On Markov intertwinings

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Top-to-random shuffle [Aldous and Diaconis, 1986]

Pitman's theorem [1975, Pitman and Rogers 1981]

One-dimensional diffusions

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Decks of cards

N cards labeled 1, 2, ..., *N*. A deck of these cards is represented by an element σ of the symmetric group S_N :



Figure: The deck of cards σ

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Top-to-random shuffle

Starting from the ordered deck (1, 2, ..., N), shuffle it as follows, where the positions I_1 , I_2 , I_3 , ... are independent and uniformly distributed on $\{1, 2, ..., N\}$:



Figure: An example of top-to-random shuffle, here $I_3 = N$

For any $n \in \mathbb{Z}_+$, let $X_n \in S_N$ be the deck at time n. The random sequence $X := (X_n)_{n \in \mathbb{Z}_+}$ is a Markov chain: to construct X_{n+1} , only the knowledge of X_n (and of I_{n+1} , which is independent from the past) is needed, not the way we ended up with X_n .

The uniform distribution \mathcal{U}_{S_N} on \mathcal{S}_N is invariant for X: if $X_0 \sim \mathcal{U}_{S_N}$, then for any time $n \in \mathbb{Z}_+$, $X_n \sim \mathcal{U}_{S_N}$. Furthermore, by irreducibility and aperiodicity, we know that the law $\mathcal{L}(X_n)$ converges to \mathcal{U}_{S_N} for large time n.

But from a practical point of view, we want to know for how long we have to shuffle the deck to be close to \mathcal{U}_{S_N} . Quantitative convergence provides such estimates.

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Different ways to measure closeness of probability measures m and μ on the same state space:

Total variation:

$$\left\|m-\mu\right\|_{\mathrm{tv}}\coloneqq \sup_{A \; \mathrm{event}} \left|m(A)-\mu(A)
ight| = rac{1}{2} \left\|rac{dm}{d\mu}-1
ight\|_{\mathbb{L}^{1}(\mu)}$$

Separation:

$$\mathfrak{s}(m,\mu) \coloneqq \mathrm{esssup}_{\mu} 1 - \frac{dm}{d\mu} \geq \|m-\mu\|_{\mathrm{tv}}$$

For top-to-random shuffle: for any c > 0,

$$\left\|\mathcal{L}(X_{\lfloor N\ln(N)+cN
floor})-\mathcal{U}_{\mathcal{S}_N}
ight\|_{\mathrm{tv}}\leq \exp(-c)$$

i.e. order $N \ln(N)$ shuffles are needed to be almost at equilibrium. Furthermore, it can be shown to be a cut-off time.

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The proof of Aldous and Diaconis [1986] used strong stationary times associated to a Markov chain X: they are finite stopping times (only the past of X and independent randomness are needed to decide to stop) τ such that

$$au \perp X_{ au}$$
 and $X_{ au} \sim \mu$

where μ is the invariant probability of X. They provide exact simulations of μ and we have

$$\forall n \in \mathbb{Z}_+, \quad \mathfrak{s}(\mathcal{L}(X_n), \mu) \leq \mathbb{P}[\tau > n]$$

In the top-to-random shuffle, a strong stationary time is given by the first time the last card reaches the top and is inserted in the deck. It is distributed as a sum of independent geometric random variables of parameters 1/N, 2/N, ..., (N-1)/N and 1. It leads to the previous bound (coupon collector problem).

Set-valued dual

How to construct a strong stationary time? Consider $\mathcal{D} := \{A \subset S_N\} \setminus \{\emptyset\}$ and for any $\sigma \in S_N$ and $k \in \{1, 2, ..., N\}$

$$egin{aligned} \mathcal{A}_{\sigma,k}\coloneqq \{\sigma'\in\mathcal{S}_{\mathcal{N}}\,:\,\sigma'(1)=\sigma(1),\sigma'(2)=\sigma(2),...,\sigma'(k)=\sigma(k)\} \end{aligned}$$

Define a set-valued process $D \coloneqq (D_n)_{n \in \mathbb{Z}_+}$ via

$$D_n := A_{X_n, \text{position of the initial last card}}$$

as long as the last card has not been reinserted, otherwise take $D_n := S_N$. Then D is a D-valued Markov chain absorbed at S_N .

Define the Markov kernel Λ from \mathcal{D} to \mathcal{S}_N via

$$\forall A \in \mathcal{D}, \forall \sigma \in \mathcal{S}_{N}, \qquad \Lambda(A, \sigma) \coloneqq \frac{\mathcal{U}_{\mathcal{S}_{N}}(\sigma)}{\mathcal{U}_{\mathcal{S}_{N}}(A)} \mathbb{1}_{\sigma \in A}$$

(=conditional expectation on subsets under the invariant probability).

Let P be the transition kernel associated to the time-homogeneous Markov chain X:

$$\forall \sigma, \sigma' \in \mathcal{S}_{N}, \quad P(\sigma, \sigma') := \mathbb{P}[X_{n+1} = \sigma' | X_n = \sigma]$$

The r.h.s. is 1/N if $X_{n+1} \circ X_n^{-1}$ is a cycle of the form $(1 \rightarrow i \rightarrow i - 1 \rightarrow i - 2 \rightarrow \cdots \rightarrow 2 \rightarrow 1)$ and 0 otherwise. Let Q be the transition kernel associated to the time-homogeneous Markov chain D. We have the intertwining relation

$$Q\Lambda = \Lambda P$$
 (1)

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which makes sense, since Q is a $\mathcal{D} \times \mathcal{D}$ matrix, Λ is a $\mathcal{D} \times S_N$ matrix and P is a $S_N \times S_N$ matrix.

Fill and Diaconis [1990] have shown that (1) is a key for the construction of strong stationary times for finite Markov processes, as soon as the transition kernel P is ergodic and Q is absorbing.

Assume (1) is satisfied for general such P and Q. Let $X := (X_n)_{n \in \mathbb{Z}_+}$ and $D := (D_n)_{n \in \mathbb{Z}_+}$ be corresponding Markov chains, with $\mathcal{L}(D_0)\Lambda = \mathcal{L}(X_0)$. Then there exists a coupling of X and D such that for all times $n \in \mathbb{Z}_+$, we have for the conditional expectations:

$$\begin{cases} \mathcal{L}(X_n|D_{\llbracket 0,n \rrbracket}) = \Lambda(D_n, \cdot) \\ \mathcal{L}(D_{\llbracket 0,n \rrbracket}|X) = \mathcal{L}(D_{\llbracket 0,n \rrbracket}|X_{\llbracket 0,n \rrbracket}) \end{cases}$$
(2)

It follows that the absorption time of D is a strong stationary time for X. Our goal is to extend such properties to diffusion processes X.

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Brownian motion

The Brownian motion is approximated (in law) by letting $\epsilon \to 0_+$ in the following picture, where each step $\pm \sqrt{\epsilon}$ is chosen with probability 1/2:



Figure: Approximation of a Brownian path

Pitman's transformation

Let $(B_t)_{t\geq 0}$ be a Brownian motion. Consider the Pitman's transformation $R := (R_t)_{t\geq 0}$ given by

$$R_t := -B_t + 2 \max_{s \in [0,t]} B_s$$

Let Λ the Markov kernel from \mathbb{R}_+ to \mathbb{R} given by

$$\forall r \geq 0, \qquad \Lambda(r, dx) \coloneqq \mathcal{U}_{[-r,r]}(dx)$$

Relation (2) can be extended to this continuous setting: for all time $t \ge 0$,

$$\begin{cases} \mathcal{L}(B_t|R_{[0,t]}) = \Lambda(R_t, \cdot) \\ \mathcal{L}(R_{[0,t]}|B) = \mathcal{L}(R_{[0,t]}|B_{[0,t]}) \end{cases}$$
(3)

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Furthermore, the process R is a Bessel-3 process, i.e. it has the law of the norm of a Brownian motion in dimension 3.

Pitman's theorem in picture



Figure: Trajectories: Brownian motion $B_{[0,t]}$, $R_{[0,t]}$, $-R_{[0,t]}$, and the segment-valued dual: $[-R_t, R_t]$

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Generators

The intertwining relation (1) holds for the generators G and G of B and R:

$$G\Lambda = \Lambda G$$

Here
$$G = \frac{1}{2}\partial_x^2$$
 (on $\mathcal{C}^2(\mathbb{R})$) and $\mathcal{G} = \frac{1}{2}\partial_r^2 + \frac{1}{r}\partial_r$ (on $\mathcal{C}^2(\mathbb{R}_+)$).

The generator G of a Markov process $(X_t)_{t\geq 0}$ is understood in the sense of martingale problems: for any function f from the domain of G (=nice observable on the underlying state space), the process $M^f := (M_t^f)_{t\geq 0}$ is a (local) martingale, where

$$M_t^f := f(X_t) - f(X_0) - \int_0^t G[f](X_s) \, ds$$

Martingales are among the main tools of probability theory, they satisfy

$$\forall t,s \geq 0, \quad \mathbb{E}[M_{t+s}^f|X_{[0,t]}] = M_t^f$$

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On the circle

Consequently of (3), we get estimates on the convergence of the Brownian motion $W := (W_t)_{t\geq 0}$ on the circle $\mathbb{T} := \mathbb{R}/(2\pi\mathbb{Z})$: let τ be the hitting time of π by R:

$$\tau := \inf\{t \ge 0 : R_t = \pi\}$$

It is a strong stationary time for W. The hitting times of Bessel processes are well-studied, we have:

$$\forall \ \lambda > 0, \qquad \mathbb{E}[\exp(-\lambda\tau)] = \frac{\sqrt{\pi}\sqrt[4]{\lambda}}{\sqrt[4]{2}\Gamma(3/2)} \frac{1}{I_{1/2}(\pi\sqrt{2\lambda})}$$

where $I_{1/2}$ is the modified Bessel function of index 1/2. A Tauberien theorem enables to deduce the behavior for large $t \ge 0$ of $\mathbb{P}[\tau > t]$ and thus of $\mathfrak{s}(\mathcal{L}(W_t), \mathcal{U}_{\mathbb{T}})$ and $\|\mathcal{L}(W_t) - \mathcal{U}_{\mathbb{T}}\|_{\mathrm{tv}}$. Top-to-random shuffle [Aldous and Diaconis, 1986]

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Elliptic real diffusions

On \mathbb{R} , consider a diffusion X whose generator is $G \coloneqq a\partial^2 + b\partial$, where a, b are smooth functions and a is positive. Consider the speed measure $\mu(dx) \coloneqq \underline{\mu}(x) dx$ with density

$$\underline{\mu}(x) \coloneqq \frac{\exp(c(x))}{a(x)} \quad \text{with} \quad c(x) \coloneqq \int_0^x \frac{b(y)}{a(y)} \, dy$$

It is invariant for X and even reversible: G can be extended into a self-adjoint operator on $\mathbb{L}^2(\mu)$. Consider

$$\mathcal{D} := \{[y, z] : y, z \in (-\infty, +\infty), y \leq z\}$$

As usual, define the associated Markov kernel Λ from $\mathcal D$ to $\mathbb R$ via

$$\forall [y, z] \in \mathcal{D}, \Lambda([y, z], dx) \coloneqq \begin{cases} \delta_y(dx) &, \text{ if } y = z \\ \frac{\mu(x)}{\mu([y, z])} \mathbb{1}_{[y, z]}(x) dx &, \text{ otherwise} \end{cases}$$

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Dual process

On $\ensuremath{\mathcal{D}}$ define the degenerate generator

$$\begin{aligned} \mathcal{G} &\coloneqq (\sqrt{a(z)}\partial_z - \sqrt{a(y)}\partial_y)^2 \\ &+ (a'(y)/2 - b(y))\partial_y + (a'(z)/2 - b(z))\partial_z \\ &+ 2\frac{\sqrt{a(y)}\mu(y) + \sqrt{a(z)}\mu(z)}{\mu([y,z])}(\sqrt{a(z)}\partial_z - \sqrt{a(y)}\partial_y), \end{aligned}$$

Its interest is in the intertwining relation $\mathcal{G}\Lambda = \Lambda G$.

Proposition

There exists a unique process [Y, Z] whose generator is \mathcal{G} , up to its explosion time ζ . The diagonal is an entrance boundary for this process.

The proof is based on the fact that $(\mu([Y_t, Z_t]))_{t \in [0, \zeta]}$ is a (stopped) Bessel-3 process, up to the time-change $(\theta_t)_{t \in [0, \varsigma]}$ given by

$$2\int_{0}^{\theta_{t}} (\sqrt{a(Y_{s})}\underline{\mu}(Y_{s}) + \sqrt{a(Z_{s})}\underline{\mu}(Z_{s}))^{2} ds = t$$

Strong stationary times in 1-dimension

Application to the existence of strong stationary times:

Theorem

Assume that X is positive recurrent. There exists a strong stationary time for X, whatever its initial distribution, if and only if $-\infty$ and $+\infty$ are entrance boundaries.

Positive recurrence or ergodicity: μ is finite and in large time X_t converges to the renormalization of μ , analytically this is characterized by

$$\int_{-\infty}^{0} \exp(-c(y)) \, dy = +\infty \quad \text{and} \quad \int_{0}^{\infty} \exp(-c(y)) \, dy = +\infty$$

Entrance boundary: e.g. for $+\infty$, it means that X can be started from $+\infty$ (comes down from infinity), it amounts to:

$$\int_0^{+\infty} \left(\int_0^x \exp\left(-c(y)\right) \, dy \right) \, \mu(dx) \, < \, +\infty$$

SQC

Consider again the generator $G := a\partial^2 + b\partial$, on \mathbb{R} or \mathbb{T} , where *a* is allowed to vanish on a finite number of points, but such that \sqrt{a} remains smooth. On the vanishing points, assume *b* does not vanish: 1-dimensional hypoellipticity in the sense of Hörmander [1967].

Up to some adjustments (modification of the measure used in Λ , dual processes which can disconnect, ...), the above approach is valid and enables to recover the density theorem (i.e. the law of X_t admits a density for all t > 0) and to obtain estimates on the speed of convergence to equilibrium (even when the invariant measure does not charge the whole state space). The Bessel-3 process is still there, the hypoellipticity is only felt at the level of the time change.

Is it possible to recover the generality of Hörmander's theorem in this probabilistic way?

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Consider a (complete) Riemannian manifold V of dimension $n \ge 2$. Examples in dimension n = 2:



The Laplacian *G* associated to *V* is the generator of the Brownian motion $X := (X_t)_{t \ge 0}$ on *V* (up to speeding of time by 2). The process *X* can be approximated as in \mathbb{R} by using independent and uniform increments on spheres of radius $\sqrt{\epsilon}$, as in the real case. The generator *G* is reversible with respect to the Riemannian measure μ . When *V* is compact, μ can be normalized into a probability measure.

Let \mathcal{D} be the set of compact subdomains of V with a smooth boundary. Consider Λ the Markov kernel from \mathcal{D} to V, corresponding to the conditioning of μ . We want:

• to find a Markov generator G intertwined with G through Λ :

$$\mathcal{G}\Lambda = \Lambda G$$
 (4)

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• To associate to \mathcal{G} a Markov evolution of domains $(\mathfrak{D}_t)_{t\geq 0}$, starting from a singleton: $\mathfrak{D}_0 = \{x_0\}$.

• To couple the evolutions X and $(\mathfrak{D}_t)_{t\geq 0}$ so that

$$\left\{ \begin{array}{rcl} \mathcal{L}(X_t | \mathfrak{D}_{[0,t]}) & = & \Lambda(\mathfrak{D}_t, \cdot) \\ \mathcal{L}(\mathfrak{D}_{[0,t]} | X) & = & \mathcal{L}(\mathfrak{D}_{[0,t]} | X_{[0,t]}) \end{array} \right.$$

• To recover density theorems and to construct strong stationary times when V is compact.

Mean curvature flow

Let be given $D_0 \in \mathcal{D}$. To any $y \in \partial D_0$, associate the exit unitary normal vector ν_y and the mean curvature $\rho(y)$. At least for small time $t \ge 0$, it is possible to make the domain evolve according to

$$\forall y_t \in \partial D_t, \qquad \dot{y}_t = -\rho(y_t)\nu_{y_t}$$

The domains D_t have a tendency to round up and to shrink to a point in finite time.





Stochastic modification of the mean curvature flow

Modify the previous deterministic evolution into an infinite-dimensional stochastic differential equation on $(\mathfrak{D}_t)_{t \in [0,\zeta]}$:

$$\forall Y_t \in \partial \mathfrak{D}_t, \ dY_t = \left(\sqrt{2}dB_t + \left(2\frac{\underline{\mu}(\partial \mathfrak{D}_t)}{\mu(\mathfrak{D}_t)} - \rho(Y_t)\right)dt\right)\nu_{Y_t}$$
(5)

where $(B_t)_{t\geq 0}$ is a one-dimensional Brownian motion and $\underline{\mu}$ is (n-1)-dimensional Hausdorff measure. The global isoperimetric ratio $\underline{\mu}(\partial \mathfrak{D}_t)/\mu(\mathfrak{D}_t)$ counters the effect of the mean curvature and prevents the evolution to collapse to a singleton.

Theorem

Starting from a non-singleton element of \mathcal{D} , it is possible to define $(\mathfrak{D}_t)_{t\in[0,\zeta]}$, where ζ is a positive random time, solving (5) and whose generator \mathcal{G} satisfies the intertwining relation (4).

The proof is based on an extension of the Doss-Sussman method to the infinite dimensional setting of \mathcal{D} .

When V has constant curvature, (5) can be solved for all times, starting from a singleton $\{x_0\}$. In this situation D_t is a ball centered at x_0 and of radius R_t solving the following stochastic differential equations:

• Euclidean space \mathbb{R}^n (null curvature):

$$dR_t = \sqrt{2}dB_t + \frac{n+1}{R_t}dt$$

(Bessel process of dimension n + 2, up to scaling time by 1/2).

Furthermore when n = 2, it can be proved that starting from any $D \in \mathcal{D}$, the normalized domain $\mathfrak{D}_t/\sqrt{\mu(\mathfrak{D}_t)}$ converges to the disk of diameter $1/\sqrt{\pi}$ for large times (under the restriction that (5) can be solved for any time $t \ge 0$).

• Spherical space \mathbb{S}^n (positive curvature=1):

$$dR_t = \sqrt{2}dB_t + \left(\frac{2\sin^{n-1}(R_t)}{\int_0^{R_t}\sin^{n-1}(z)\,dz} - (n-1)\cot(R_t)\right)dt$$

Enable to construct strong stationary times (should lead to a cut-off phenomenon with respect to the dimension).

• Poincaré's model of hyperbolic space \mathbb{H}^n (negative curvature=-1)

$$dR_t = \sqrt{2}dB_t + \left(\frac{2\sinh^{n-1}(R_t)}{\int_0^{R_t}\sinh^{n-1}(z)\,dz} - (n-1)\coth(R_t)\right)\,dt$$

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To define rigorously the generator \mathcal{G} , we must have at our disposal "nice observables".

• Elementary observables:

$$F_f : \mathcal{D} \ni D \mapsto F_f(D) \coloneqq \int_D f d\mu$$

associated to the functions $f \in C^{\infty}(V)$, the space of smooth mappings on V.

• **Composite observables**: the functionals of the form $\mathfrak{F} := \mathfrak{f}(F_{f_1}, ..., F_{f_n})$, where $n \in \mathbb{Z}_+$, $f_1, ..., f_n \in \mathcal{C}^{\infty}(V)$ and $\mathfrak{f} : \mathcal{R} \to \mathbb{R}$ is a \mathcal{C}^{∞} mapping, with \mathcal{R} an open subset of \mathbb{R}^n containing the image of \mathcal{D} by $(F_{f_1}, ..., F_{f_n})$.

On elementary observables:

$$\forall D \in \mathcal{D}, \quad \mathcal{G}[F_f](D) := \int_D G[f] \, d\mu + 2 \frac{\mu(\partial D)}{\mu(D)} \int_{\partial D} f \, d\underline{\mu}$$

For the extension to composite observables, the **carré du champs** is also required:

$$\forall D \in \mathcal{D}, \quad \Gamma_{\mathcal{G}}[F_f, F_g](D) = \left(\int_{\partial D} f d\underline{\mu}\right) \left(\int_{\partial D} g d\underline{\mu}\right)$$

Then on composite observables \mathfrak{F} as above:

$$\mathcal{G}[\mathfrak{F}] = \sum_{j \in \llbracket 1, n \rrbracket} \partial_j \mathfrak{f}(F_{f_1}, \dots, F_{f_n}) \mathcal{G}[F_{f_j}] + \sum_{k, l \in \llbracket 1, n \rrbracket} \partial_{k, l} \mathfrak{f}(F_{f_1}, \dots, F_{f_n}) \Gamma_{\mathcal{G}}[F_{f_k}, F_{f_l}]$$

(consequence of the continuity of the trajectories of \mathfrak{D}).

The above constructions can be extended to any elliptic second order differential generator G on a manifold V admitting an invariant measure μ . The definition of the generator G is exactly the same, but there is a difference in the description of the infinitesimal evolution of the boundaries.

The operator \mathcal{G} induces on V a Riemannian structure so that $G = \triangle + b$, where b is a vector field. Write $\exp(U)$ the density of μ with respect to the Riemannian measure. Then b admits a (weighted Hodge) decomposition $\nabla U + \beta$. The s.d.e. (5) must be replaced by

$$dY_{t} = \left(\sqrt{2}dB_{t} + \left(2\frac{\underline{\mu}(\partial\mathfrak{D}_{t})}{\mu(\mathfrak{D}_{t})} + \langle\beta - \nabla U, \nu\rangle(Y_{t}) - \rho(Y_{t})\right)dt\right)\nu(Y_{t})$$

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Up to the stopping time until which everything is well-defined, we always have:

Theorem

The volume process $(\mu(\mathfrak{D}_{\theta_t}))_{t\geq 0}$ is a Bessel process of dimension 3, where the time change is given by

$$2\int_0^{\theta_t} (\underline{\mu}(\partial\mathfrak{D}_s))^2 \, ds = t$$

The ubiquity of the Bessel-3 process suggests that hypoellipticity in general could be investigated in a similar probabilistic way.

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