

Introduction to set-valued intertwining duality for Markov processes

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Plan of the talk

- 1 Top-to-random shuffle [Aldous and Diaconis, 1986]
- 2 Strong stationary times and Markov intertwining relations
- 3 Set-valued duals
- 4 Dubins' example [1968] and Pitman's theorem [1975, Pitman and Rogers 1981]
- 5 One-dimensional diffusions
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Decks of cards

N cards labeled $1, 2, \dots, N$. A deck of these cards is represented by an element σ of the symmetric group \mathcal{S}_N :

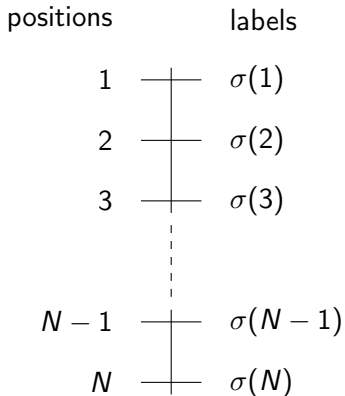


Figure: The deck of cards σ

Top-to-random shuffle

Starting from the ordered deck $(1, 2, \dots, N)$, shuffle it as follows, where the positions l_1, l_2, l_3, \dots are independent and uniformly distributed on $\{1, 2, \dots, N\}$:

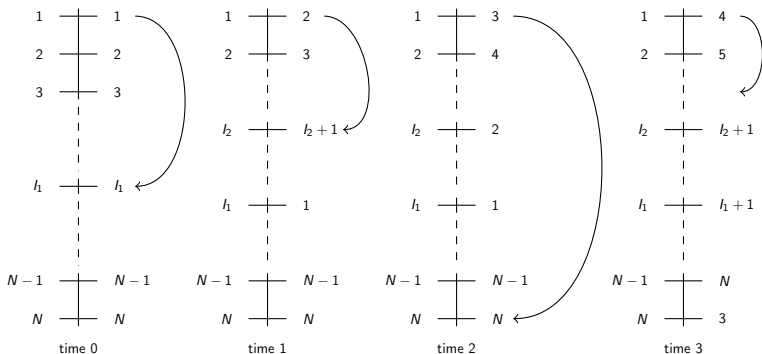


Figure: An example of top-to-random shuffle, here $l_3 = N$

Let $X(n) \in \mathcal{S}_N$ be the deck at time $n \in \mathbb{Z}_+$. The random sequence $X := (X(n))_{n \in \mathbb{Z}_+}$ is a **Markov chain**: to construct $X(n+1)$, only the knowledge of $X(n)$ (and of I_{n+1} , which is independent from the past) is needed, not the way we ended up with $X(n)$.

The uniform distribution $\mathcal{U}_{\mathcal{S}_N}$ on \mathcal{S}_N is **invariant** for X : if $X(0) \sim \mathcal{U}_{\mathcal{S}_N}$, then for any time $n \in \mathbb{Z}_+$, $X(n) \sim \mathcal{U}_{\mathcal{S}_N}$. Furthermore, by irreducibility and aperiodicity, the law $\mathcal{L}(X(n))$ converges to $\mathcal{U}_{\mathcal{S}_N}$ for large time n .

But from a practical point of view, we want to know for how long we have to shuffle the deck to be close to $\mathcal{U}_{\mathcal{S}_N}$: we are looking for **quantitative convergence** estimates.

Quantitative convergence (1)

For the top-to-random Markov chain, there is a probabilistic way to get such an estimate. Look at the last card N at time 0. This card goes up, until it reaches the top of the deck, say at random time T . At time $\tau := T + 1$ the card N is sent to a uniformly chosen position.

At time τ the deck is uniformly distributed. Furthermore τ is a stopping time independent from the ordering of the deck $X(\tau)$. This property implies that in total variation,

$$\forall n \in \mathbb{Z}_+, \quad \|\mathcal{L}(X(n)) - \mathcal{U}_{S_N}\|_{\text{tv}} \leq \mathbb{P}[\tau > n]$$

To deduce a quantitative bound, observe that τ is distributed as a sum of independent geometric random variables of parameters $1/N$, $2/N$, ..., $(N-1)/N$ and 1.

In particular, we get

$$\begin{aligned}\mathbb{E}[\tau] &= N + \frac{N}{2} + \frac{N}{3} + \cdots + \frac{N}{N} \\ &\sim N \ln(N)\end{aligned}$$

which suggests that $N \ln(N)$ shuffles are needed to be at equilibrium.

Recognizing the coupon collector problem, one can be more precise: for any $c > 0$,

$$\|\mathcal{L}(X(\lfloor N \ln(N) + cN \rfloor)) - \mathcal{U}_{S_N}\|_{\text{tv}} \leq \exp(-c)$$

Going even further, one can deduce an instance of the **cut-off** phenomenon at time $N \ln(N)$, in separation and total variation.

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Strong stationary times

The above considerations extend into a much more general theory, see in particular [Diaconis and Fill, 1990].

Consider $X := (X(n))_{n \in \mathbb{Z}_+}$ a Markov chain taking values in a finite state space V and admitting a unique invariant probability π . A finite stopping time τ (with respect to the filtration generated by X and some independent randomness) is said to be a **strong stationary time** when

$$\tau \perp\!\!\!\perp X(\tau) \quad \text{and} \quad X(\tau) \sim \pi$$

An example is provided by the first time the last card is inserted in the deck in the top-to-random model.

Strong stationary times enable probabilistic exact simulations of π .

The **separation** discrepancy between two probability measures μ and π on the same state space is defined by:

$$\mathfrak{s}(\mu, \pi) := \operatorname{ess\,sup}_{\mu} 1 - \frac{d\mu}{d\pi} \geq \|\mu - \pi\|_{\text{tv}}$$

(recall the total variation: $\|\mu - \pi\|_{\text{tv}} := \sup_{A \text{ event}} |\mu(A) - \pi(A)|$, coinciding with $\frac{1}{2} \left\| \frac{d\mu}{d\pi} - 1 \right\|_{\mathbb{L}^1(\pi)}$ when $\mu \ll \pi$).

A general property of any strong stationary time τ :

$$\forall n \in \mathbb{Z}_+, \quad \mathfrak{s}(\mathcal{L}(X(n)), \pi) \leq \mathbb{P}[\tau > n]$$

Furthermore, as soon as X is ergodic, there exist strong stationary times τ such that all these inequalities are equalities, and such τ are said to be a **sharp strong stationary time**.

Intertwining relations (1)

Denote P the transition matrix of X on V and consider a transition matrix Q on another finite state space W , as well a Markov kernel Λ from W to V , called the **link**. There is an **algebraic intertwining Markov relation** from Q to P when the following commutation relation holds:

$$Q\Lambda = \Lambda P \quad (1)$$

Let $Y := (Y(n))_{n \in \mathbb{Z}_+}$ be a Markov chain whose transitions are dictated by Q . When in addition to (1), we have

$$\mathcal{L}(X(0)) = \mathcal{L}(Y(0))\Lambda$$

we say there is a **algebraic intertwining Markov relation** from Y to X . The Markov chain Y is also called an **intertwining dual** for the **primal** Markov chain X .

Intertwining relations (2)

Assume there is an algebraic intertwining Markov relation from Y to X . Then there exists a coupling of X and Y such that for all times $n \in \mathbb{Z}_+$, we have for the conditional expectations:

$$\mathcal{L}(X(n)|Y(\llbracket 0, n \rrbracket)) = \Lambda(Y(n), \cdot) \quad (2)$$

$$\mathcal{L}(Y(\llbracket 0, n \rrbracket)|X) = \mathcal{L}(Y(\llbracket 0, n \rrbracket)|X(\llbracket 0, n \rrbracket)) \quad (3)$$

where $Y(\llbracket 0, n \rrbracket)$ stands for the partial trajectory $(Y(l))_{l \in \llbracket 0, n \rrbracket}$. The equalities (3) imply that Y can be constructed from X in an “adapted” way. In particular any stopping time for Y is also a stopping for X under such a coupling.

In this situation we say that Y and X satisfies a **probabilistic intertwining Markov relation** or that they are **coupled by intertwining**.

Assume that there exists an **absorbing state** $\omega \in W$ for Q , $Q(\omega, \omega) = 1$, and that furthermore $\{\omega\}$ is the unique recurrence class for Q . In particular, whatever the initial distribution of $Y(0)$, the absorbing time

$$\tau := \inf\{t \in \mathbb{Z}_+ : Y(n) = \omega\}$$

is a.s. finite.

When X and Y are coupled by intertwining, then the absorbing time τ is a strong stationary time for X . Indeed, an important observation is that $\Lambda(\omega, \cdot) = \pi$, coming from (1) applied at ω . This is probably the most common method to construct strong stationary times.

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Set-valued duals

A particular instance of the previous framework corresponds to

$$W := \{S \subset V, S \neq \emptyset\}$$
$$\forall S \in W, \quad \Lambda(S, \cdot) = \frac{\pi(S \cap \cdot)}{\pi(S)}$$

at least when π is positive (typically for irreducible Markov chains X). The corresponding dual chain Y is then set-valued.

The whole set V usually turns out to be absorbing for Y and we recover that $\Lambda(V, \cdot) = \pi$.

The top-to-random example can be written under this form, by considering cylindrical subsets of \mathcal{S}_N corresponding to decks of cards coinciding with the current deck on the positions above the last card.

Diaconis and Fill [1990] have shown on a three-point example, that sharp set-valued duals may not exist.

Assume that π is positive and consider P^* the adjoint Markov kernel of P is $\mathbb{L}^2(\pi)$ (it is the transition kernel of the time-reversed Markov chain of X at equilibrium):

$$\forall x, y \in V, \quad P^*(x, y) := \frac{\pi(y)}{\pi(x)} P(y, x)$$

A transition kernel J on $\bar{W} := W \sqcup \{\emptyset\}$ is constructed as follows: given $S \subset V$ and a random variable U uniformly distributed on $[0, 1]$, consider the random subset

$$\Phi(S) := \{y \in V : P^*(y, S) \geq U\}$$

We define

$$\forall S, S' \in \bar{W}, \quad J(S, S') := \mathbb{P}[\Phi(S) = S']$$

The **evolving set processes** of Morris and Peres are the \bar{W} -valued Markov chains $Z := (Z(n))_{n \in \mathbb{Z}_+}$ whose transition kernel is J .

Deduction of a set-valued dual

The Markov chain Z is absorbed at \emptyset and at V . The mapping $\pi : \bar{W} \ni S \mapsto \pi(S)$ is harmonic for J , in the sense that $J\pi = \pi$. It leads to consider the **Doob transform** of J by π , which is the $W \times W$ transition matrix Q given by

$$\forall S, S' \in W, \quad Q(S, S') := \frac{\pi(S')}{\pi(S)} J(S, S')$$

It is the transition kernel of Z conditioned not to be absorbed at \emptyset . Let $Y := (Y(n))_{n \in \mathbb{Z}_+}$ be a Markov chain whose transition kernel is Q and such that $\mathcal{L}(Y(0))\Lambda = \mathcal{L}(X(0))$. Then Y is a set-valued intertwining dual for X .

Random mappings (1)

The previous construction can be generalised via an approach similar to the coupling-from-the-past algorithm of Propp and Wilson [1996]. For any $S \in \bar{W}$, let be given a random mapping $\psi_S : V \rightarrow V$ such that

$$\forall x \in V, \forall x' \in S, \quad \mathbb{P}[\psi_S(x) = x'] = \frac{P^*(x, x')}{\zeta(S)}$$

where ζ is a positive mapping on \bar{W} . The family $(\psi_S)_{S \in \bar{W}}$ enables to define a random mapping Ψ from \bar{W} to \bar{W} via

$$\forall S \in \bar{W}, \quad \Psi(S) := \{y \in V : \psi_S(y) \in S\}$$

Consider the transition matrix K from \bar{W} to \bar{W} given by

$$\forall S, S' \in \bar{W}, \quad K(S, S') := \mathbb{P}[\Psi(S) = S']$$

and its modified Doob transform via

$$\forall S, S' \in W, \quad Q(S, S') := \frac{\pi(S')\zeta(S)}{\pi(S)} K(S, S')$$

Random mappings (2)

Introduce the following conditioned transition: for $x, x' \in V$ such that $P(x, x') > 0$ and any $S \in W$ containing x ,

$$\forall S' \in W, \quad K_{x, x'}(S, S') := \mathbb{P}[\Psi(S) = S' | \psi_S(x') = x]$$

Consider

$$\mathcal{W} := \{(x, S) \in V \times W : x \in S\}$$

and let \mathcal{P} be the set of probability measures m on \mathcal{W} which can be written under the form

$$\forall (x, S) \in \mathcal{W}, \quad m(x, S) = \mu(S) \Lambda(S, x)$$

where μ is the marginal of m on W .

Random mappings (3)

Define a Markov kernel Q on \mathcal{W} via

$$\forall (x, S), (x', S') \in \mathcal{W}, \quad Q((x, S), (x', S')) := P(x, x')K_{x, x'}(S, S')$$

Theorem

Let $(X(n), Y(n))_{n \in \mathbb{Z}_+}$ be a Markov chain on \mathcal{W} whose initial distribution $\mathcal{L}(X_0, Y_0)$ belongs to \mathcal{P} and whose transitions are given by Q . Then $X := (X(n))_{n \in \mathbb{N}}$ and $Y := (Y(n))_{n \in \mathbb{N}}$ are Markov chains whose respective transitions are given by P and Q and Y is a set-valued intertwining dual for X .

This result was used to construct strong stopping times for random walks on discrete Heisenberg groups.

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Dubins' example [1968]

Consider for X the Brownian motion on $[0, 1]$, reflected at 0 and 1 and starting from $1/2$. It is positive recurrent with the restriction of the Lebesgue measure as invariant probability. A strong stationary time can be constructed as follows: let τ_1 be the first time X hits $1/4$ or $3/4$. Next let τ_2 be the first time after τ_1 that $X(\tau_1) \pm 1/8$ is reached. Iteratively, τ_{n+1} is the first time after τ_n that $X(\tau_n) \pm 1/2^{n+2}$ is hit. The limit $\tau := \lim_{n \rightarrow \infty} \tau_n$ exists a.s. and is a strong stationary time for X .

This construction can be extended to any initial distribution, by first waiting that $1/2$ is reached (not always smart, for instance if $X(0)$ was already at equilibrium).

Pitman's transformation

Let $B := (B(t))_{t \geq 0}$ be a standard Brownian motion. Consider the Pitman's transformation $R := (R(t))_{t \geq 0}$ given by

$$R(t) := -B(t) + 2 \max_{s \in [0, t]} B(s)$$

Let Λ the Markov kernel from \mathbb{R}_+ to \mathbb{R} given by

$$\forall r \geq 0, \quad \Lambda(r, dx) := \mathcal{U}_{[-r, r]}(dx)$$

Relations (2) and (3) can be extended to this continuous setting: for all time $t \geq 0$,

$$\begin{cases} \mathcal{L}(B(t) | R([0, t])) = \Lambda(R(t), \cdot) \\ \mathcal{L}(R([0, t]) | B) = \mathcal{L}(R([0, t]) | B([0, t])) \end{cases} \quad (4)$$

Furthermore, the process R is a **Bessel-3 process**, i.e. it has the law of the norm of a Brownian motion in dimension 3.

Pitman's theorem in pictures

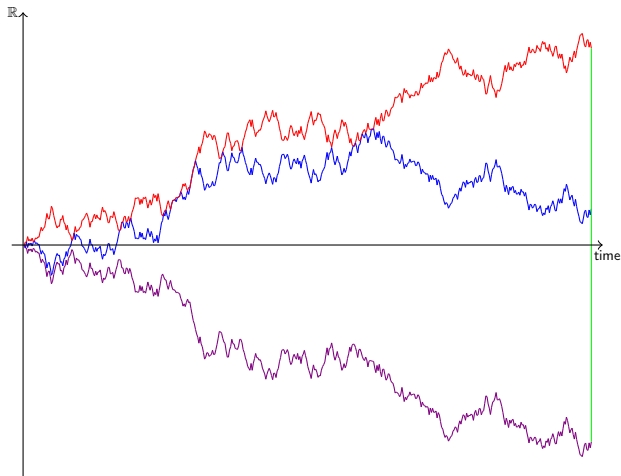


Figure: Trajectories: Brownian motion $B([0, t])$, $R([0, t])$, $-R([0, t])$, and the segment-valued dual: $[-R(t), R(t)]$

Consequently of (4), we get estimates on the convergence of the Brownian motion $W := (W(t))_{t \geq 0}$ on the circle $\mathbb{T} := \mathbb{R}/(2\pi\mathbb{Z})$: let τ be the hitting time of π by R :

$$\tau := \inf\{t \geq 0 : R(t) = \pi\}$$

It is a strong stationary time for W . The hitting times of Bessel processes are well-studied, we have:

$$\forall \lambda > 0, \quad \mathbb{E}[\exp(-\lambda\tau)] = \frac{\sqrt{\pi} \sqrt[4]{\lambda}}{\sqrt[4]{2} \Gamma(3/2)} \frac{1}{I_{1/2}(\pi\sqrt{2\lambda})}$$

where $I_{1/2}$ is the modified Bessel function of index $1/2$.

A Tauberien theorem enables to deduce the behavior for large $t \geq 0$ of $\mathbb{P}[\tau > t]$ and thus of $\mathfrak{s}(\mathcal{L}(W(t)), \mathcal{U}_{\mathbb{T}})$ and $\|\mathcal{L}(W(t)) - \mathcal{U}_{\mathbb{T}}\|_{\text{tv}}$.

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Elliptic real diffusions

On \mathbb{R} , consider a diffusion X whose generator is $G := a\partial^2 + b\partial$, where a, b are smooth functions and a is positive. Consider the **speed measure** $\mu(dx) := \underline{\mu}(x) dx$ with density

$$\underline{\mu}(x) := \frac{\exp(c(x))}{a(x)} \quad \text{with} \quad c(x) := \int_0^x \frac{b(y)}{a(y)} dy$$

It is invariant for X and even **reversible**: G can be extended into a self-adjoint operator on $\mathbb{L}^2(\mu)$.

Consider

$$\mathcal{D} := \{[y, z] : y, z \in (-\infty, +\infty), y \leq z\}$$

As usual, define the associated Markov kernel Λ from \mathcal{D} to \mathbb{R} via

$$\forall [y, z] \in \mathcal{D}, \Lambda([y, z], dx) := \begin{cases} \delta_y(dx) & , \text{ if } y = z \\ \frac{\underline{\mu}(x)}{\mu([y, z])} \mathbb{1}_{[y, z]}(x) dx & , \text{ otherwise} \end{cases}$$

Dual process

On \mathcal{D} define the degenerate generator

$$\begin{aligned} \mathcal{G} &:= (\sqrt{a(z)}\partial_z - \sqrt{a(y)}\partial_y)^2 \\ &\quad + (a'(y)/2 - b(y))\partial_y + (a'(z)/2 - b(z))\partial_z \\ &\quad + 2 \frac{\sqrt{a(y)}\underline{\mu}(y) + \sqrt{a(z)}\underline{\mu}(z)}{\mu([y, z])} (\sqrt{a(z)}\partial_z - \sqrt{a(y)}\partial_y), \end{aligned}$$

Its interest is in the intertwining relation $\mathcal{G}\Lambda = \Lambda G$.

Proposition

There exists a unique process $[Y, Z]$ whose generator is \mathcal{G} , up to its explosion time ζ . The diagonal is an entrance boundary for this process.

The proof is based on the fact that $(\mu([Y(t), Z(t)]))_{t \in [0, \zeta]}$ is a (stopped) Bessel-3 process, up to the time-change $(\theta(t))_{t \in [0, \zeta]}$ given by

$$2 \int_0^{\theta(t)} (\sqrt{a(Y(s))}\underline{\mu}(Y(s)) + \sqrt{a(Z(s))}\underline{\mu}(Z(s)))^2 ds = t$$

Strong stationary times in 1-dimension

Application to the existence of strong stationary times:

Theorem

Assume that X is positive recurrent. There exists a strong stationary time for X , whatever its initial distribution, if and only if $-\infty$ and $+\infty$ are entrance boundaries.

Positive recurrence or ergodicity: μ is finite and in large time $X(t)$ converges to the renormalization of μ , analytically this is characterized by

$$\int_{-\infty}^0 \exp(-c(y)) dy = +\infty \quad \text{and} \quad \int_0^{+\infty} \exp(-c(y)) dy = +\infty$$

Entrance boundary: e.g. for $+\infty$, it means that X can be started from $+\infty$ (comes down from infinity), it amounts to:

$$\int_0^{+\infty} \left(\int_0^x \exp(-c(y)) dy \right) \mu(dx) < +\infty$$

1-dimension hypoelliptic diffusions

Consider again the generator $G := a\partial^2 + b\partial$, on \mathbb{R} or \mathbb{T} , where a is allowed to vanish on a finite number of points, but such that \sqrt{a} remains smooth. On the vanishing points, assume b does not vanish: 1-dimensional hypoellipticity in the sense of [Hörmander, 1967].





Up to some adjustments (modification of the measure used in Λ , dual processes which can disconnect, ...), the above approach is valid and enables to recover the density theorem (i.e. the law of $X(t)$ admits a density for all $t > 0$) and to obtain estimates on the speed of convergence to equilibrium (even when the invariant measure does not charge the whole state space). The Bessel-3 process is still there, the hypoellipticity is only felt at the level of the time change.

Is it possible to recover the generality of Hörmander's theorem in this probabilistic way?




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