

Macroscopic models for collective sperm-cells dynamics and cells inclination

Laurent Navoret

Joint work with Pierre Degond

ANR Motimo

Workshop “Collective dynamics of active particles, swimmers, motile cells”
IMFT, Toulouse, 23 September 2015

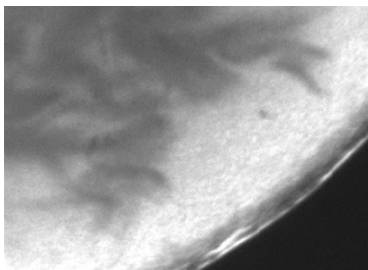
ANR Motimo

→ spermatozoa in seminal liquid

- $5 \cdot 10^9$ cells per cm^3
- volume ratio 50%

→ collective movement

- mass motility
- dark waves
- fertility measure



[X. Druart]

Spermatozoa : collective movement and cell geometry

- cell geometry
 - head : flat ellipsoids
 - flagella : neglected
- head inclination \Rightarrow dark zone

Collective displacements

system of interacting particles

$$(x_k(t), v_k(t)), k \in \{1, \dots, N\}$$



macroscopic model

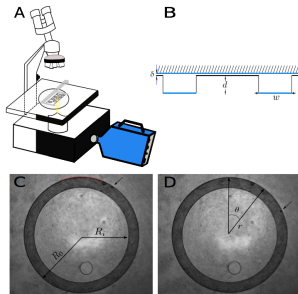
$$(\rho(x, t), u(x, t))$$

Interactions

- hydrodynamic interactions
- volume-exclusion interactions
predominant at large density

⇒ **local alignment** for self-propelled elongated particles [Peruani, 2006]

⇒ experiment validation : spontaneous rotation [Creppy et al.]



[Creppy et al., to appear]

Modeling assumptions

- self-propulsion
- local alignment interactions with neighbors (volume-exclusion interactions)
- geometry : flat disks

Context

- system of interacting particles
- self-propelled dynamics with alignment interactions : Vicsek dynamics, phase transition, band formation [Vicsek, Chaté, etc.]
- derivation of macroscopic model with kinetic theory [Degond, Motsch, 2008]

Plan

- ① Particle model
- ② Derivation of the macroscopic model
- ③ Macroscopic model
- ④ Numerical simulations
- ⑤ Conclusion

Plan

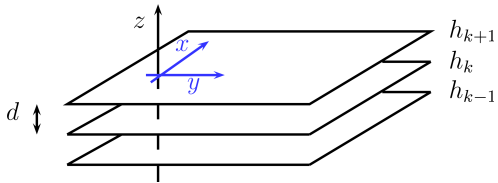
- 1 Particle model
- 2 Derivation of the macroscopic model
- 3 Macroscopic model
- 4 Numerical simulations
- 5 Conclusion

Particle model

- **2D flow**

Displacements in layers

Layers : horizontal planes
with inter-distance d



- Motion inside layers

(altitude) $h_k \in \mathbb{R}$

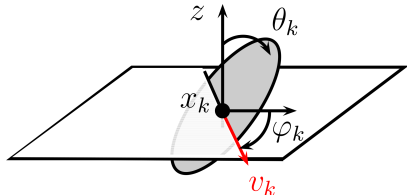
(position) $x_k(t) \in \mathbb{R}^2$

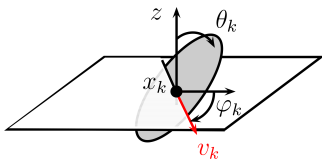
(velocity) $v_k(t) \in \mathbb{S}^1 \rightarrow \varphi_k(t) \in [0, 2\pi]$

- Spatial configuration of disks

$v_k \in$ Disk plane

(inclination) $\theta_k(t) \in [0, \pi]$





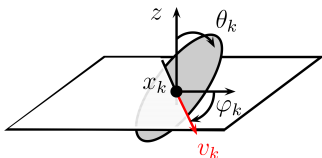
Microscopic model

$$\frac{dx_k}{dt} = v(\varphi_k),$$

$$d\varphi_k = -\nu \sin(\varphi_k - \bar{\varphi}_k) dt + \sqrt{2D} dB_t^{\varphi,k} \quad \text{modulo } 2\pi,$$

$$d\theta_k = (-K \sin(2(\theta_k - \bar{\theta}_k)) + \mu T_k) dt + \sqrt{2\delta} dB_t^{\theta,k} \quad \text{modulo } \pi,$$

- **self-propulsion : unit speed** ($v_k = v(\varphi_k) = (\cos \varphi_k, \sin \varphi_k)^T$)



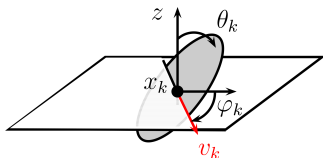
Microscopic model

$$\frac{dx_k}{dt} = v(\varphi_k),$$

$$d\varphi_k = -\nu \sin(\varphi_k - \bar{\varphi}_k) dt + \sqrt{2D} dB_t^{\varphi,k} \quad \text{modulo } 2\pi,$$

$$d\theta_k = (-K \sin(2(\theta_k - \bar{\theta}_k)) + \mu T_k) dt + \sqrt{2\delta} dB_t^{\theta,k} \quad \text{modulo } \pi,$$

- **self-propulsion : unit speed** ($v_k = v(\varphi_k) = (\cos \varphi_k, \sin \varphi_k)^T$)
- **Alignment in velocity : relaxation to averaged $\bar{\varphi}_k$**
 - local alignment with the disks of the same layer
 - local alignment with the disks the neighboring layers



Microscopic model

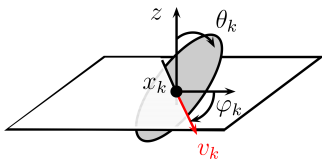
$$\frac{dx_k}{dt} = v(\varphi_k),$$

$$d\varphi_k = -\nu \sin(\varphi_k - \bar{\varphi}_k) dt + \sqrt{2D} dB_t^{\varphi,k} \quad \text{modulo } 2\pi,$$

$$d\theta_k = (-K \sin(2(\theta_k - \bar{\theta}_k)) + \mu T_k) dt + \sqrt{2\delta} dB_t^{\theta,k} \quad \text{modulo } \pi,$$

- **self-propulsion : unit speed** ($v_k = v(\varphi_k) = (\cos \varphi_k, \sin \varphi_k)^T$)
- **Alignment in velocity : relaxation to averaged $\bar{\varphi}_k$**
 - local alignment with the disks of the same layer
 - local alignment with the disks the neighboring layers
- **Inclination dynamics**
 - nematic alignment with disks of the same layer : relax to $\bar{\theta}_k$
 - with disks of the neighboring layers : torque w.r.t to $\langle v_k \rangle$: T_k

$$T_k = \sum_{\substack{j, h_j = h_k \pm 1, \\ |X_j - X_k| \leq R}} T_{kj} \quad T_{kj} = \text{sign}(h_j - h_k) (v_k \times v_j) \cdot \hat{z}$$



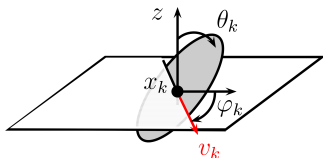
Microscopic model

$$\frac{dx_k}{dt} = v(\varphi_k),$$

$$d\varphi_k = -\nu \sin(\varphi_k - \bar{\varphi}_k) dt + \sqrt{2D} dB_t^{\varphi,k} \quad \text{modulo } 2\pi,$$

$$d\theta_k = (-K \sin(2(\theta_k - \bar{\theta}_k)) + \mu T_k) dt + \sqrt{2\delta} dB_t^{\theta,k} \quad \text{modulo } \pi,$$

- **self-propulsion : unit speed** ($v_k = v(\varphi_k) = (\cos \varphi_k, \sin \varphi_k)^T$)
- **Alignment in velocity** : relaxation to averaged $\bar{\varphi}_k$
 - local alignment with the disks of the same layer
 - local alignment with the disks the neighboring layers
- **Inclination dynamics**
 - nematic alignment with disks of the same layer : relax to $\bar{\theta}_k$
 - with disks of the neighboring layers : torque w.r.t to $\langle v_k \rangle$: T_k
- **diffusion** in velocity and inclination : $B_t^{\varphi,k}$, $B_t^{\theta,k}$ Brownian motion



Microscopic model

$$\frac{dx_k}{dt} = v(\varphi_k),$$

$$d\varphi_k = -\nu \sin(\varphi_k - \bar{\varphi}_k) dt + \sqrt{2D} dB_t^{\varphi,k} \quad \text{modulo } 2\pi,$$

$$d\theta_k = (-K \sin(2(\theta_k - \bar{\theta}_k)) + \mu T_k) dt + \sqrt{2\delta} dB_t^{\theta,k} \quad \text{modulo } \pi,$$

- **self-propulsion : unit speed** ($v_k = v(\varphi_k) = (\cos \varphi_k, \sin \varphi_k)^T$)
- **Alignment in velocity** : relaxation to averaged $\bar{\varphi}_k$
 - local alignment with the disks of the same layer
 - local alignment with the disks the neighboring layers
- **Inclination dynamics**
 - nematic alignment with disks of the same layer : relax to $\bar{\theta}_k$
 - with disks of the neighboring layers : torque w.r.t to $\langle v_k \rangle$: T_k
- **diffusion** in velocity and inclination : $B_t^{\varphi,k}$, $B_t^{\theta,k}$ Brownian motion

⇒ several parameters : ν , μ , K , D , δ

Plan

- ① Particle model
- ② Derivation of the macroscopic model
- ③ Macroscopic model
- ④ Numerical simulations
- ⑤ Conclusion

Mean-field kinetic equation

- $f(x, \varphi, \theta, h, t)$ distribution function in phase space
 $x \in \mathbb{R}^2, \varphi \in [0, 2\pi], \theta \in [0, \pi], h \in \mathbb{N}$
- Mean-field model

$$\partial_t f + \nabla_x \cdot (v(\varphi)f) = -\partial_\varphi \left(-\nu \sin(\varphi - \bar{\varphi}_f) f \right) + D \partial_\varphi^2 f \\ - \partial_\theta \left((-K \sin(2(\theta - \bar{\theta}_f)) + \mu T_f) f \right) + \delta \partial_\theta^2 f$$

with $v(\varphi) = (\cos \varphi, \sin \varphi)^T$.

Mean inclination angle $\bar{\theta}_f$:

$$e^{2i\bar{\theta}_f} = \frac{J_{R,f}^\theta}{|J_{R,f}^\theta|} \quad J_{R,f}^\theta(x, h, t) = \int_{\substack{\theta \in [0, \pi], \varphi \in [0, 2\pi], \\ y \in \mathbb{R}^2, |y-x| \leq R}} e^{2i\theta} f(y, \varphi, \theta, h, t) dy d\varphi d\theta$$

Rescaling

- Large time and space scale (hydrodynamic rescaling)

$$\tilde{x} = \varepsilon x, \quad \tilde{t} = \varepsilon t, \quad \varepsilon \ll 1$$

- interaction frequency inside layers
 \gg interaction frequency between neighboring layers

$$\rightarrow \mu = \varepsilon \mu', \text{ with } \mu' = O(1)$$

$$\rightarrow \beta = \varepsilon \beta', \text{ with } \beta' = O(1)$$

$$\begin{aligned} & \varepsilon \left(\partial_t f^\varepsilon + \nabla_x \cdot (v(\varphi) f^\varepsilon) + \overbrace{\partial_\theta (\mu T_{f^\varepsilon} f^\varepsilon) + \partial_\varphi (S_{f^\varepsilon} f^\varepsilon)}^{\text{interactions with nb. layers}} \right) \\ & = \underbrace{K \partial_\theta (\sin(2(\theta - \bar{\theta}_{f^\varepsilon})) f^\varepsilon) + \delta \partial_\theta^2 f^\varepsilon + \nu \partial_\varphi (\sin(\varphi - \bar{\varphi}_{f^\varepsilon}) f^\varepsilon) + D \partial_\varphi^2 f^\varepsilon}_{\text{interactions inside layers}} \end{aligned}$$

Rescaling

- Large time and space scale (hydrodynamic rescaling)

$$\tilde{x} = \varepsilon x, \quad \tilde{t} = \varepsilon t, \quad \varepsilon \ll 1$$

- interaction frequency inside layers
 \gg interaction frequency between neighboring layers

$$\rightarrow \mu = \varepsilon \mu', \text{ with } \mu' = O(1)$$

$$\rightarrow \beta = \varepsilon \beta', \text{ with } \beta' = O(1)$$

$$\begin{aligned} & \varepsilon \left(\partial_t f^\varepsilon + \nabla_x \cdot (v(\varphi) f^\varepsilon) + \overbrace{\partial_\theta (\mu T_{f^\varepsilon} f^\varepsilon) + \partial_\varphi (S_{f^\varepsilon} f^\varepsilon)}^{\text{interactions with nb. layers}} \right) \\ &= \underbrace{Q(f)}_{\text{interactions inside layers}} \end{aligned}$$

- Operator $Q(f)$:

$$Q(f^\varepsilon) = K \partial_\theta (\sin(2(\theta - \bar{\theta}_{f^\varepsilon})) f^\varepsilon) + \delta \partial_\theta^2 f^\varepsilon + \nu \partial_\varphi (\sin(\varphi - \bar{\varphi}_{f^\varepsilon}) f^\varepsilon) + D \partial_\varphi^2 f^\varepsilon$$

Equilibria

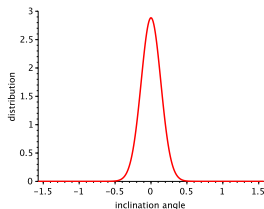
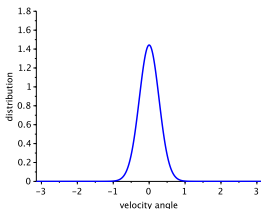
- $Q(f)$ non-linear tensorial Fokker-Planck operator
- Convergence to **local equilibria** for velocity/inclination distribution

$$\mathcal{E} = \{f, Q(f) = 0\} = \{\rho M_{\bar{\varphi}, \bar{\theta}} \mid \rho \in \mathbb{R}^+, \bar{\varphi} \in [0, 2\pi], \bar{\theta} \in [0, \pi]\}$$

with $M_{\bar{\varphi}, \bar{\theta}}$ product of von-Mises function

$$M_{\bar{\varphi}, \bar{\theta}}(\varphi, \theta) = \frac{1}{Z} \exp\left(\frac{\nu}{D} \cos(\varphi - \bar{\varphi})\right) \exp\left(\frac{K}{\delta} \cos(2(\theta - \bar{\theta}))\right)$$

with Z : renormalization constant.



- macroscopic dynamics of $\rho, \bar{\varphi}, \bar{\theta}$

Moment method

- Dynamics of ρ : mass conservation

$$\partial_t \rho + \nabla_x \cdot (c_1 \rho v(\bar{\varphi})) = 0$$

- Dynamics of $\bar{\varphi}, \bar{\theta}$?
 - No conservation of momentum \Rightarrow lack of collisional invariants
dimension of $\mathcal{E} = 3 >$ dimension of collisional invariants = 1
 - [Degond, Motsch, 2008] Generalized collisional invariants

Plan

- ① Particle model
- ② Derivation of the macroscopic model
- ③ Macroscopic model**
- ④ Numerical simulations
- ⑤ Conclusion

Theorem (Macroscopic system)

$$[\text{density}] \quad \partial_t \rho + c_1 \nabla_x \cdot (\rho v(\bar{\varphi})) = 0$$

$$[\text{mean velocity}] \quad \rho \left(\partial_t \bar{\varphi} + c_2 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\varphi} \right) + \frac{1}{\kappa_1} v(\bar{\varphi})^\perp \cdot \nabla_x \rho = \frac{\nu \beta'}{c_3} v(\bar{\varphi})^\perp \cdot \mathcal{S}$$

$$[\text{mean inclination}] \quad \rho \left(\partial_t \bar{\theta} + c_1 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\theta} \right) = \frac{\mu'}{c_4} (c_1 \rho v(\bar{\varphi}))^\perp \cdot \mathcal{T}$$

- Left-hand side :
Self-Organized Hydrodynamics (SOH) model $(\rho, \bar{\varphi})$ + transport $(\bar{\theta})$
- Right-hand side : interaction between layers

$$\mathcal{S} = \sum_{k, k-h=\pm 1} \langle g M_2 M_2 \rangle (\bar{\theta}, \bar{\theta}_k) c_1 \rho_k v(\bar{\varphi}_k)$$

$$\mathcal{T} = \sum_{k, k-h=\pm 1} \text{sign}(k-h) \langle g M_2 M_2 \partial_\theta I_2 \rangle (\bar{\theta}, \bar{\theta}_k) c_1 \rho_k v(\bar{\varphi}_k)$$

Theorem (Macroscopic system)

$$[\text{density}] \quad \partial_t \rho + c_1 \nabla_x \cdot (\rho v(\bar{\varphi})) = 0$$

$$[\text{mean velocity}] \quad \rho \left(\partial_t \bar{\varphi} + c_2 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\varphi} \right) + \frac{1}{\kappa_1} v(\bar{\varphi})^\perp \cdot \nabla_x \rho = \frac{\nu \beta'}{c_3} v(\bar{\varphi})^\perp \cdot \mathcal{S}$$

$$[\text{mean inclination}] \quad \rho \left(\partial_t \bar{\theta} + c_1 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\theta} \right) = \frac{\mu'}{c_4} (c_1 \rho v(\bar{\varphi}))^\perp \cdot \mathcal{T}$$

- Left-hand side :
Self-Organized Hydrodynamics (SOH) model $(\rho, \bar{\varphi})$ + transport $(\bar{\theta})$
- Right-hand side : interaction between layers

$$\mathcal{S} = \sum_{k, k-h=\pm 1} \langle g M_2 M_2 \rangle (\bar{\theta}, \bar{\theta}_k) c_1 \rho_k v(\bar{\varphi}_k)$$

$$\mathcal{T} = \sum_{k, k-h=\pm 1} \text{sign}(k-h) \langle g M_2 M_2 \partial_\theta I_2 \rangle (\bar{\theta}, \bar{\theta}_k) c_1 \rho_k v(\bar{\varphi}_k)$$

- c_1, c_2, c_3, c_4 : weighted averages of von-Mises equilibria

Theorem (Macroscopic system)

$$[\text{density}] \quad \partial_t \rho + c_1 \nabla_x \cdot (\rho v(\bar{\varphi})) = 0$$

$$[\text{mean velocity}] \quad \rho \left(\partial_t \bar{\varphi} + c_2 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\varphi} \right) + \frac{1}{\kappa_1} v(\bar{\varphi})^\perp \cdot \nabla_x \rho = \frac{\nu \beta'}{c_3} v(\bar{\varphi})^\perp \cdot \mathcal{S}$$

$$[\text{mean inclination}] \quad \rho \left(\partial_t \bar{\theta} + c_1 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\theta} \right) = \frac{\mu'}{c_4} (c_1 \rho v(\bar{\varphi}))^\perp \cdot \mathcal{T}$$

- Left-hand side :
Self-Organized Hydrodynamics (SOH) model $(\rho, \bar{\varphi})$ + transport $(\bar{\theta})$

- Right-hand side : interaction between layers

$$\mathcal{S} = \sum_{k, k-h=\pm 1} \langle g M_2 M_2 \rangle (\bar{\theta}, \bar{\theta}_k) c_1 \rho_k v(\bar{\varphi}_k)$$

$$\mathcal{T} = \sum_{k, k-h=\pm 1} \text{sign}(k-h) \langle g M_2 M_2 \partial_\theta I_2 \rangle (\bar{\theta}, \bar{\theta}_k) c_1 \rho_k v(\bar{\varphi}_k)$$

- g overlap function \Rightarrow strong coupling between velocity and inclination

Plan

- ① Particle model
- ② Derivation of the macroscopic model
- ③ Macroscopic model
- ④ Numerical simulations**
- ⑤ Conclusion

Numerical method

$$\partial_t \rho + c_1 \nabla_x \cdot (\rho v(\bar{\varphi})) = 0$$

$$\rho \left(\partial_t \bar{\varphi} + c_2 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\varphi} \right) + \frac{1}{\kappa_1} v(\bar{\varphi})^\perp \cdot \nabla_x \rho = \frac{\nu \beta'}{c_3} v(\bar{\varphi})^\perp \cdot \mathcal{S}$$

$$\rho \left(\partial_t \bar{\theta} + c_1 (v(\bar{\varphi}) \cdot \nabla_x) \bar{\theta} \right) = \frac{\mu'}{c_4} (c_1 \rho v(\bar{\varphi}))^\perp \cdot \mathcal{T}$$

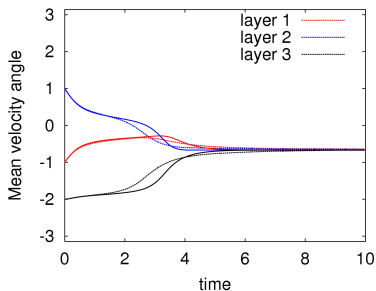
Difficulties :

- no momentum conservation \Rightarrow non-conservative hyperbolic system
- [Motsch, Navoret] Relaxation scheme : at each time step
 - ① solve the left-hand side with a finite volume scheme
 - ② add the source term
 - ③ renormalized the velocity vector field

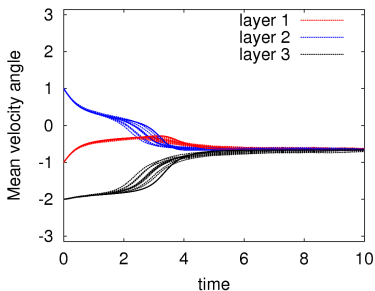
Comparison between particle and macroscopic simulations

micro-macro comparisons

Space homogeneous simulation : 3 layers, **velocity angle** $\bar{\varphi}$



macro / micro (averaged over 20 runs)



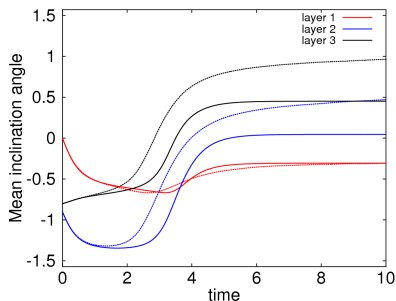
macro / micro (10 runs)

- relaxation to a unique velocity angle
- good agreement micro/macro

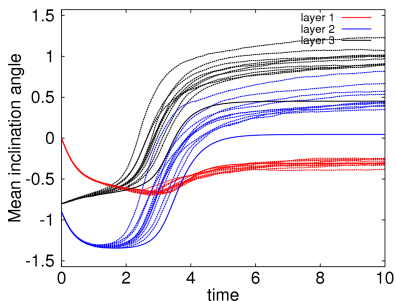
Parameters : 25000 particles per layer

micro-macro comparisons

Space homogeneous simulation : 3 layers, **inclination angle** $\bar{\theta}$



macro / micro (averaged over 20 runs)



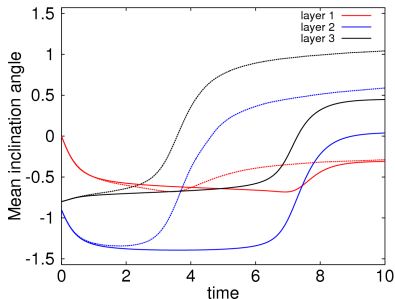
macro / micro (10 runs)

- relaxation in two steps
- stochastic fluctuations (finite number of particles)
→ second relaxation earlier for micro simulations
- after second relaxation : constant inclination

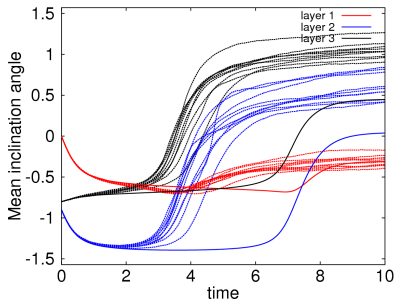
Parameters : 25000 particles per layer, $d = 0.02 = R$

micro-macro comparisons

Space homogeneous simulation : 3 layers, inclination angle $\bar{\theta}$



macro / micro (averaged over 20 runs)



macro / micro (10 runs)

- relaxation in two steps
- stochastic fluctuations (finite number of particles)
→ second relaxation earlier for micro simulations
- after second relaxation : constant inclination

Parameters : 25000 particles per layer, $d = 0.0205 > R = 0.02$

micro-macro comparisons

2D Taylor-Green test-case

→ velocity : normalized and translated Taylor Green vortex

$$v(\bar{\varphi}) = \frac{(u_1, u_2)}{\|(u_1, u_2)\|}$$

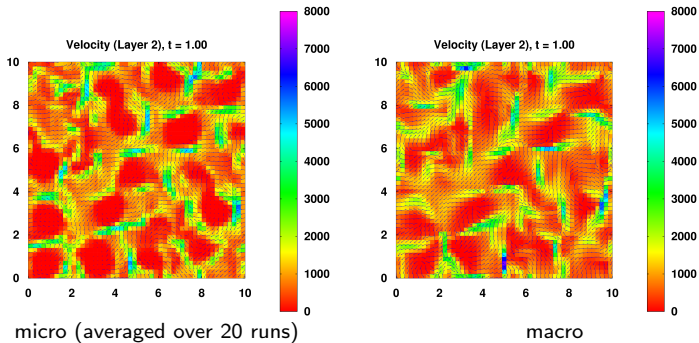
$$u_1(x, y) = \frac{1}{3} \sin\left(\frac{\pi}{5}x\right) \cos\left(\frac{\pi}{5}y\right) + \frac{1}{3} \sin\left(\frac{3\pi}{10}x\right) \cos\left(\frac{3\pi}{10}y\right) + \frac{1}{3} \sin\left(\frac{\pi}{2}x\right) \cos\left(\frac{\pi}{2}y\right)$$

$$u_2(x, y) = -\frac{1}{3} \cos\left(\frac{\pi}{5}x\right) \sin\left(\frac{\pi}{5}y\right) - \frac{1}{3} \cos\left(\frac{3\pi}{10}x\right) \sin\left(\frac{3\pi}{10}y\right) - \frac{1}{3} \cos\left(\frac{\pi}{2}x\right) \sin\left(\frac{\pi}{2}y\right)$$

→ inclination : space homogeneous

micro-macro comparisons

2D Taylor-Green test-case : 3 layers, **density**

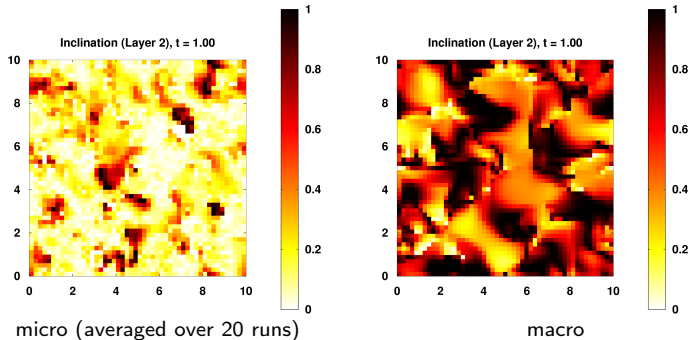


- Layer 2 : intermediate layer
- good agreement
- except in low density region
- macro : more diffuse

Parameters : 10^5 particles per layer

micro-macro comparisons

2D Taylor-Green test-case : **cosine of the inclination angle**



- uniform at $t = 0 \Rightarrow$ non uniform
- large differences between micro/macro
- aligned inclinations do not necessarily match regions of uniform densities

Parameters : 10^5 particles per layer

Plan

- ① Particle model
- ② Derivation of the macroscopic model
- ③ Macroscopic model
- ④ Numerical simulations
- ⑤ Conclusion

Conclusion

- Tensorial alignment problem (velocity, inclination)
- Macroscopic model : hyperbolic system with source terms
- Some deviations for the inclination dynamics but general good agreement between particle/macroscopic simulations

To do :

- characterize the macroscopic motion
- model refinement (density constraint, etc.)
- annular ring geometry

Advertisement

Master 2 of Cell Physics (Université de Strasbourg)

$$2 \rho v_{\text{cell}} = \left(1 + \frac{D}{Dt} \right) \left(\sigma_{\text{cell}} + \Delta \mu \rho_{\text{cell}} + \rho v_{\text{cell}} - \frac{D}{Dt} \left(\rho_1 \rho_2 + \rho_2 \rho_3 - \frac{2}{3} \rho_1 \rho_2 \rho_3 \right) \right)$$



$$\frac{D}{Dt} \rho_{\text{cell}} = \frac{1}{\gamma_1} \rho_{\text{cell}} + \lambda_1 \rho_{\text{cell}} \Delta \mu - \rho_1 v_{\text{cell}} \rho_2 - \rho_2 v_{\text{cell}} \rho_3 - \rho_1 v_{\text{cell}} \rho_2 + \Delta \mu + \rho_1 \rho_2 v_{\text{cell}} + \rho_2 v_{\text{cell}} + \lambda_2 \rho_2 \rho_3$$

Direction : D. Riveline (IGBMC)