt + \sqrt{2\gamma T_1} dB_1(t)$$

$$dp_j(t) = (-\partial_{q_j} V(q_t))dt \quad 2 \leq j \leq d-1$$

$$dp_d(t) = (-\gamma p_d(t) - \partial_{q_d} V(q_t))dt + \sqrt{2\gamma T_d} dB_d(t).$$

For short:

$$\begin{cases} dq(t) = p(t)dt \\ dp(t) = (-\gamma \Lambda p(t) - \nabla_q V(q(t)))dt + \sqrt{2\gamma T} dB(t) \end{cases}$$

where  $\Lambda : \mathbb{R}^d \rightarrow \mathbb{R}^2$  is the projection  $\Lambda(x_1, \dots, x_d) = (x_1, x_d)$ ,  
 $\sqrt{T} : (x_1, x_d) \rightarrow (\sqrt{T_1}x_1, \sqrt{T_d}x_d)$  and  $B(t) = (B_1(t), B_d(t))$  is a B.M.  
 in  $\mathbb{R}^2$ .

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 in  $\mathbb{R}^2$ .

# Langevin dynamics

More generally, consider the following system:

$$\begin{cases} dX_t = Y_t dt \\ dY_t = (2\beta^{-1})^{1/2} \sigma(X_t) dW_t - (\sigma^2(X_t) Y_t + \nabla V(X_t)) dt. \end{cases}$$

Observe that the diffusion term,  $\tilde{\sigma}(x) = (2\beta^{-1})^{1/2} \sigma(x)$ , depends only on the  $x$  coordinate and on an unknown parameter  $\beta$ . Moreover, the damping force has the form  $c(x, y) = \frac{\beta}{2} \tilde{\sigma}^2(x)$ .

Since Einstein, this last relationship between the damping force and the diffusion term is known as the **fluctuation-dissipation relation** and has numerous applications, e.g.,

- it appears as a tool for the simulation of molecular dynamics (see, e.g., Lelièvre *et al.* (2010) [Section 2.2.3] ),
- it also appears as limit of the Ehrenfest nuclei dynamics and is called Langevin dynamics (see Szepeszy (2011) ).

# A stochastic neuronal model (1/3)

See, e.g., León and Samson (2017) and references therein.

## Data

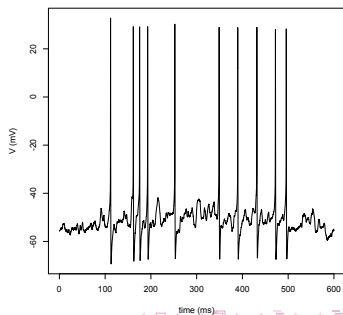
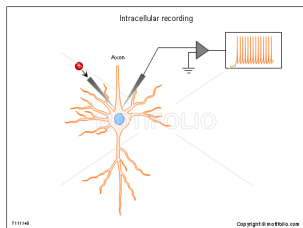
- Membrane potential: difference in voltage between the interior and exterior of the cell
- High frequency records available ( $\delta = 0.1$  ms)

## Objective

- Prediction of spike emission

## Tools

- Neuronal modeling with stochastic models
- Estimation

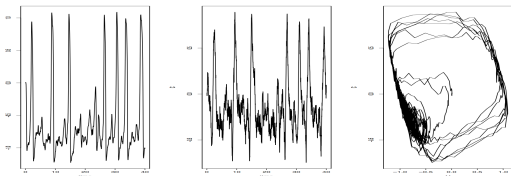


# A stochastic neuronal model (2/3)

The **stochastic FitzHugh-Nagumo model** is defined as follows:

$$\begin{cases} dV_t = \frac{1}{\varepsilon}(V_t - V_t^3 - C_t - s)dt \\ dC_t = (\gamma V_t - C_t + \beta)dt + \tilde{\sigma}dB_t \end{cases}$$

- $V_t$  the membrane potential of a single neuron,
- $C_t$  a recovery variable / channel kinetics,
- $\varepsilon$  the time scale separation,  $s$  the stimulus input,
- $\beta, \gamma$  positive constants determining the position of the fixed point and the duration of the excitation,
- $B_t$  a Brownian motion,  $\tilde{\sigma} > 0$  the diffusion coefficient.



## A stochastic neuronal model (3/3)

Defining  $X_t = V_t$  and  $Y_t = \frac{1}{\varepsilon}(V_t - V_t^3 - C_t - s)$ , we get:

$$\begin{cases} dX_t = Y_t dt \\ dY_t = \frac{1}{\varepsilon}(Y_t(1 - \varepsilon - 3X_t^2) - X_t(\gamma - 1) - X_t^3 - (s + \beta))dt + \frac{\tilde{\sigma}}{\varepsilon}dB_t. \end{cases}$$

Let  $\nabla V(x) = \frac{1}{\varepsilon}(x(\gamma - 1) + x^3 + (s + \beta))$  and  $c(x) = \frac{1}{\varepsilon}(-1 + \varepsilon + 3x^2)$ .

We recognize the system governing the dynamics of a particle with  $X_t$  referring to its position and  $Y_t$  to its velocity, whose movement is guided by a potential  $V(x)$  and a damping force  $c(x)$ .

$$\begin{cases} dX_t = Y_t dt \\ dY_t = -(c(X_t)Y_t + \nabla V(X_t))dt + \sigma dB_t \end{cases}$$

with  $\sigma = \frac{\tilde{\sigma}}{\varepsilon}$ .

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# Assumptions

$Z_t = (X_t, Y_t) \in \mathbb{R}^{2d}$  governed by:

$$\begin{cases} dX_t = Y_t dt \\ dY_t = -(c(X_t, Y_t)Y_t + \nabla V(X_t))dt + \sigma dB_t, \end{cases}$$

with  $\sigma > 0$ .

$\mathcal{H}_1$  The potential  $V(x)$  is smooth over  $\mathbb{R}^d$  and lower bounded.

$\mathcal{H}_2$  The damping force  $c(x, y)$  is smooth, bounded, and there exist  $c, M > 0$  s.t.  $c^s(x, y) \geq cld > 0, \forall (|x| > M, y \in \mathbb{R}^d)$ .

From **Wu (2001)**, we know that for every initial state  $z = (x, y) \in \mathbb{R}^{2d}$ , the system admits a unique weak solution, and that this solution is non-explosive.

The infinitesimal generator writes:

$$L = \frac{\sigma^2}{2} \partial_{yy} + y \partial_x - (c(x, y)y + \nabla_x V(x)) \partial_y.$$



# Local properties, hypoellipticity (1/4)

$L$  can be written in Hörmander's form

$$L = \frac{\sigma^2}{2} \sum_{i=1}^d L_i^2 + L_0$$

with vector fields  $L_i$  defined by

(1) pour  $1 \leq i \leq d$ ,  $L_i = \frac{\partial}{\partial y_i}$ ,

(2)

$$L_0 = \sum_{k=1}^d y_k \frac{\partial}{\partial x_k} - \sum_{k=1}^d \left( (c(x, y)y)_k + \frac{\partial V}{\partial x_k} \right) \frac{\partial}{\partial y_k}.$$

It holds

$$[L_i, L_0] = L_i L_0 - L_0 L_i = \frac{\partial}{\partial x_i} - \sum_{k=1}^d \frac{\partial((c(x, y)y)_k)}{\partial y_i} \frac{\partial}{\partial y_k}$$

so that  $\{L_i, 1 \leq i \leq d; [L_i, L_0], 1 \leq i \leq d\}(z)$  span  $\mathbb{R}^{2d}$ , for all  $z$ .

## Local properties, hypoellipticity (2/4)

$\Rightarrow$  hypoellipticity (Hörmander sum of squares theorem).

**Consequence:**  $\forall z, \forall t > 0$ ,

the distribution  $P_t(z, \cdot)$  of  $Z_t$  starting at  $z$  at time 0 admits a  $C^\infty$  density  $p_t(z, \cdot)$  w.r.t. Lebesgue.

Hence,  $\mu(dz) = p_s(z)dz$  with  $p_s C^\infty$ , and thus the strong Feller property.

Small time behavior of  $p_t(z, \cdot)$  ?

**Example:**  $d = 1, c = V = 0$ . Then  $Z_t$  is a two dimensional gaussian vector, with mean  $(x_0 + y_0 t, y_0)$  and covariance matrix

$$\text{Var}(X_t) = \frac{t^3}{3}, \text{Var}(Y_t) = t, \text{Cov}(X_t, Y_t) = \frac{t^2}{2}.$$

So the transition density behaves, for small  $t$ , as

$$\frac{\sqrt{3}}{\pi} \frac{1}{t^2} e^{-\frac{y_0^2}{6t}} \quad \text{instead of} \quad \frac{1}{2\pi} \frac{1}{t}$$

which is the classical small time explosion for elliptic diffusions (like the B.M.).

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# Local properties, hypoellipticity (3/4)

Theorem (Konakov, Menozzi & Molchanov, 2010)

$$\begin{cases} dX_t = Y_t dt \\ dY_t = \sigma dW_t + b(X_t, Y_t) dt, \end{cases}$$

with  $b \in C^\infty$ , bounded as well as all its derivatives. Let  $T > 0$ . Then  $\forall z = (x, y)$ ,  $\forall t > 0$ , the distribution of  $Z_t = (X_t, Y_t)$  has a density  $q_t(z, \cdot)$  with respect to Lebesgue and  $\exists C, C' > 0$  t.q. for  $0 < t < T$ ,

$$q_t(z, z') \leq C' \frac{1}{t^{2d}} \exp \left( -C \left[ \frac{|y - y'|^2}{4t} + \frac{3 \left| x' - x - \frac{t(y+y')}{2} \right|^2}{t^3} \right] \right).$$

De plus,  $\exists t_0 > 0, \exists C'' > 0$  t.q.  $\forall 0 < t < t_0$ ,

$$q_t((x, y), (x + ty, y)) \geq C'' \frac{1}{t^{2d}}.$$

# Local properties, hypoellipticity (4/4)

We can generalize that result to a non bounded drift term.

Corollary (Cattiaux, León & Prieur, 2014)

*We do no more assume boundedness.*

*$\forall z$ , for any open neighborhood  $U$  of  $z$ , one can write:*

$$\forall z' \in U, \forall 0 < t < T, p_t(z, z') \leq q_t(z, z') + C(U)e^{-\frac{C'(U)}{t}}$$

*for constants  $C(U)$  and  $C'(U) > 0$ .*

*We also prove*

$$\forall (z, z'), \exists 0 < C(z') \text{ s.t. } \forall t \geq 0, p_t(z, z') \leq D(z') < +\infty.$$

# Long time behavior, coercivity and mixing (1/2)

We now add the following assumption

$\mathcal{H}_3$   $V$  and  $\nabla V$  have polynomial growth at infinity with

$$+\infty \geq \liminf_{|x| \rightarrow +\infty} \frac{x \cdot \nabla V(x)}{|x|} \geq \nu > 0 \quad (\text{drift's condition}).$$

The force  $-\nabla V(x)$  is "strong enough" for  $|x|$  large to ensure a quick return of the system to compact subsets of  $\mathbb{R}^{2d}$ .

*Under  $\mathcal{H}_i$ ,  $i = 1, 2, 3$ , the process  $Z_t = (X_t, Y_t)$  is positive recurrent with a unique invariant probability measure  $\mu$ . Moreover, moments of any order of  $\mu$  exist: for all  $k_1, k_2 \in \mathbb{N}$ ,*

$$\mathbb{E}(X_t^{k_1} Y_t^{k_2}) = \int x^{k_1} y^{k_2} d\mu(x, y) < +\infty.$$

**Scheme of proof:** the proof involves the construction of a Lyapunov function  $\Psi(x, y)$ , such that there exist a compact  $K \in \mathbb{R}^{2d}$  and constants  $C, \xi > 0$ , such that  $-\frac{L\Psi}{\Psi} \geq \xi \mathbb{1}_{K^c} - C \mathbb{1}_K$ . The choice of the Lyapunov function is not trivial. See, e.g., Wu (2001).

## Long time behavior, coercivity and mixing (2/2)

For any  $z$ , let's  $P_t f(z) = \mathbb{E}_z(f(Z_t))$  for bounded  $f$ 's.

$\psi \in \mathbb{L}^1(\mu)$ . There exist  $D > 0$  and  $\rho < 1$  s.t. for all  $z$ , all  $f$  s.t.  
 $\sup_z \frac{|f(z)|}{\psi(z)} < +\infty$ ,

$$\left| P_t f(z) - \int f d\mu \right| \leq D \sup_a \left( \frac{|f(a) - \int f d\mu|}{\psi(a)} \right) \psi(z) \rho^t.$$

It follows that  $(Z_t := (X_t, Y_t), t \geq 0)$  is  $\beta$ -mixing.

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# Invariant density estimators

**Complete observations:** we observe both coordinates  $X_t$  and  $Y_t$  at discrete times  $i\delta$ ,  $i = 1, \dots, n$ . Let  $K$  be a kernel function,  $b = (b_1, b_2)$  a bandwidth.

$$\check{p}_s(x, y) := \frac{1}{nb_1^d b_2^d} \sum_{i=1}^n K\left(\frac{x - X_{i\delta}}{b_1}, \frac{y - Y_{i\delta}}{b_2}\right).$$

Now we do not observe  $y$  anymore.

**Partial observations:**

$$\hat{p}_s(x, y) := \frac{1}{nb_1^d b_2^d} \sum_{i=1}^{n-1} K\left(\frac{x - X_{i\delta}}{b_1}, \frac{X_{(i+1)\delta} - X_{i\delta}}{\delta}\right).$$

**Main issue:** the choice of the bandwidth  $b = (b_1, b_2)$ .

# Invariant density estimators

**Complete observations:** we observe both coordinates  $X_t$  and  $Y_t$  at discrete times  $i\delta$ ,  $i = 1, \dots, n$ . Let  $K$  be a kernel function,  $b = (b_1, b_2)$  a bandwidth.

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**Main issue:** the choice of the bandwidth  $b = (b_1, b_2)$ .

# Adaptive estimation (1/2)

**Data-driven procedure** [Comte, Prieur, Samson, 2017]

Our selection criterion is based on Goldenshluger and Lepski (2011).

Let  $\check{p}_{b,b'} = K_{b'} \star \check{p}_b(x, y)$ , with  $K_{b'}(u, v) = \frac{1}{b'_1 b'_2} K(\frac{u}{b'_1}, \frac{v}{b'_2})$ . Let  $p_b = K_b \star p$ . In the following,  $d = 1$ .

$$\tilde{b} = \arg \min_{b \in \mathcal{B}_n} (A(b) + U(b)), \quad \text{with}$$

- $\mathcal{B}_n = \{(b_{1,k}, b_{2,\ell}) = (1/k, 1/\ell), k, \ell = 1, \dots, B_n\}$ ,
- $A(b)$  mimicking the bias ( $= \sup_{b' \in \mathcal{B}_n} (\|\check{p}_{b,b'} - \check{p}_{b'}\|^2 - U(b'))_+$ )
- $V(b)$  mimicking the variance ( $= \kappa \frac{\|K\|_1^2 \|K\|^2}{n b_1 b_2} \sum_{i=0}^{n-1} \beta(i\delta)$ )

$$\mathbb{E} (\|\check{p}_{\tilde{b}} - p\|^2) \leq C \inf_{b \in \mathcal{B}_n} (\|p - p_b\|^2 + U(b)) + C' \frac{\log(n)}{n\delta}.$$

$\kappa$  can then be calibrated by the slope heuristic (see Arlot and Massart, 2009, Lacour *et al.*, 2016). Same results in the partial observations case.

## Adaptive estimation (2/2)

That procedure is numerically demanding due to the double convolutions  $\check{p}_{b,b'}$ , especially in the multidimensional case.

In practice, we implement the selection procedure in Lacour, Massart and Rivoirard (2016):

$$\hat{b} = \arg \min_{b \in \mathcal{B}_n} (\|\check{p}_b - \check{p}_{b_{min}}\|^2 + U(b)) \text{ with}$$

$$b_{min} = (\min_{1 \leq k \leq B_n} b_{1,k}, \min_{1 \leq \ell \leq B_n} b_{2,\ell}).$$

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# Numerical results (1/4)

## Harmonic Oscillator:

$$\begin{cases} dX_t = Y_t dt \\ dY_t = -(\alpha X_t + \gamma Y_t) dt + \sigma dB_t \end{cases}$$

with  $\alpha > 0, \gamma > 0$ . In the following, we choose  $\alpha = 4, \gamma = 0.5, \sigma = 0.5$ . The potential is then  $V(x) = \alpha/2x^2$ . The stationary distribution is Gaussian, with mean zero and explicit diagonal variance matrix:

$$p(x, y) = \frac{\gamma\sqrt{\alpha}}{\pi\sigma^2} \exp\left(-\frac{2\gamma}{2\sigma^2}y^2 - \frac{2\gamma\alpha}{2\sigma^2}x^2\right)$$

with diagonal variances equal to 1/16 and 1/4, respectively in our case.

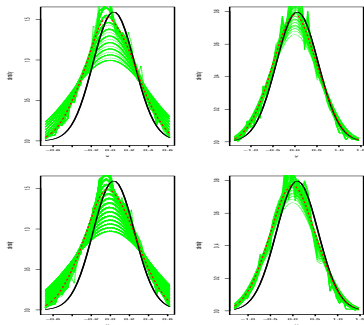
# Numerical results (2/4)

Kernel estimation of the invariant density:

- complete observations (top)
- partial observations (bottom)

$n = 2000, \delta = 0.2$ .

100 trajectories simulated with a Euler scheme with step size  $\delta/10$ .



$\mathcal{B}_n = \{(b_1, b_2) \in \{1/\sqrt{4n}, 2/\sqrt{4n}, \dots, 30/\sqrt{4n}\}^2\}$ . Anisotropic selected bandwidth  $\hat{b} = (8/\sqrt{4n}, 17/\sqrt{4n})$  (complete),  $\hat{b} = (9/\sqrt{4n}, 19/\sqrt{4n})$  (partial).

# Numerical results (3/4)

## Van Der Pol Oscillator:

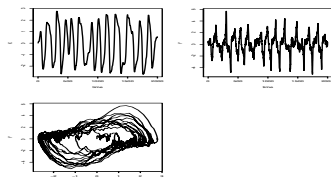
$$\begin{cases} dX_t = Y_t dt \\ dY_t = -((c_1 X_t^2 - c_2)Y_t + \omega_0^2 X_t) dt + \sigma dB_t \end{cases}$$

with  $\sigma, c_1, c_2, \omega_0^2 > 0$ . In the following, we choose  $\sigma = c_1 = c_2 = \omega_0^2 = 1$ . The potential is then  $V(x) = \omega_0^2/2x^2$ . The invariant density  $p$  satisfies Fokker-Planck equation:

$$\frac{1}{2} \frac{\partial^2 p(x, y)}{\partial y^2} - y \frac{\partial p(x, y)}{\partial x} + c(x)p(x, y) + (c(x)y + \nabla D(x)) \frac{\partial p(x, y)}{\partial y} = 0$$

solved with finite difference scheme (see Kumar *et al.*, 2006).

Sample  $(X_{i\delta})_{i=0, \dots, n}$  (top left),  
 $(Y_{i\delta})_{i=0, \dots, n}$  (top right) and state  
 phase (bottom) for  $\delta = 0.5$  and  
 $n = 2000$ .





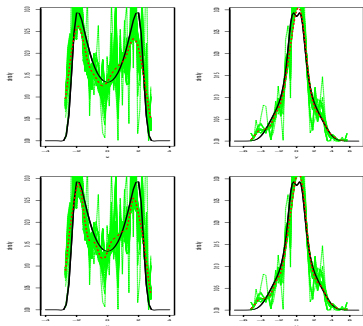
# Numerical results (4/4)

Kernel estimation of the invariant density:

- complete observations (top)
- partial observations (bottom)

$n = 2000, \delta = 0.05$ .

100 trajectories simulated with a Euler scheme with step size  $\delta/10$ .



# Conclusion, perspectives

**Conclusion:** we obtained

- non parametric (recursive) estimation for the invariant density (Cattiaux *et al.*, 2014a, 2015 ),
- a data-driven procedure for the selection of the bandwidth (see Comte *et al.*, 2016 ),
- see also Cattiaux *et al.* (2014b,2016,2017) for the estimation of the drift and of the diffusion matrix.

We have considered the more realistic non trivial case of partial observations.

**Perspectives:**

- to consider more complex models which are more realistic for environmental modeling (non linear Fokker-Planck equations, confined models, degenerated variances, ... ),
- adaptive estimation in higher dimension,
- adaptivity with respect to  $\delta$ ,
- ...

## Some references I



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Thanks for your attention

Happy birthday  
Parabéns !

