



Final MOTIMO workshop, Collective dynamics of active particles, swimmers, motile cells,  
September 22-23, 2015, Institut de Mecanique des Fluides de Toulouse (IMFT).

# Floaterers on Faraday waves: Clustering and heterogeneous flow

@CeydaSanli

September 23, 2015, Toulouse.

CompleXity Networks



Physics of Fluids

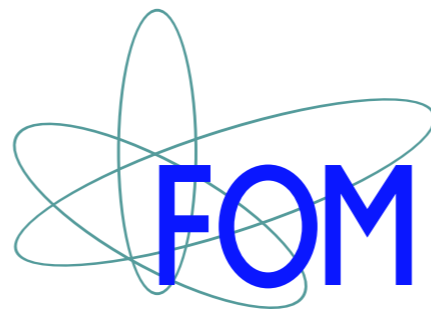


# Greetings, ...

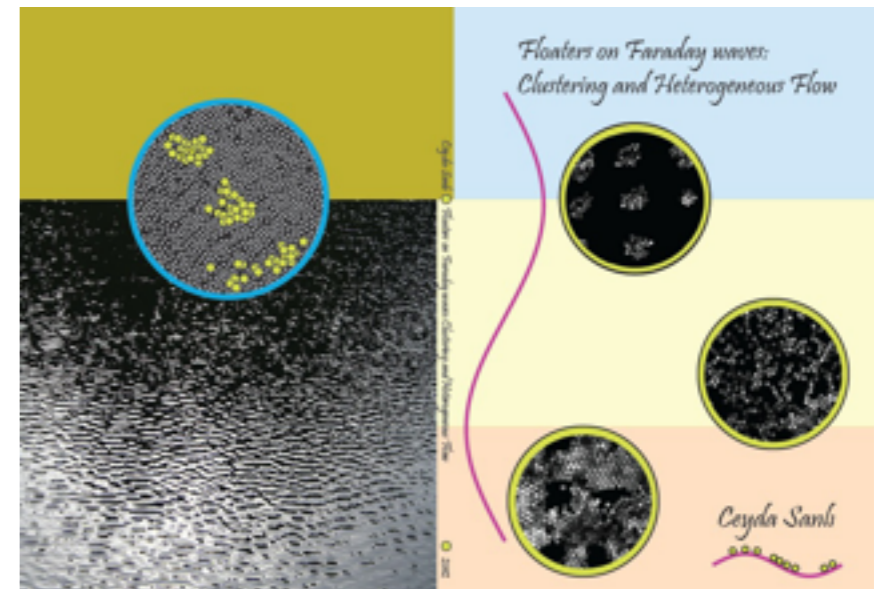


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# Two Directions of my PhD Thesis:

- Clustering on **standing** Faraday waves

- C. Sanlı *et al*,  
Phys. Rev. E **89**, 053011 (2014).

- Heterogeneous flow on **capillary** Faraday waves

- C. Sanlı *et al*,  
Phys. Rev. E **90**, 033018 (2014).

# Observation

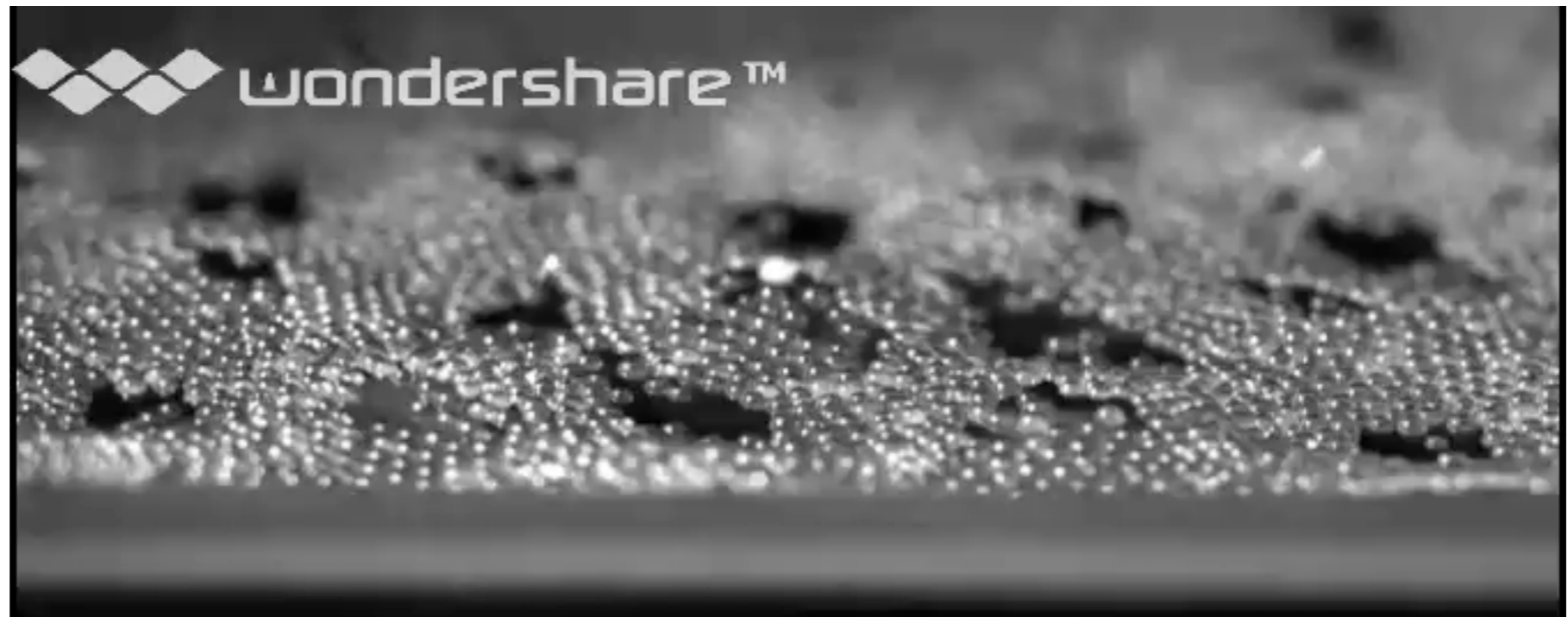
• low  $\phi$ :

- $a \sim 0.8$  mm
- $f = 19$  Hz

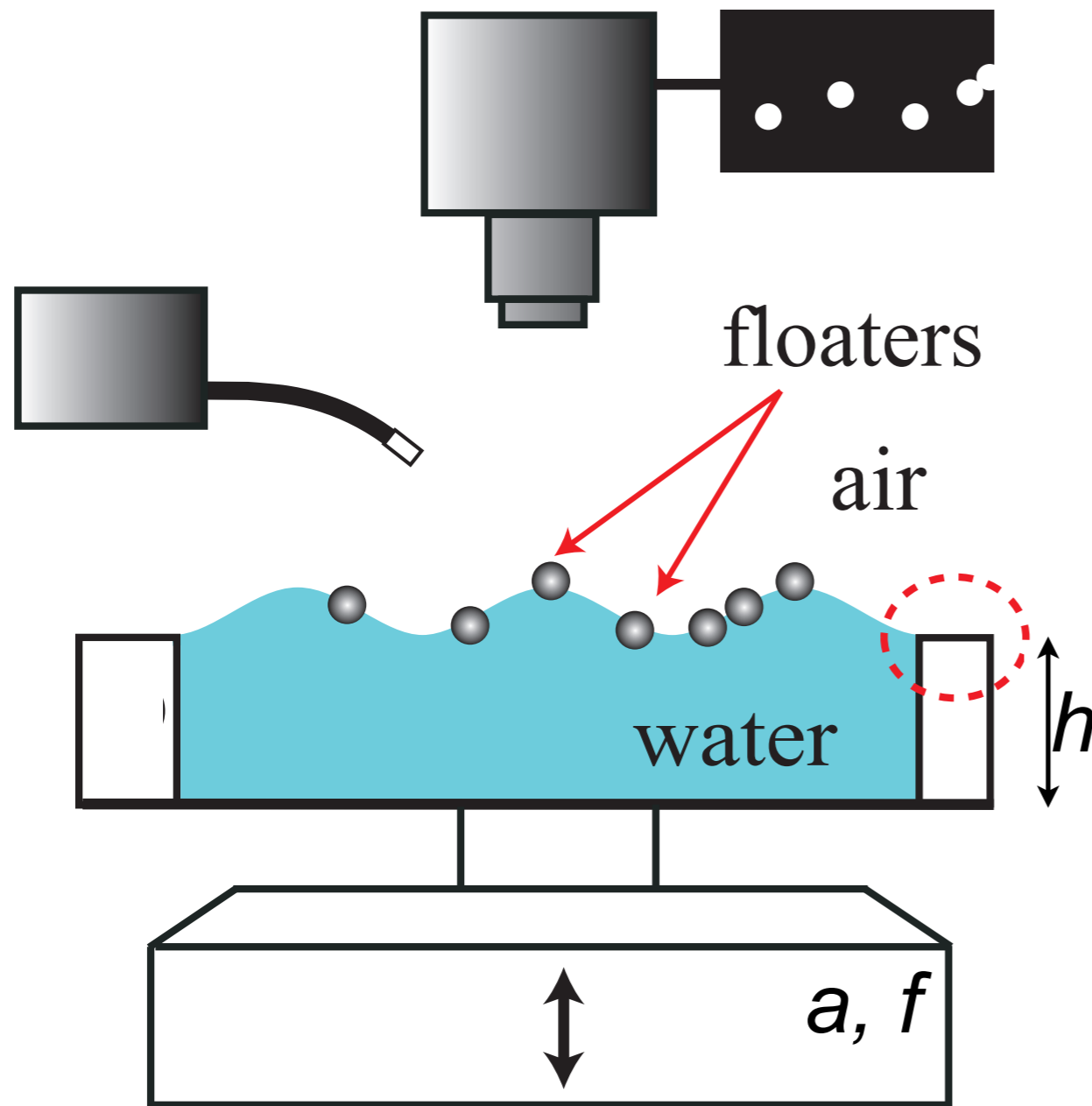


• high  $\phi$ :

- $a \sim 1.2$  mm
- $f = 20$  Hz



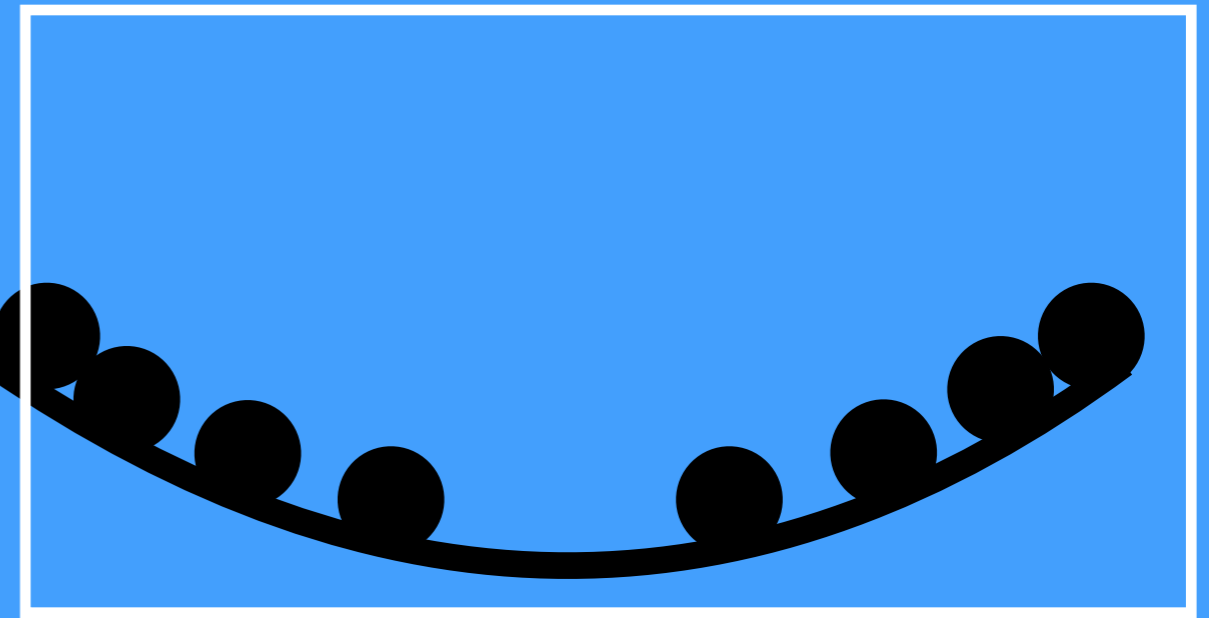
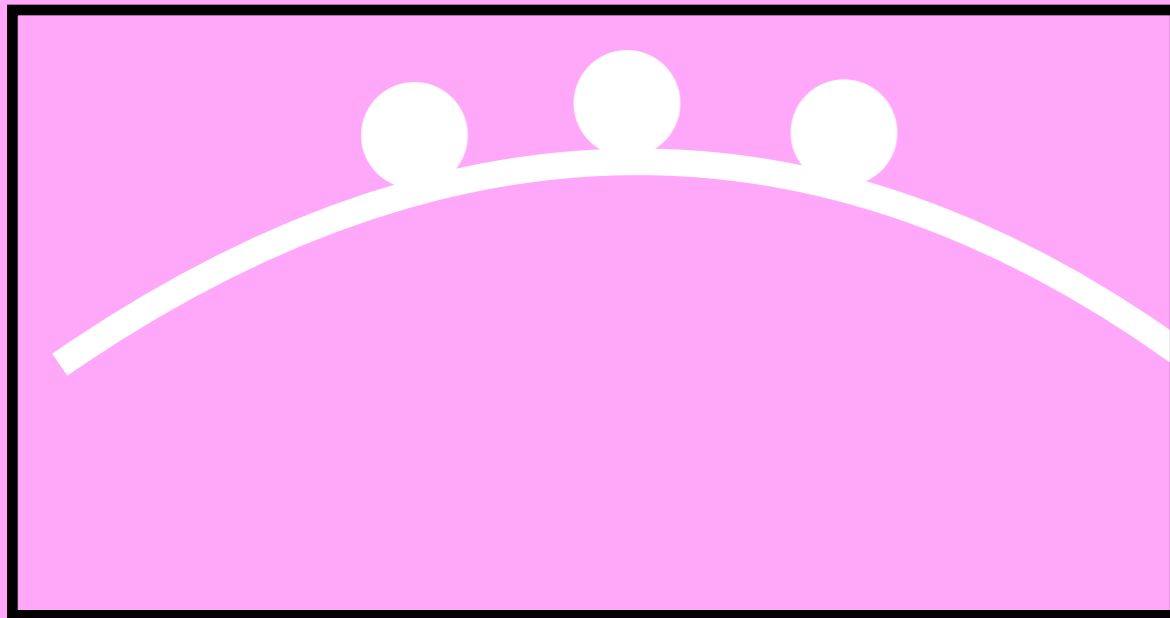
# Experimental set-up



- Control parameters
    - $\theta$ : contact angle of spheres
    - $R$ : radius of spheres
    - $a$ : shaking amplitude
    - $f$ : shaking frequency
    - $h$ : depth of water layer
    - $\phi$ : area fraction
- $\phi = \text{bead area} / \text{total area}$

# Part-1:

- Why the antinode clusters at low  $\phi$  ?

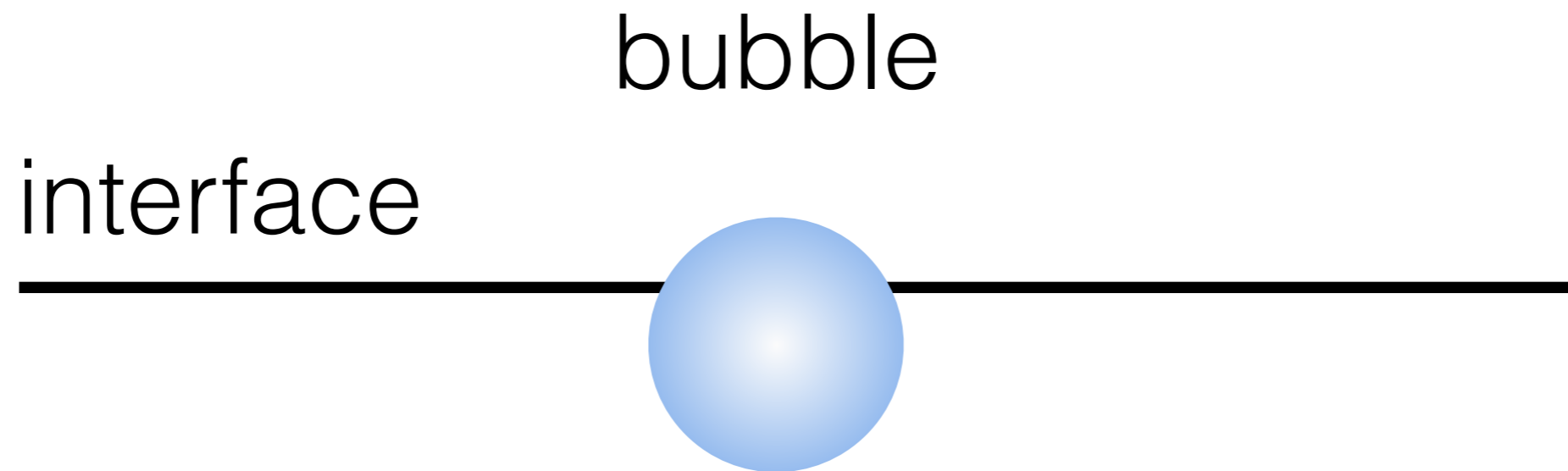


- Why the node clusters at high  $\phi$  ?

# Floater on a static surface

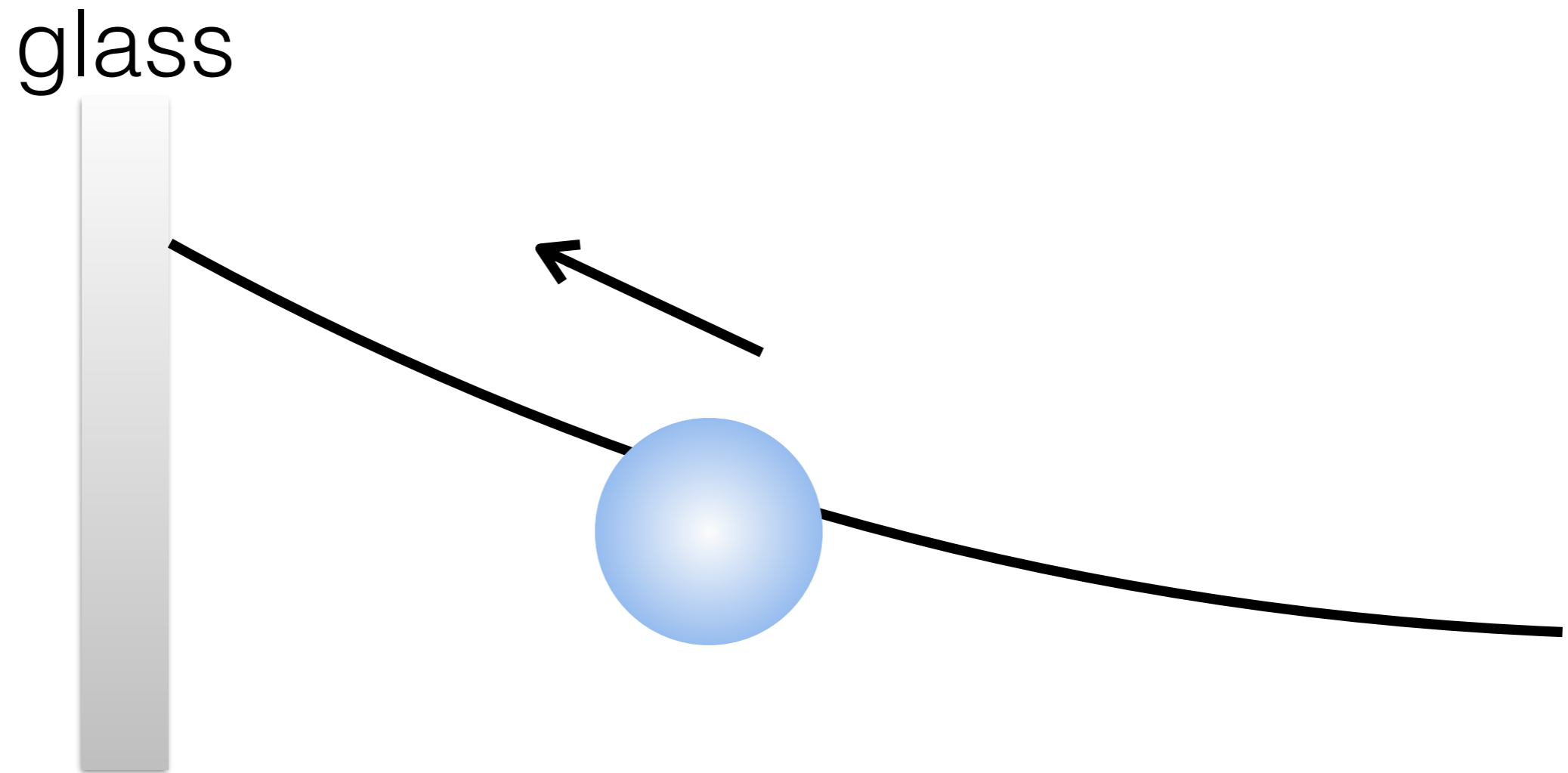


# Bubble in an equilibrium

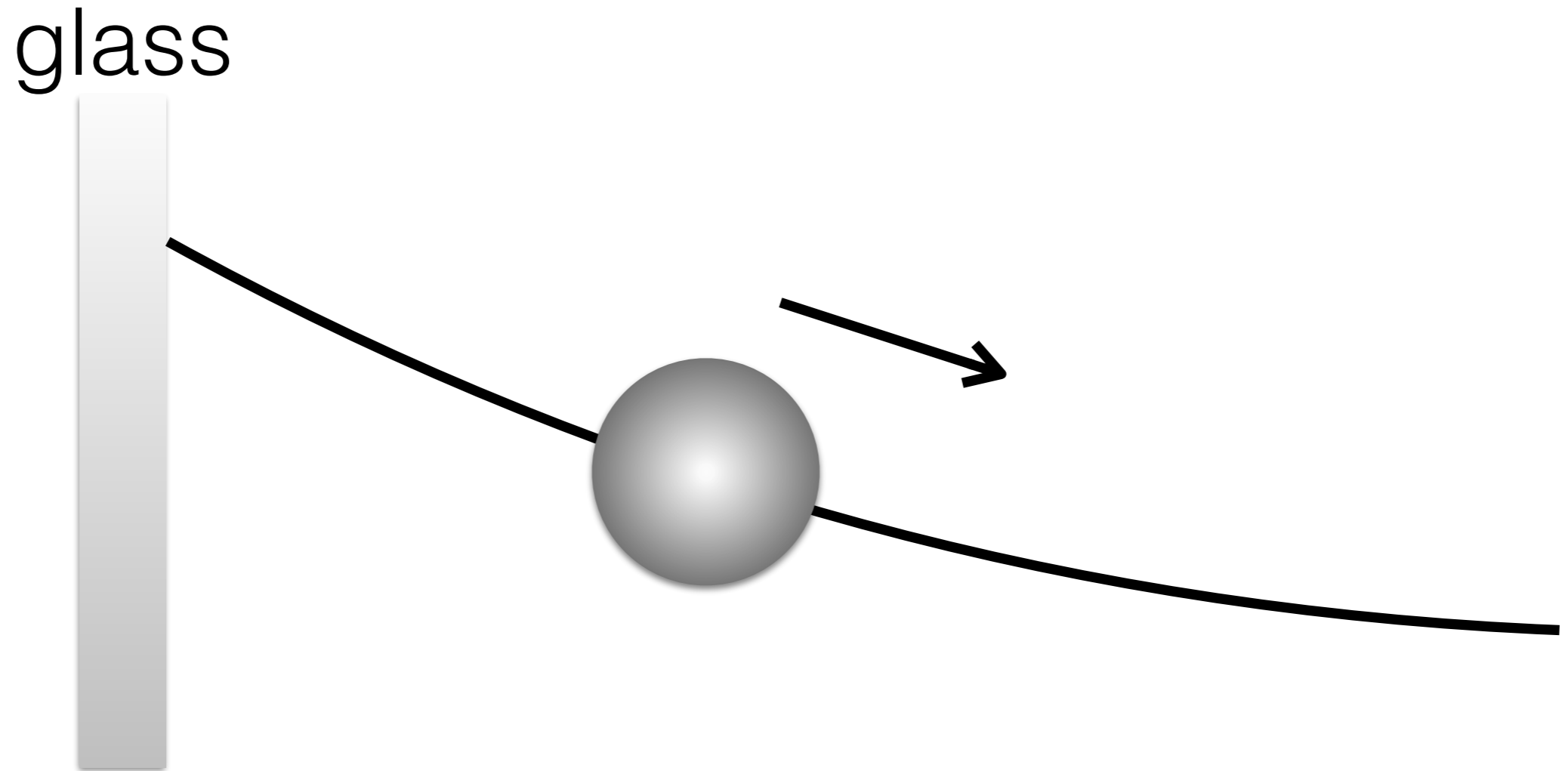




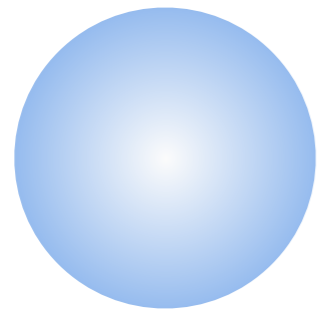
# Bubble in a curved interface



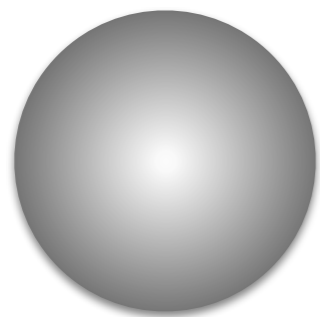
# Heavy particle in a curved interface



# Phenomenological conclusion



- bubbles (**light** particles) drift to a local surface **maximum**

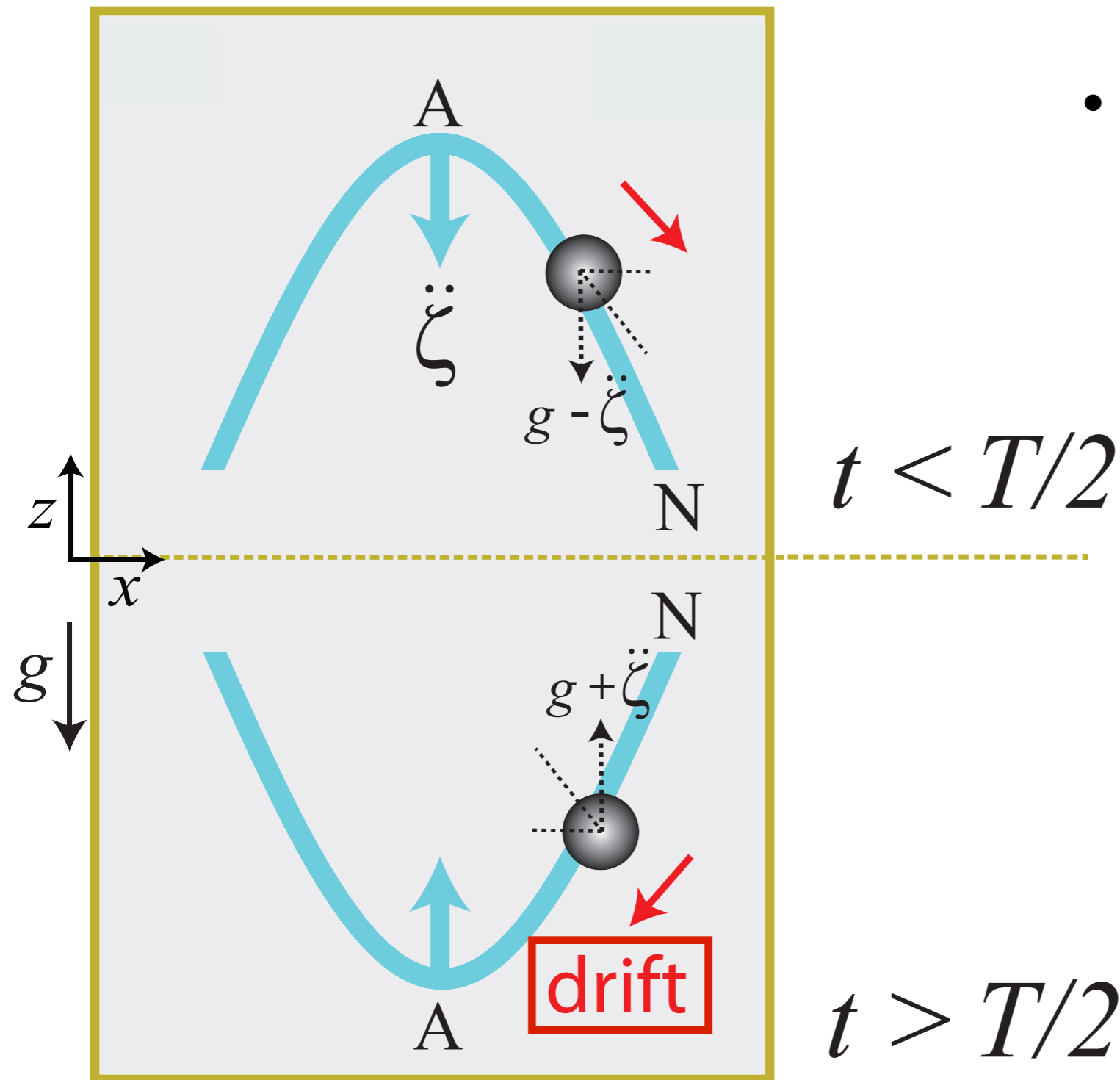


- **heavy** particles drift to a local surface **minimum**

# Floaters on a dynamic surface



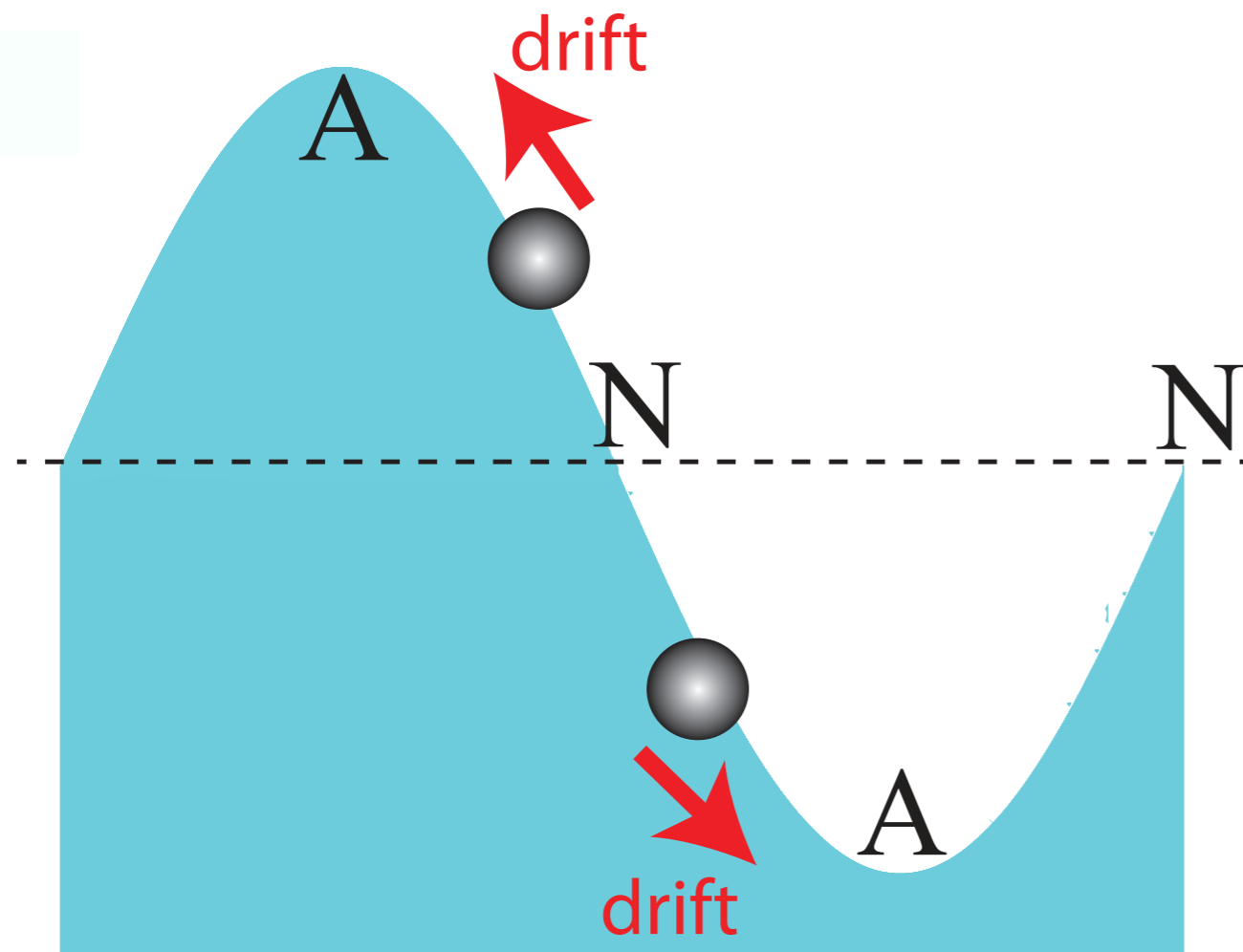
# Heavy particle on a standing wave



- $T$ : the wave period

**wave elevator!**

# Antinode clusters at low $\phi$



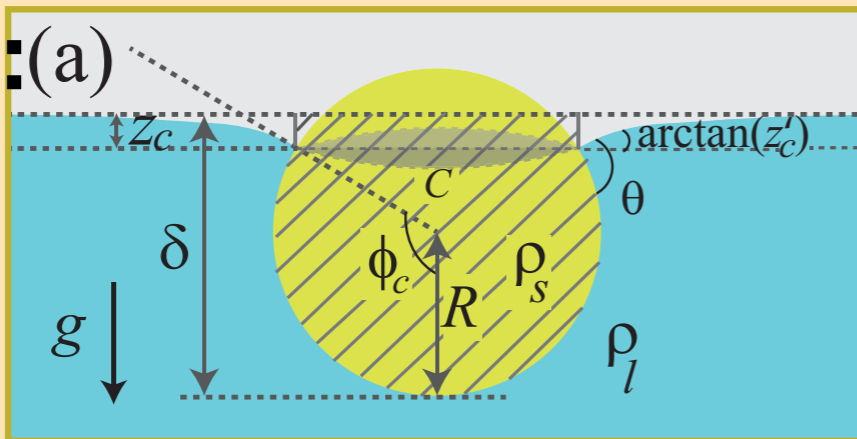
- The drift force is always towards the **antinodes** for our floaters.
- The drift force is a **single** floater force.

- G. Falkovich, A. Weinberg, P. Denissenko, and S. Lukaschuk, “Floater clustering in a standing wave”, *Nature (London)* **435**, 1045–1046 (2005).

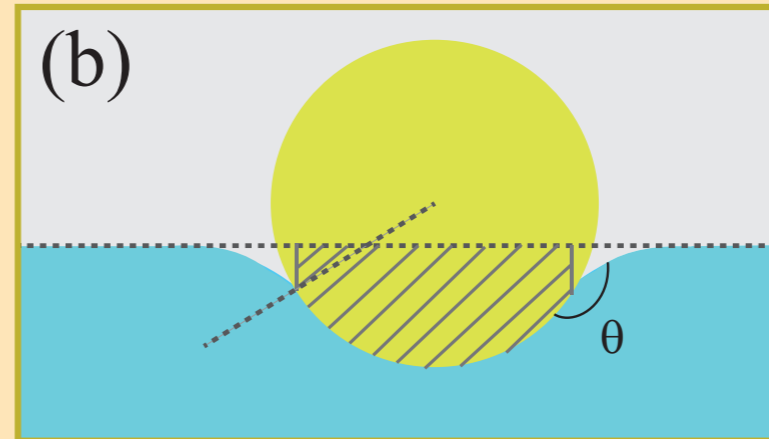
# Full story at low $\phi$

**our floaters:**

hydrophilic ( $\theta < \pi/2$ )



hydrophobic ( $\theta > \pi/2$ )

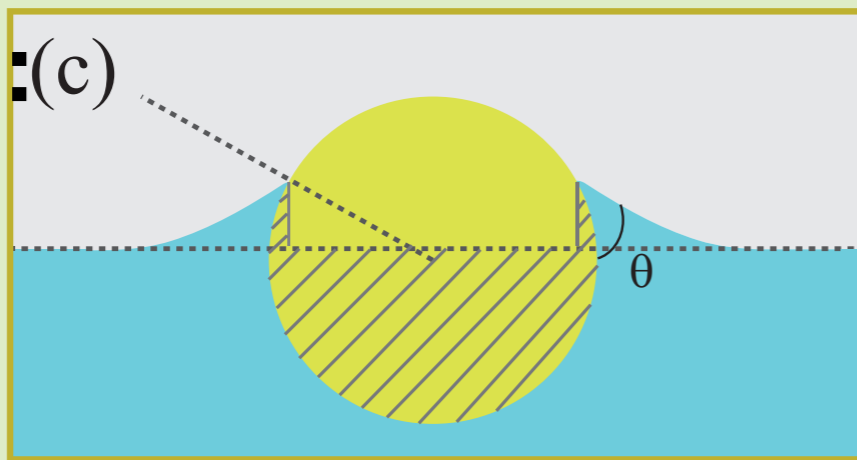


heavy  
( $\rho_s > \rho_l$ )

$\vec{F}_C \uparrow$

**antinode clusters**

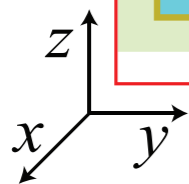
**bubbles:**



light  
( $\rho_s < \rho_l$ )

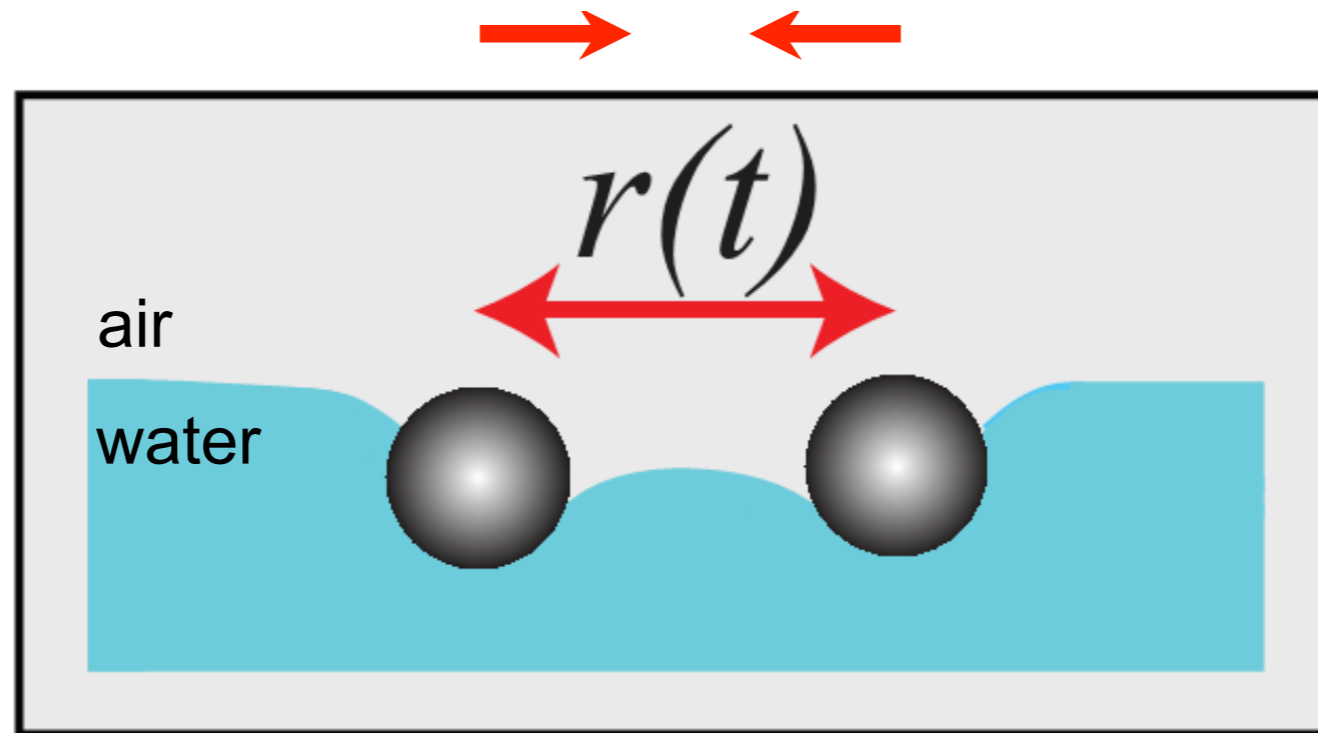
$\vec{F}_C \downarrow$

**node clusters**



- J. B. Keller, "Surface tension force on a partly submerged body", Phys. Fluids **10**, 3009–3010 (1998).

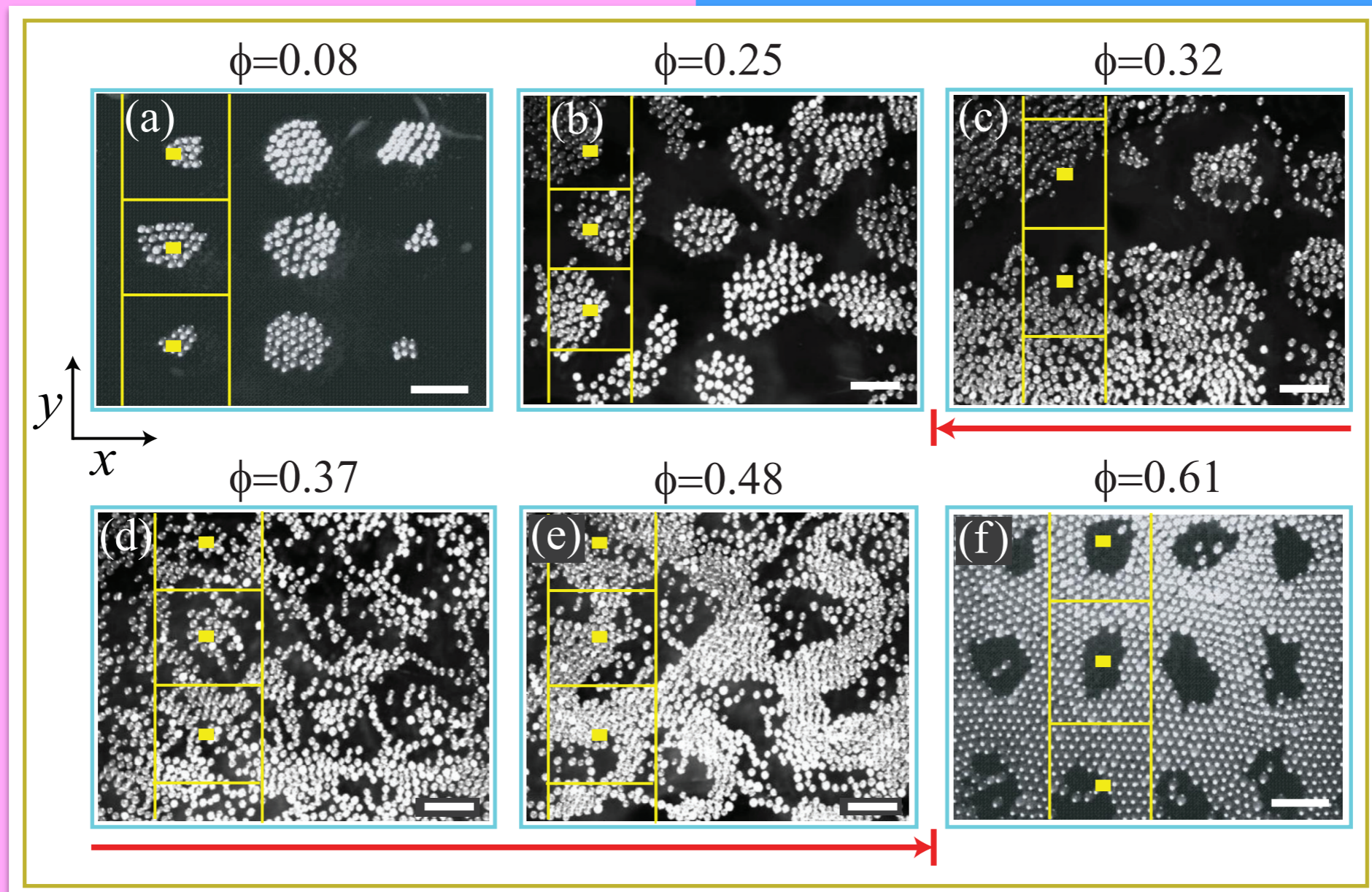
# Attractive capillary interaction



- D. Vella and L. Mahadevan, “The Cheerios effect”, Am. J. Phys. **73**, 817–825 (2005).
- D. Y. C. Chan, J. D. Henry, Jr., and L. R. White, “The interaction of colloidal particles collected at fluid interfaces”, J. Colloid Interface Sci. **79**, 410–418 (1981).
- N. D. Vassileva, D. van den Ende, F. Mugele, and J. Mellema, “Capillary forces between spherical particles floating at a liquid-liquid interface”, Langmuir **21**, 11190–11200 (2005). **(PhD Thesis, UTwente)**



# Experiment: Clusters



arXiv:1405.2027

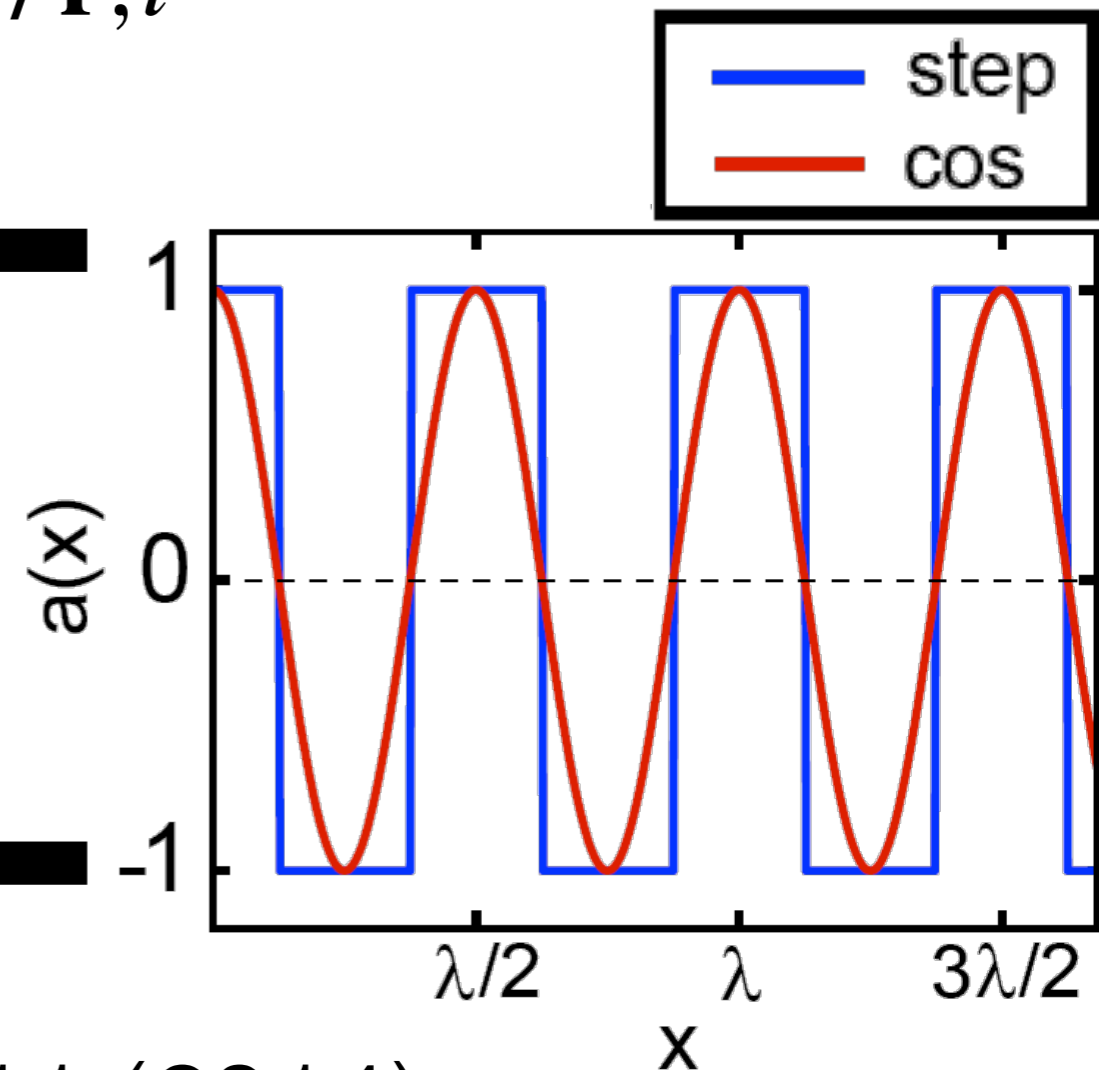
# Correlation number $c$

- $$c \equiv \frac{\langle \phi(\mathbf{r}, t) a(\mathbf{r}) \rangle_{\mathbf{r}, t}}{\langle \phi(\mathbf{r}, t) \rangle_{\mathbf{r}, t}}$$

- $\mathbf{r} = (x, y)$
- $t$  is time
- $\lambda$  is the wavelength

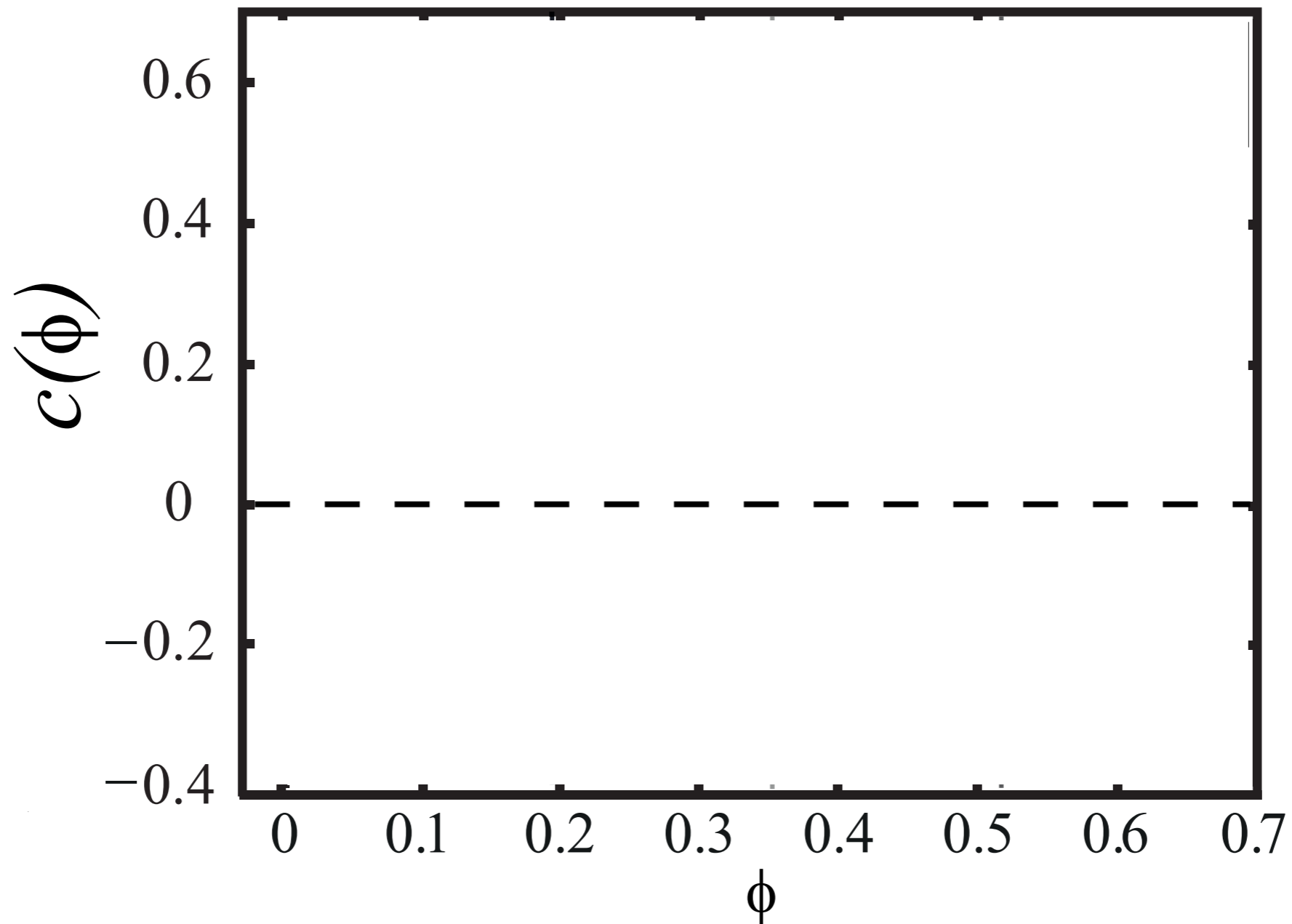
antinodes ←

nodes ←

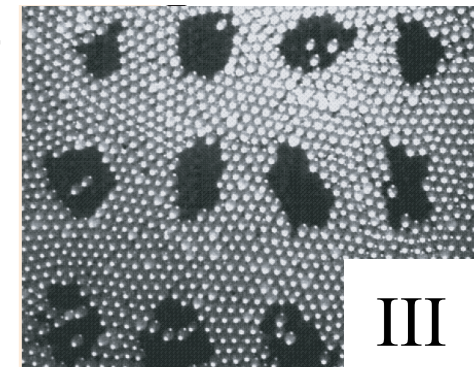
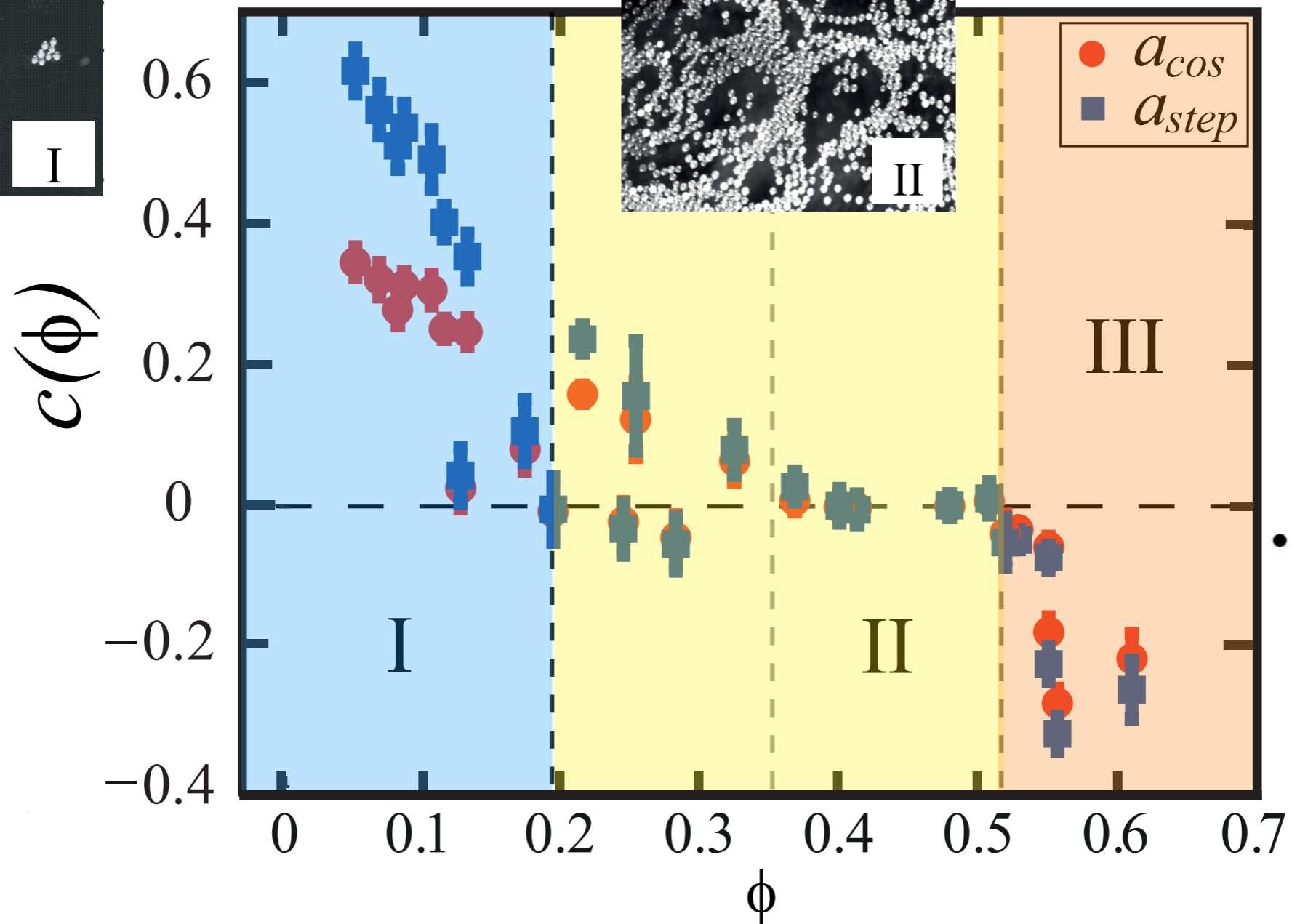
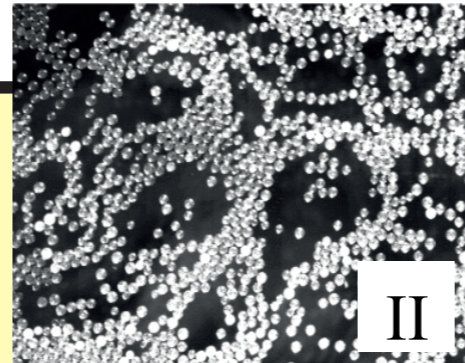
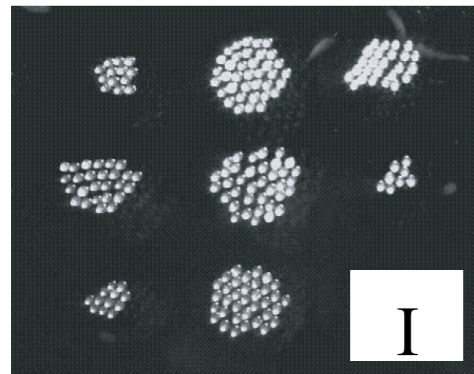


- C. Sanlı *et al.* PRE **89**, 053011 (2014).

# Correlation number $c$



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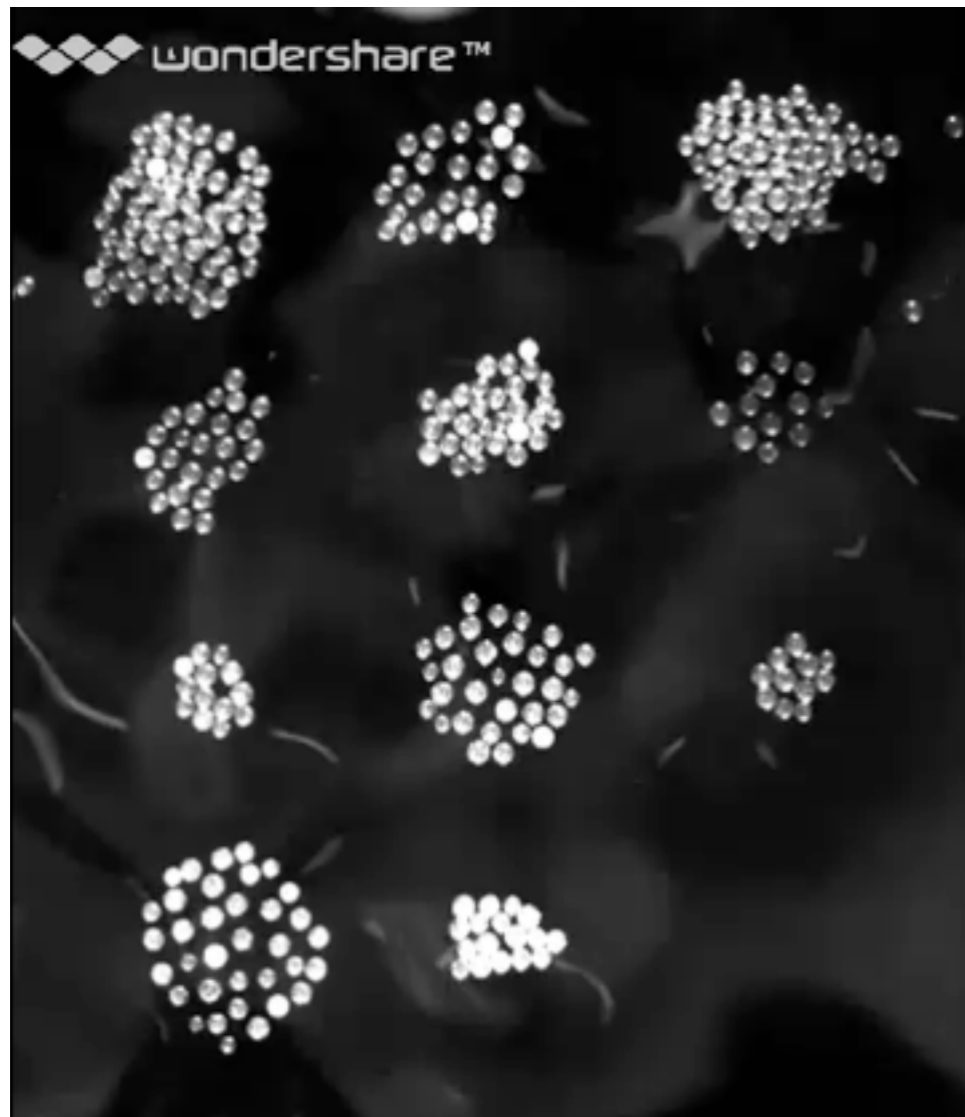


# Why the node clusters at high $\phi$ ?

- Let's look at the experiment more carefully!

**breathing**

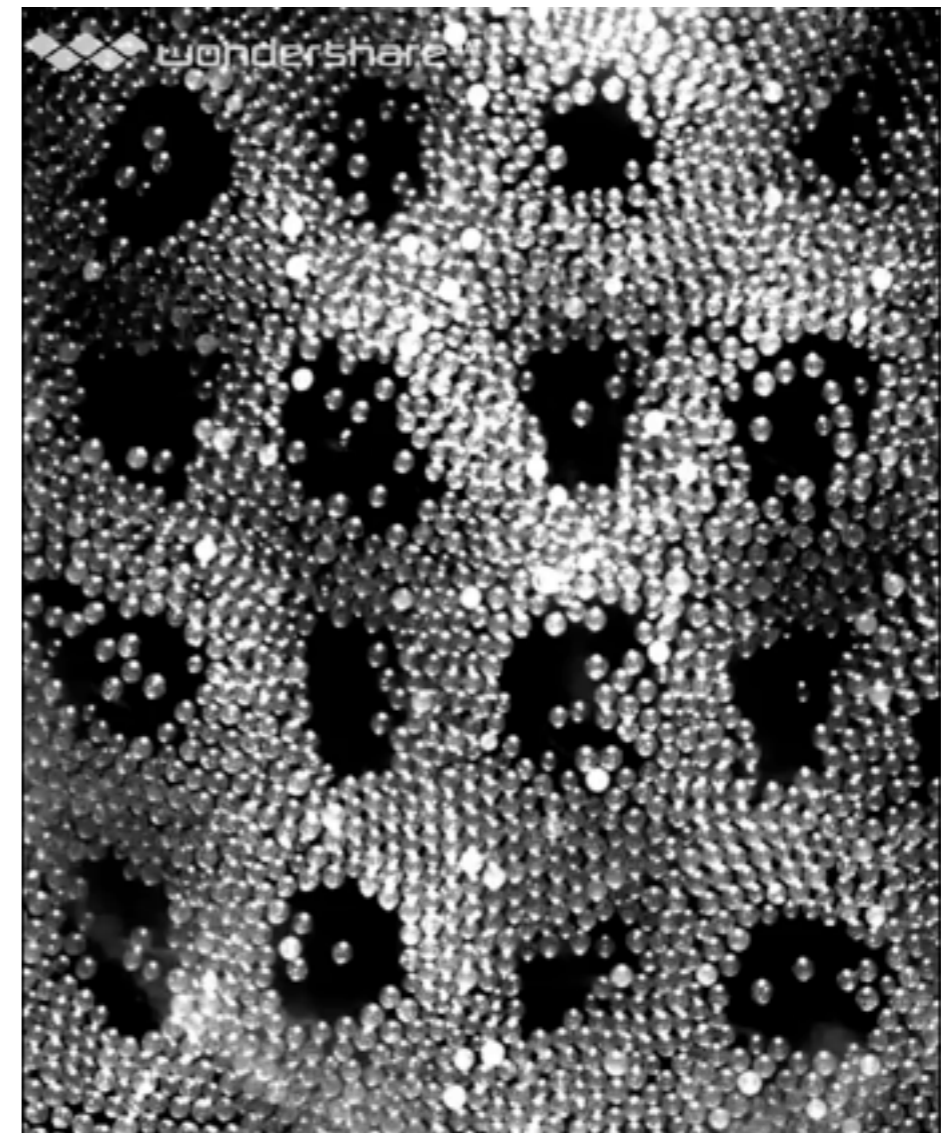
10 mm



**antinode clusters**

**non-breathing**

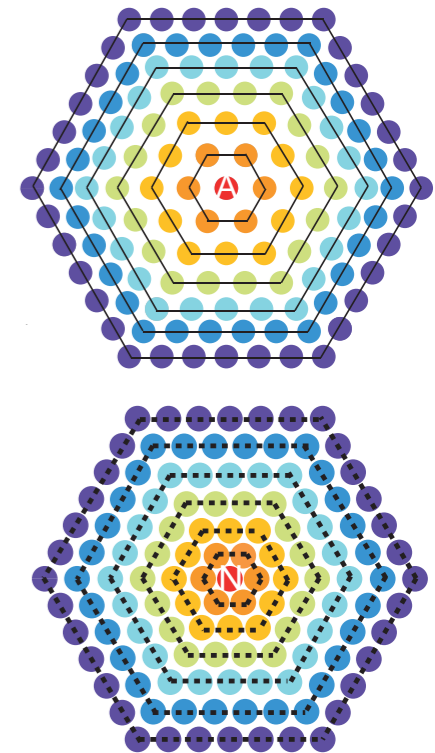
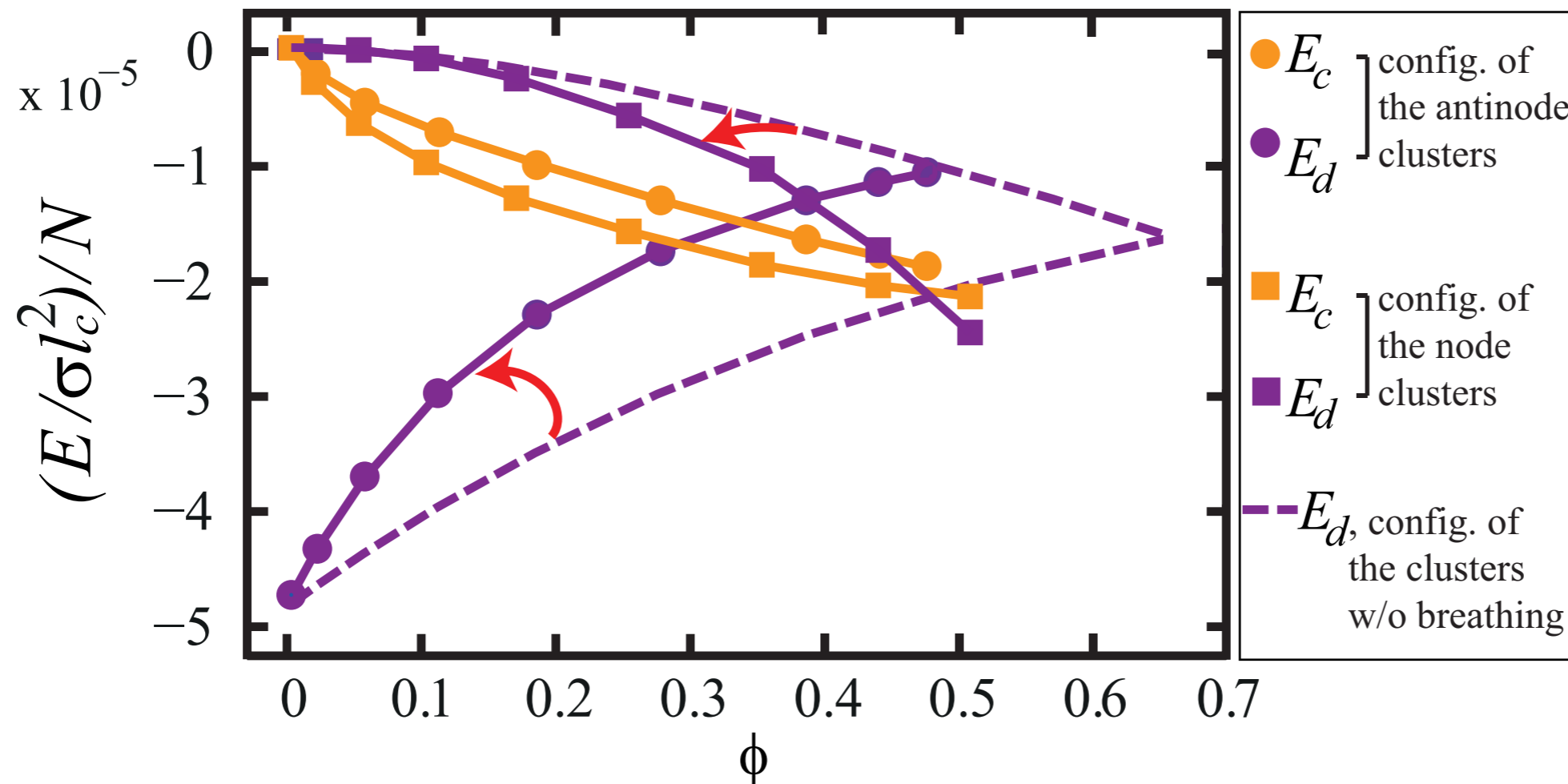
10 mm



**node clusters**

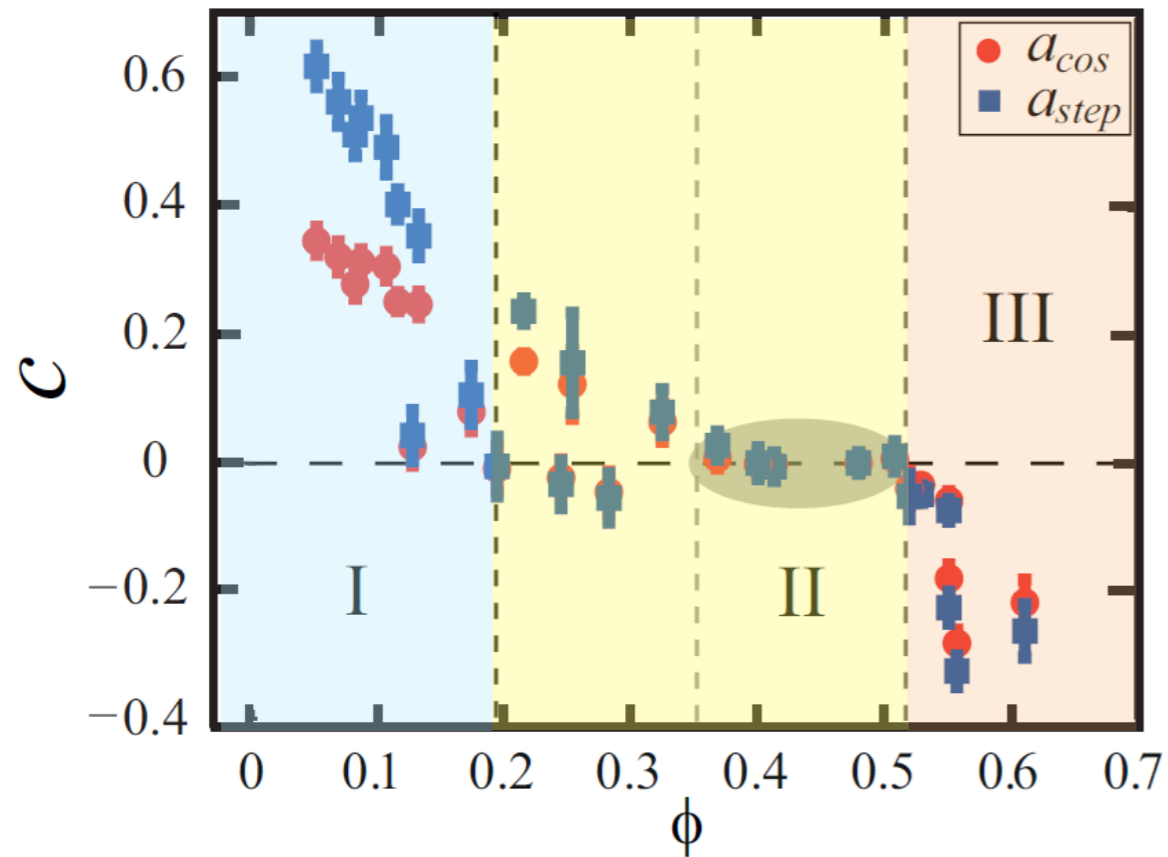
# Potential energy approach

## drift to the antinodes + capillary attraction + breathing



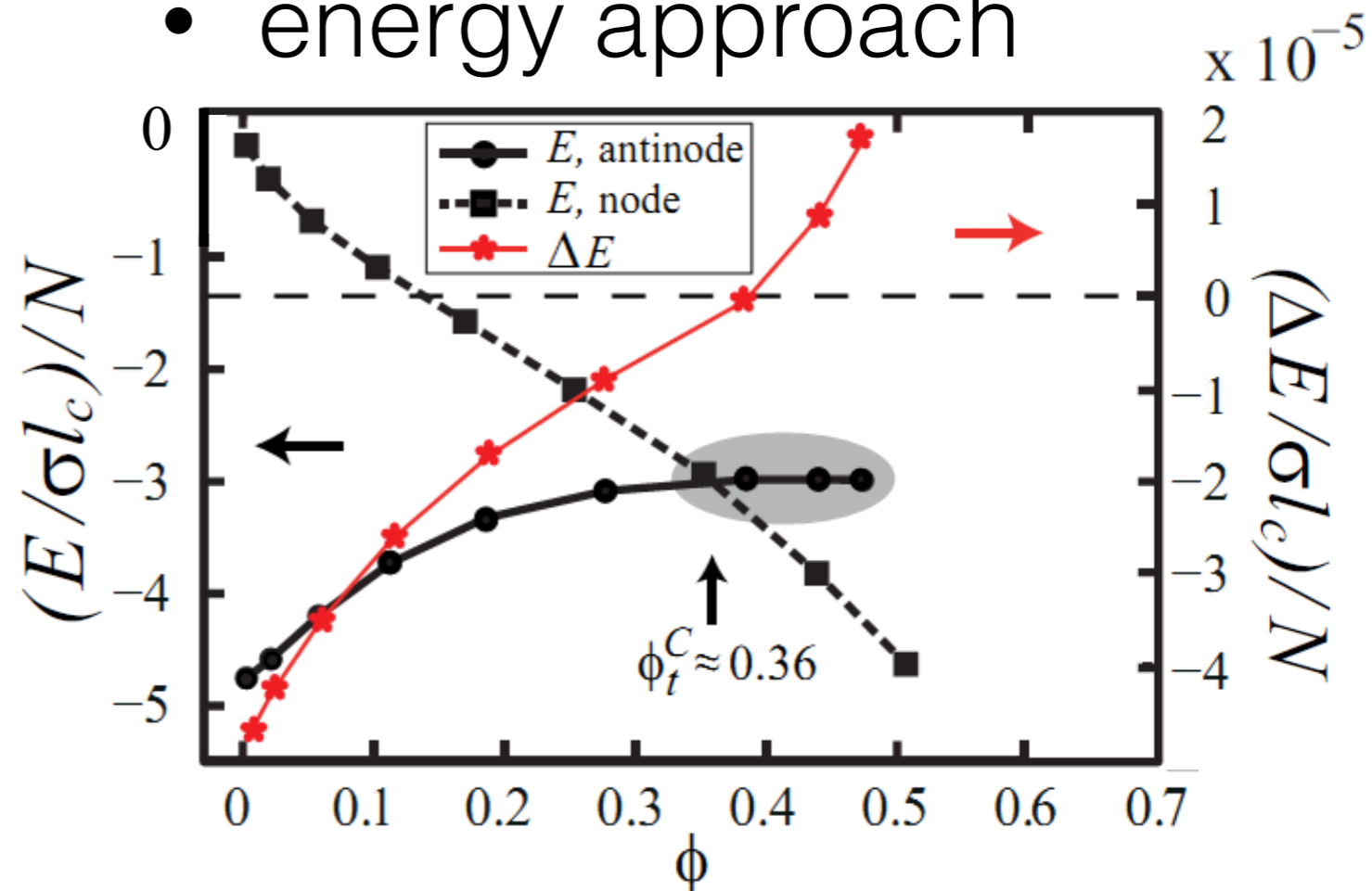
- $E_c$ : **capillary energy**
- $E_d$ : **drift energy**
- $\sigma$ : surface tension
- $l_c$ : capillary length
- $N$ : number of floaters

# Comparison with experiment



- experiment

- energy approach



# Part- 2:

Floater on a

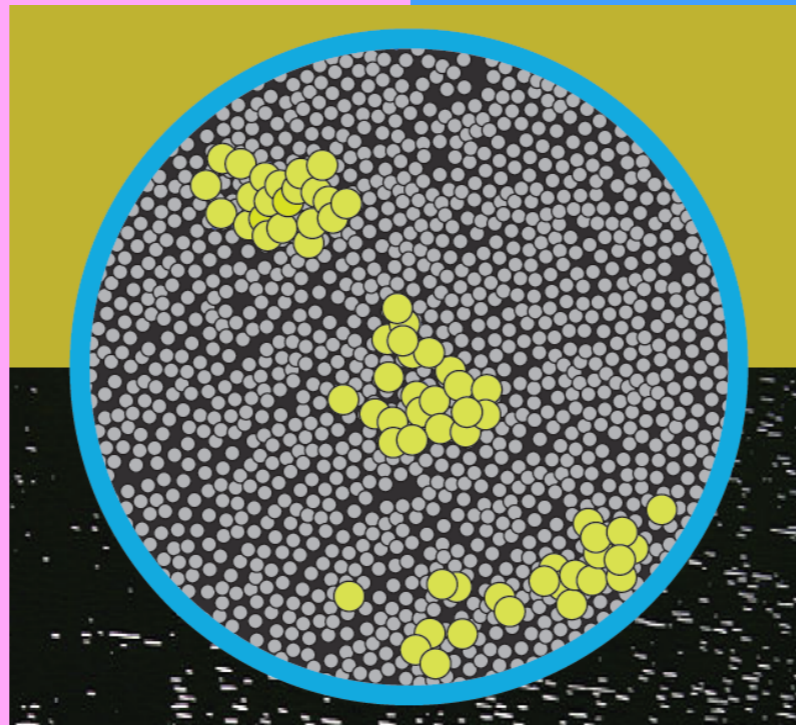
Faraday wave:

capillary

Heterogeneous

flows and

group formations





# Observation

- $\phi = 0.63$

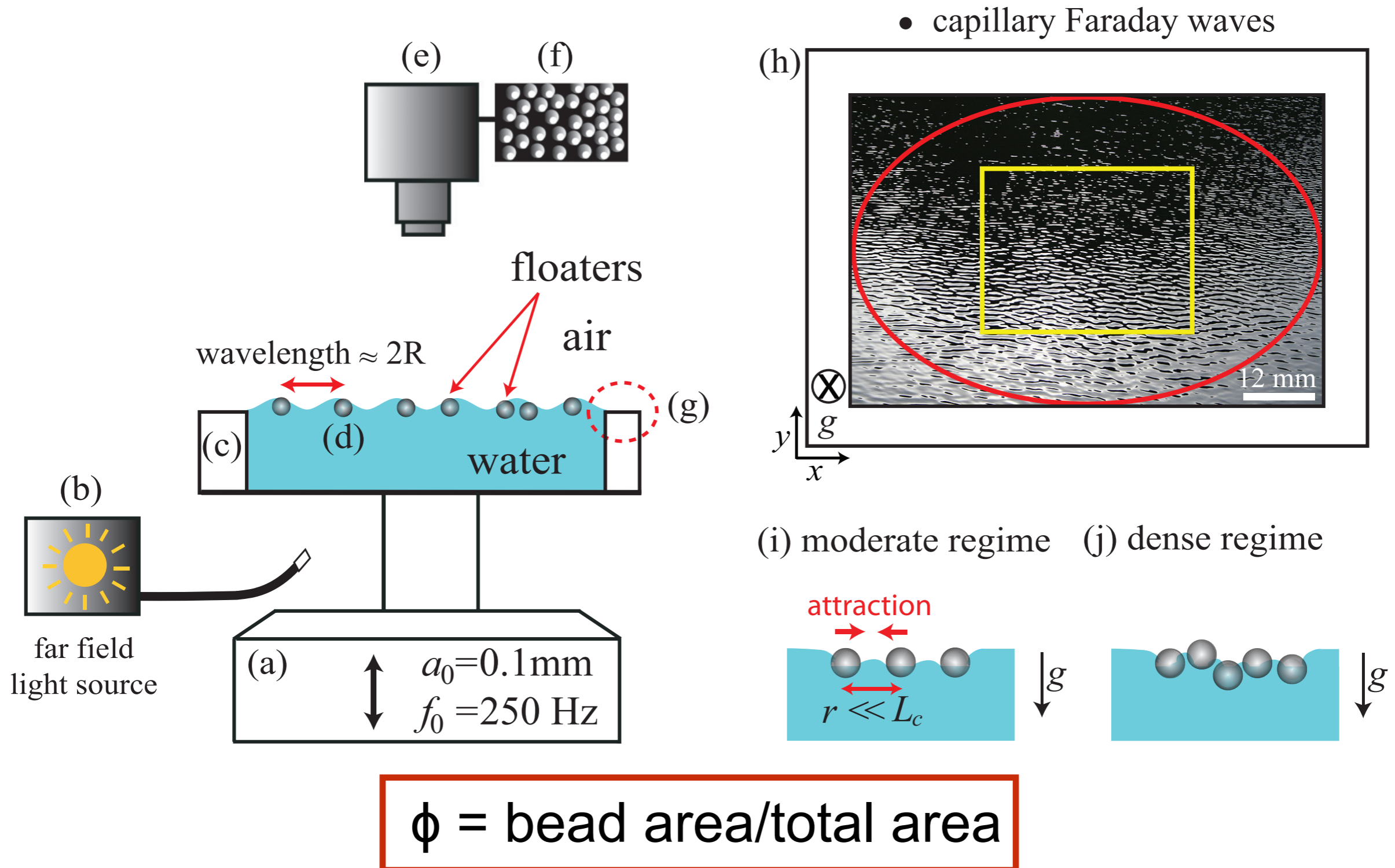
- $a = 1 \text{ mm}$
- $f = 250 \text{ Hz}$

- 4 times slower than real time



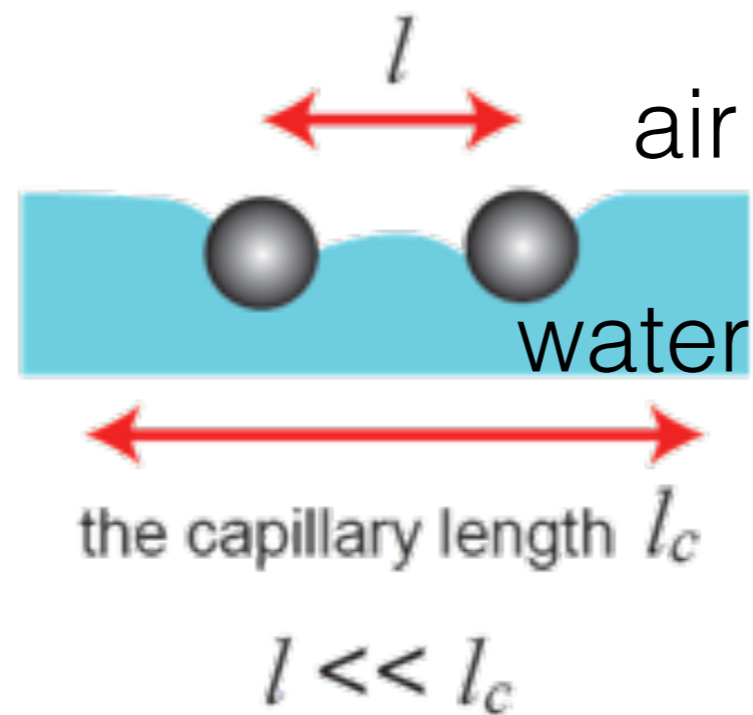
2 mm

# Experimental set-up

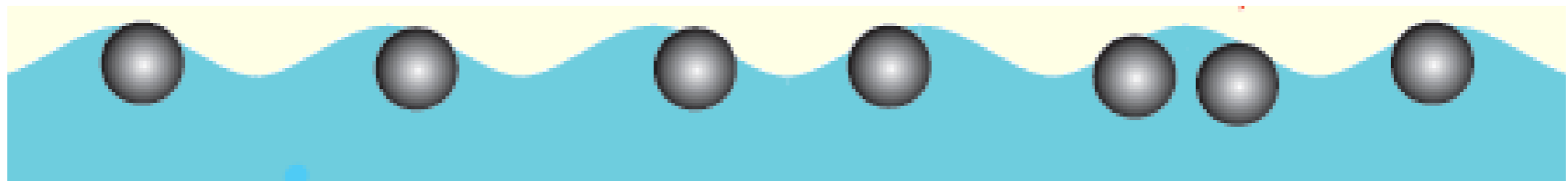


# Driving interactions and forces

- attractive capillary interactions

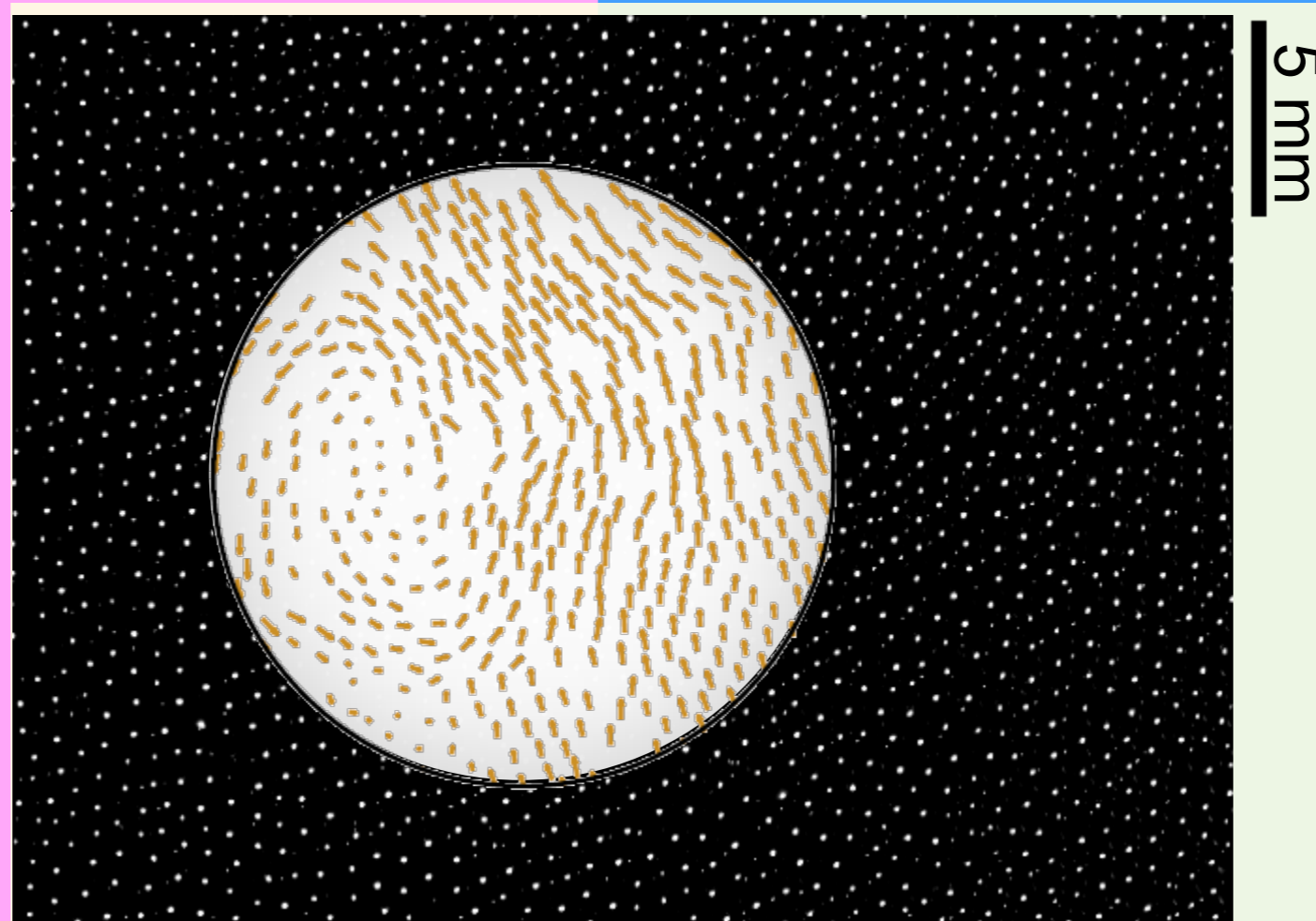


- erratic capillary (Faraday) waves



Four-point  
dynamic  
susceptibility:

To quantify the  
heterogeneous  
dynamics



# Self-overlap order parameter

time  $\leftarrow$

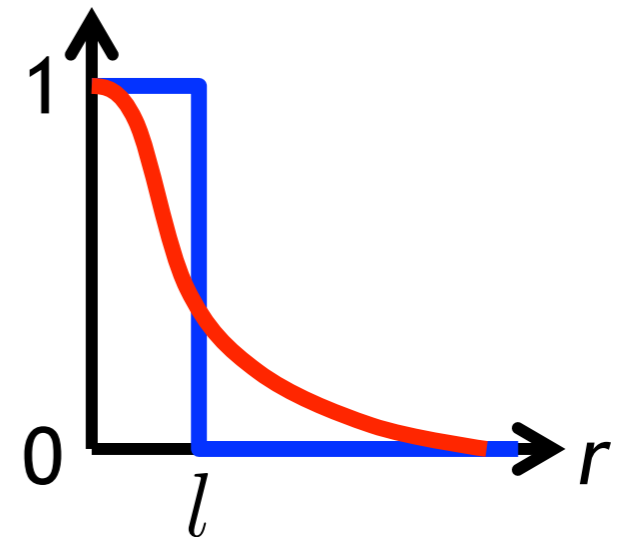
$$Q_t(l, \tau) = \frac{1}{N} \sum_{i=1}^N \omega_l(r_i)$$

number of spheres  $\leftarrow$

$\rightarrow$  a cutoff function

$\rightarrow$  sphere index

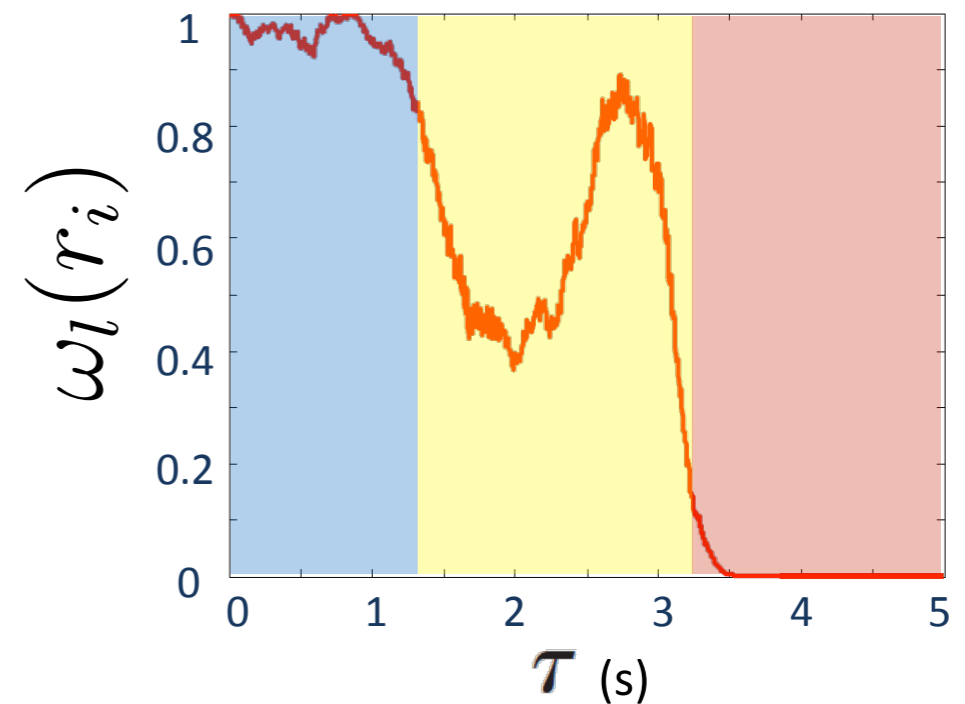
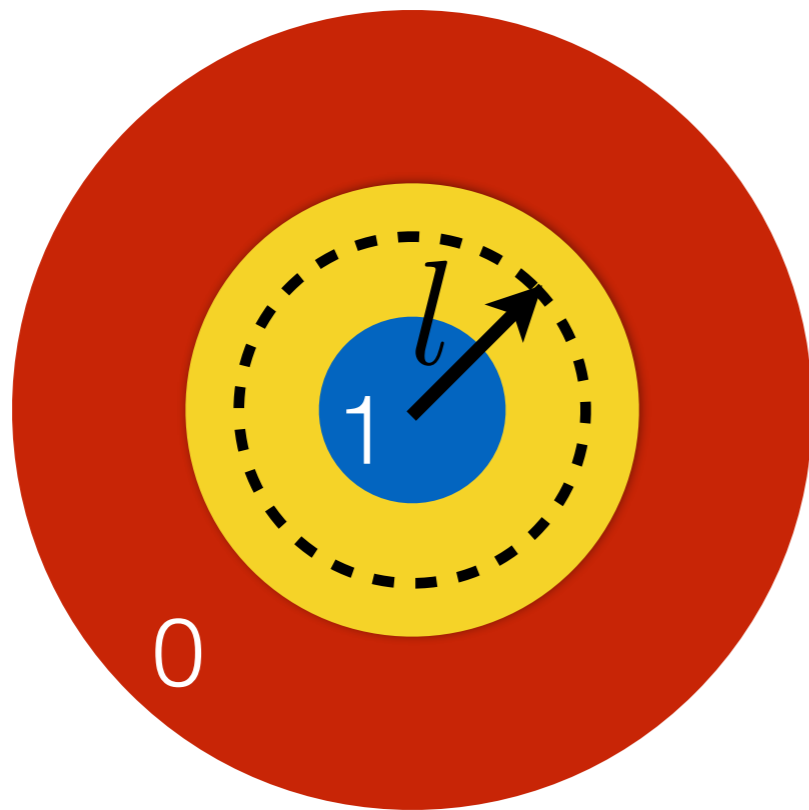
- $\omega_l(r_i) = e^{-r_i^2 / (2l)^2}$
- $r_i = |\mathbf{r}_i(t + \tau) - \mathbf{r}_i(t)|$



- A. R. Abate and D. J. Durian, "Topological persistence and dynamical heterogeneities near jamming", Phys. Rev. E **76**, 021306-1-9 (2007).

# A measure of mobility

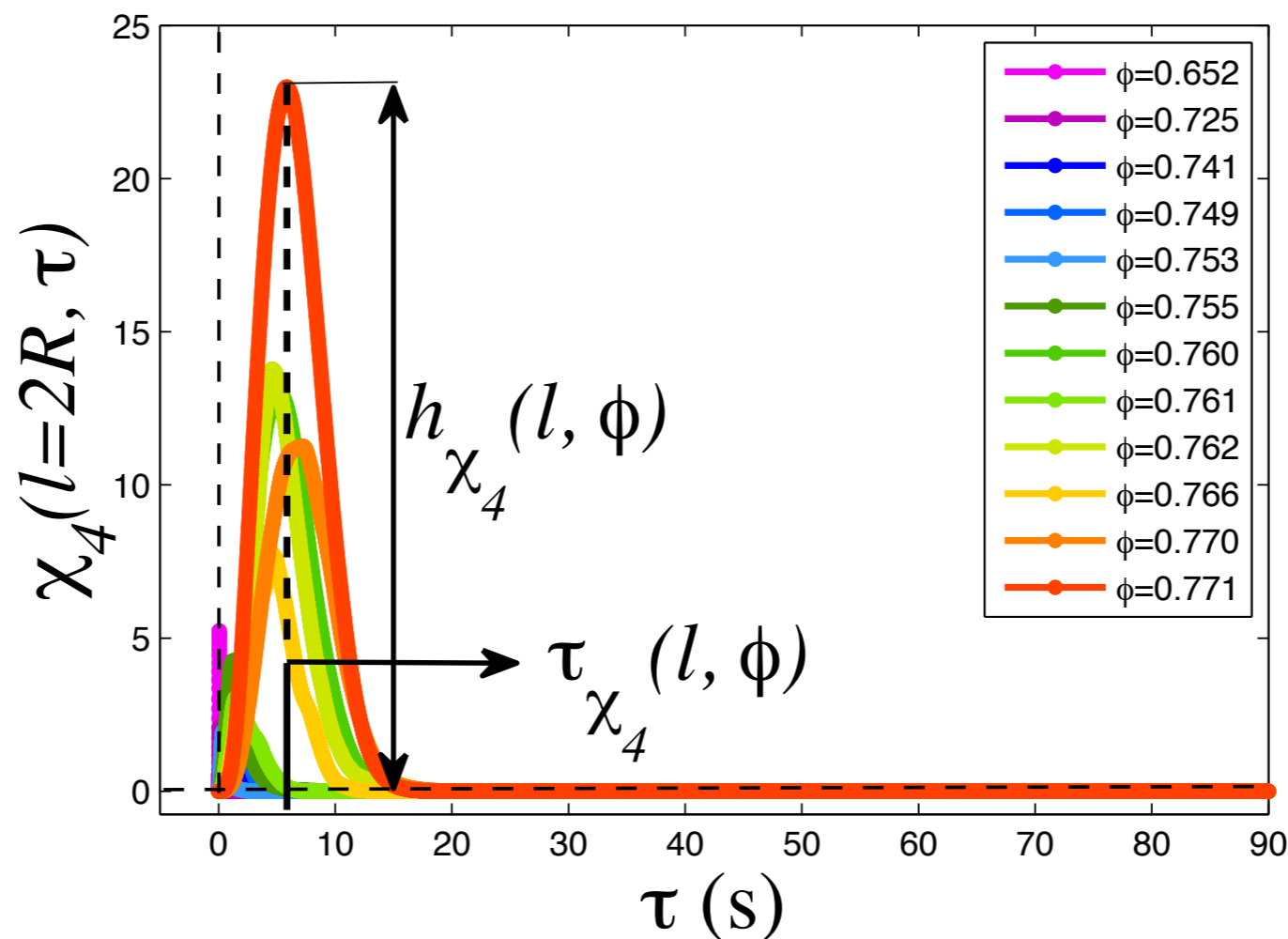
- $\omega_l(r_i) = \begin{cases} 1, & \text{if a sphere } i \text{ travels } \ll l \text{ within } \mathcal{T} \\ 0, & \text{if a sphere } i \text{ travels } \gg l \text{ within } \mathcal{T} \end{cases}$



# Four-point dynamic susceptibility

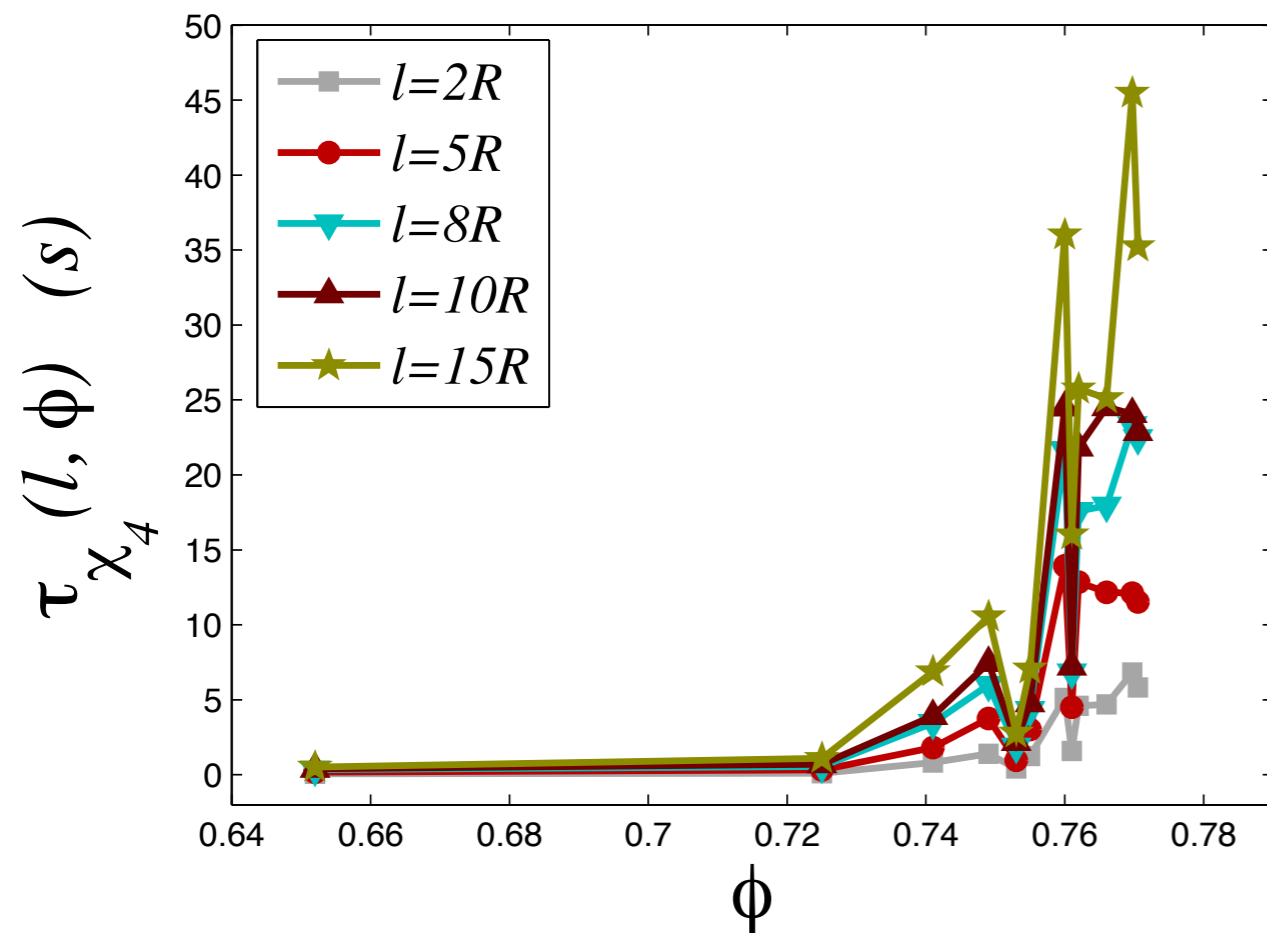
- To quantify the heterogeneous dynamics by the mobility (self-overlap order parameter):

- $$\chi_4(l, \tau) = N \left[ \langle Q_t^2(l, \tau) \rangle_t - \langle Q_t(l, \tau) \rangle_t^2 \right]$$

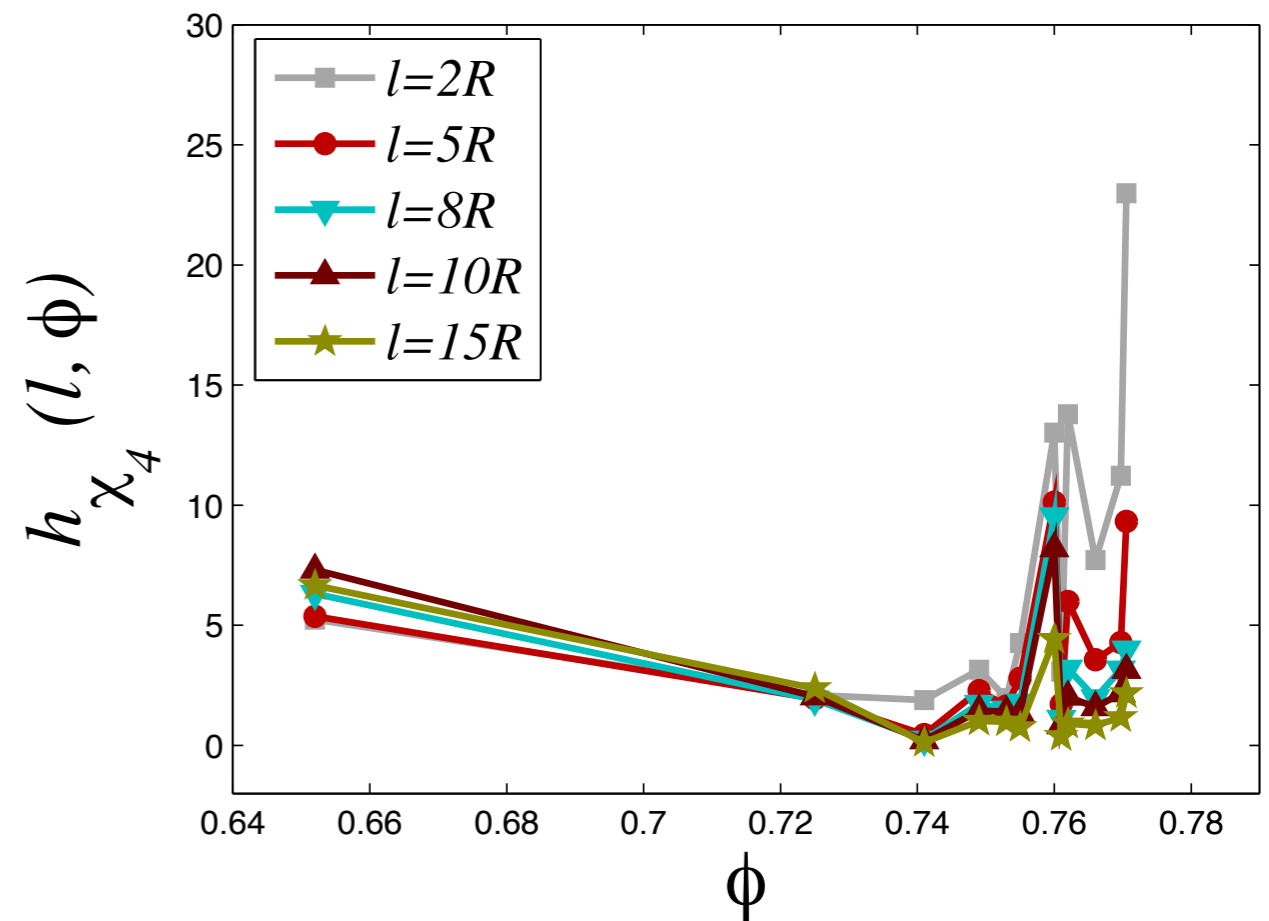


# Quantifying the heterogeneous flow

- time-scale of the heterogeneity



- amount of the heterogeneity



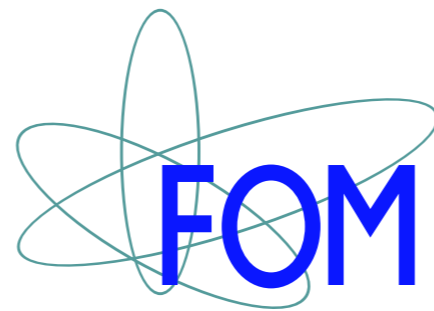


# Greetings, ...

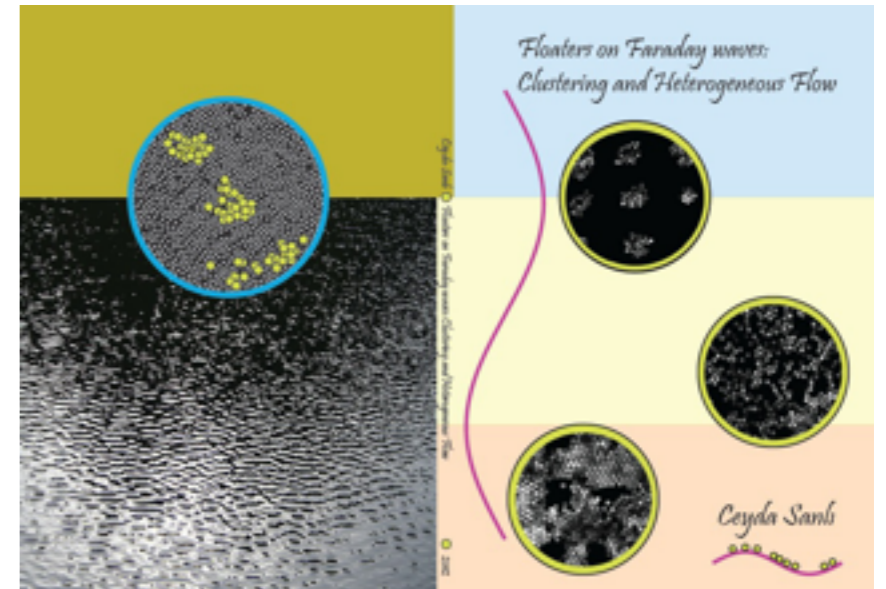


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**UNIVERSITY OF TWENTE.**



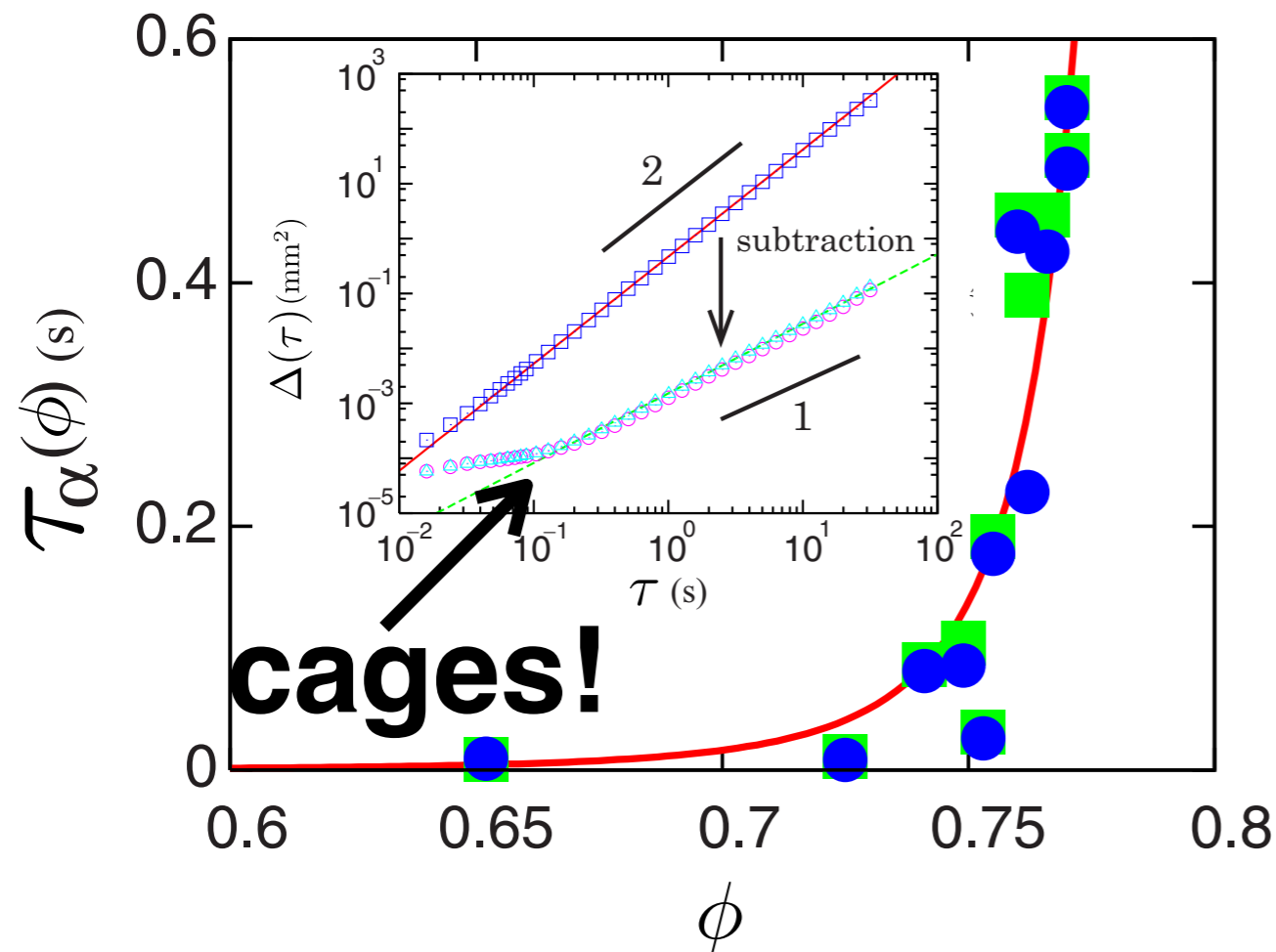
**cedaysan@gmail.com**

**s.luding  
@utwente.nl**

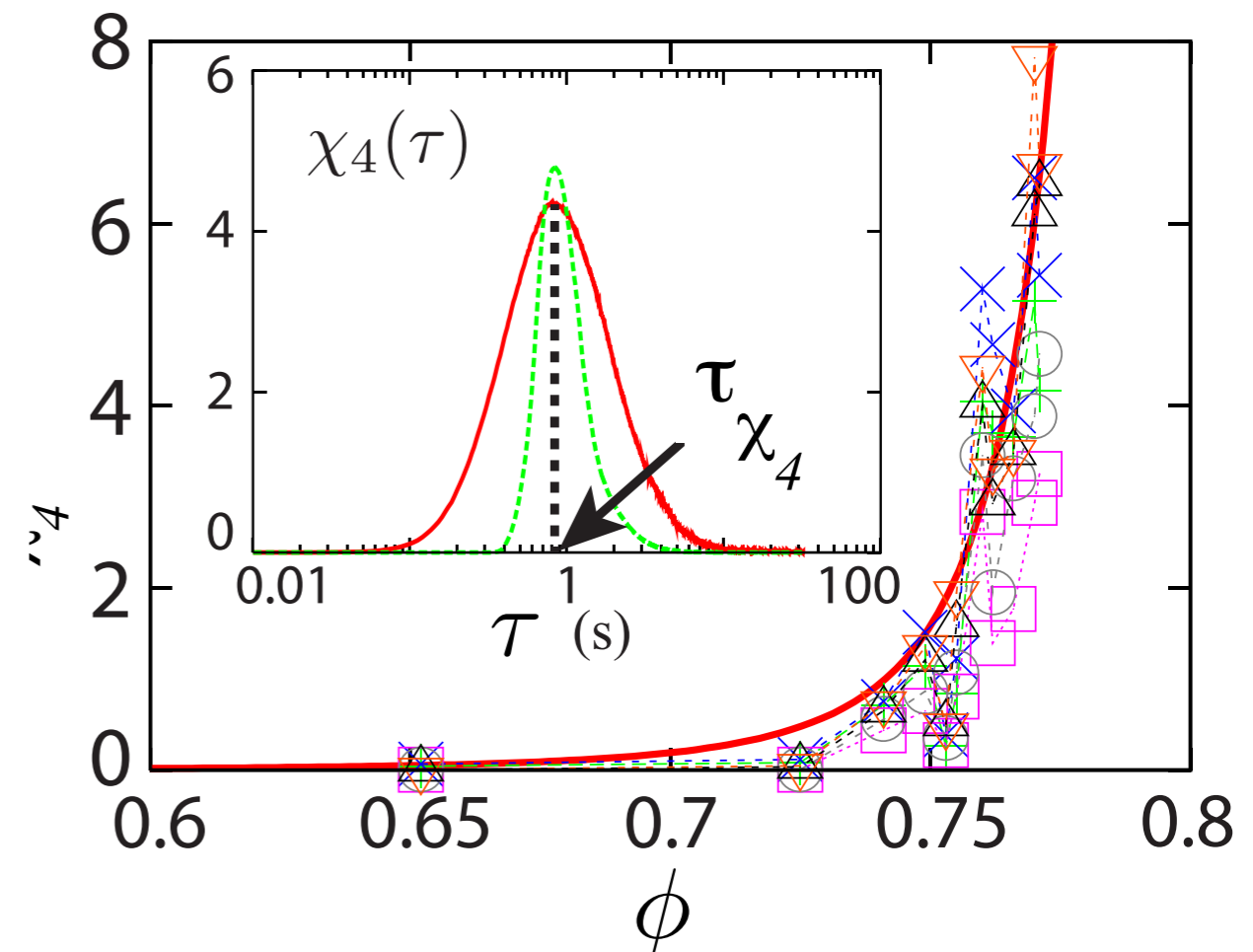


# What does $\chi_4(l, \tau)$ measure?

- time-scale of the flow by the diffusion



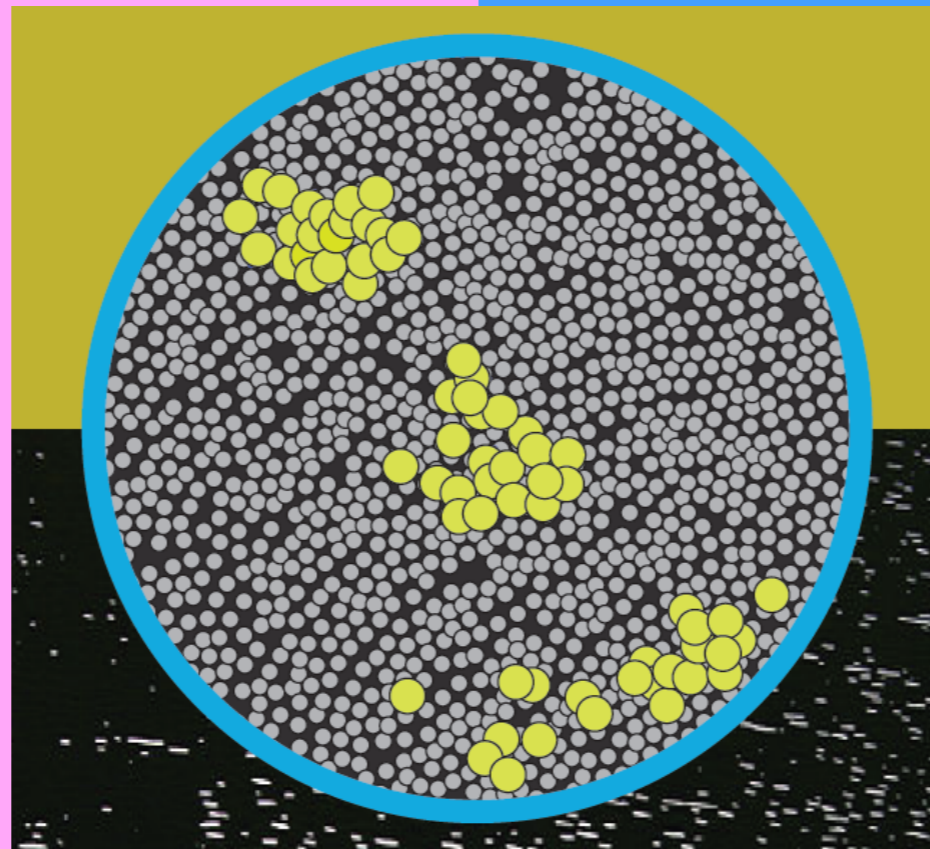
- time-scale of the heterogeneous flow



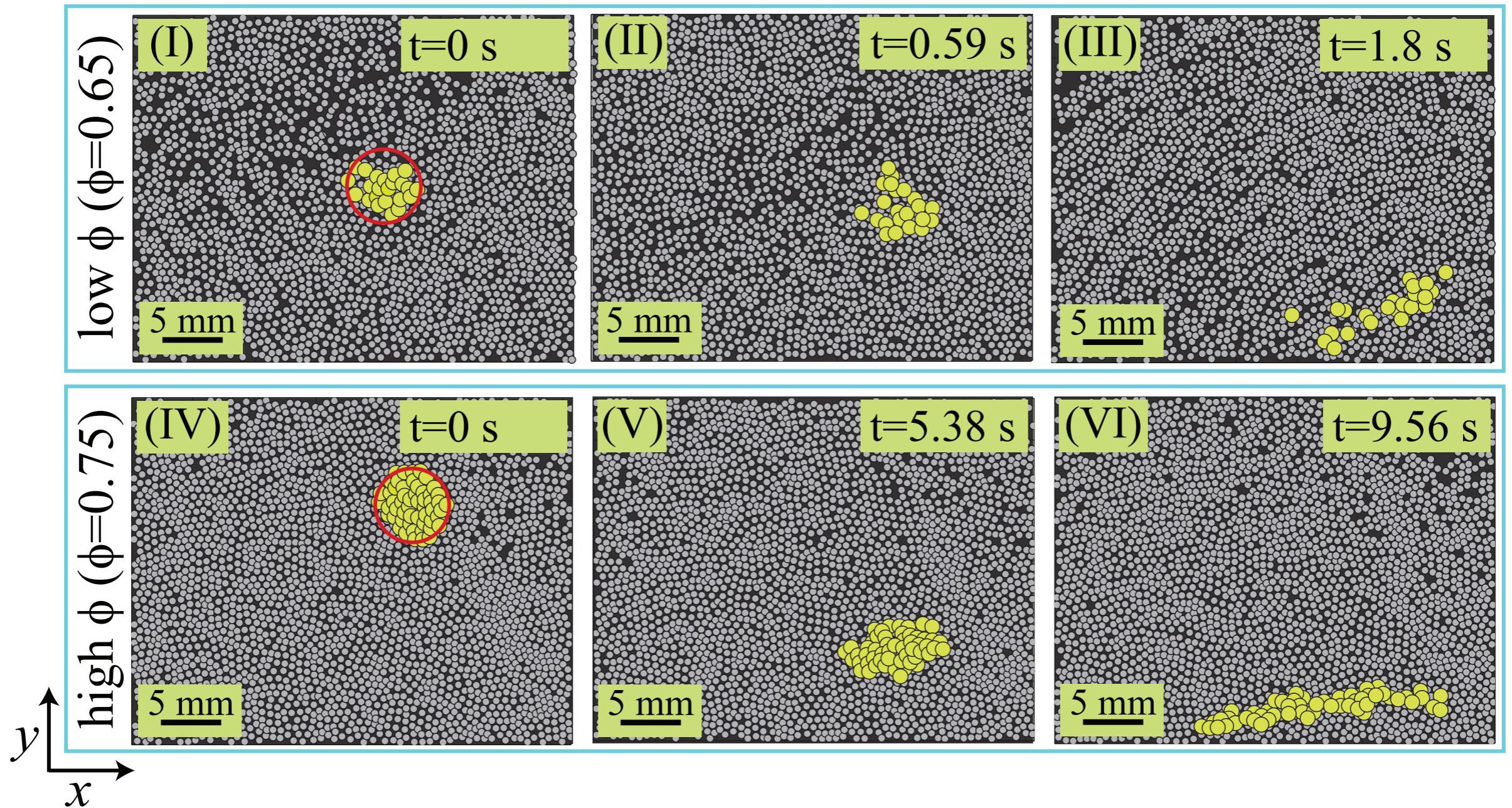
- C. Sanlı, Kuniyasu Saitoh *et al.* PRE **90**, 033018 (2014).

Morphological  
analysis:

Quantifying  
dynamics of the  
group deformation

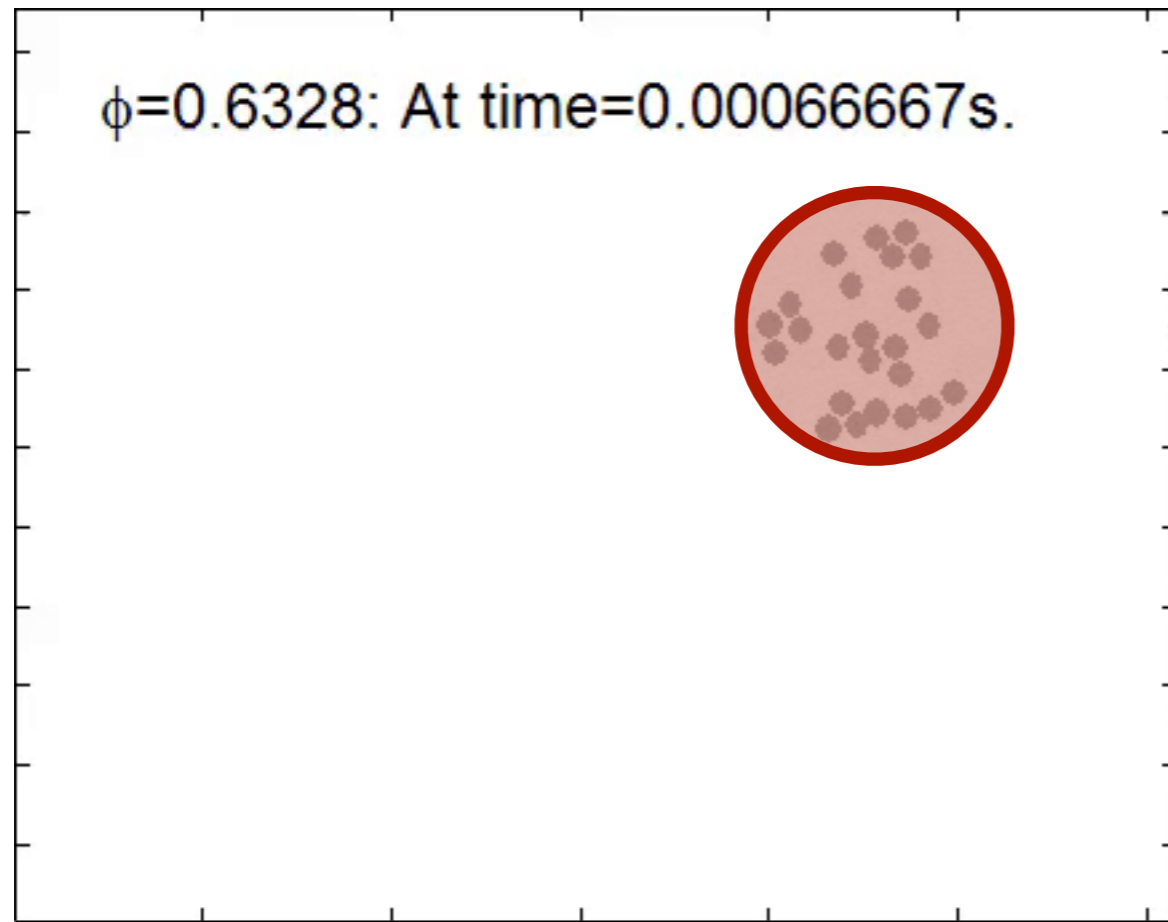


# Observation: Detail looking



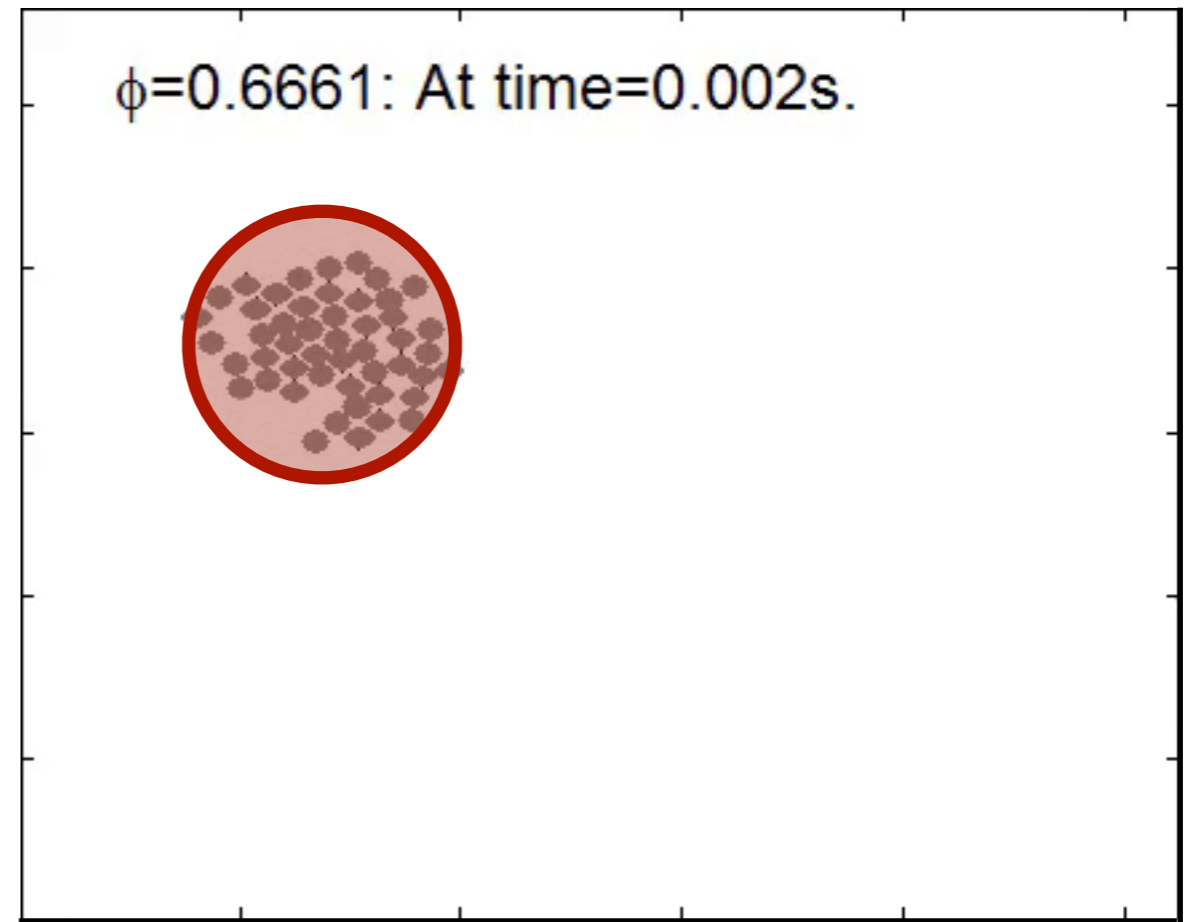
# Construction of a subgroup

low  $\phi$



5 mm

high  $\phi$

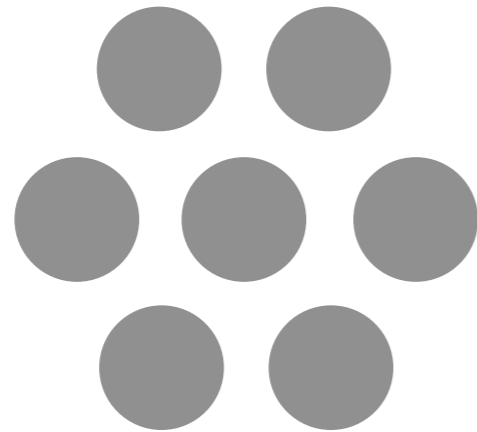


5 mm

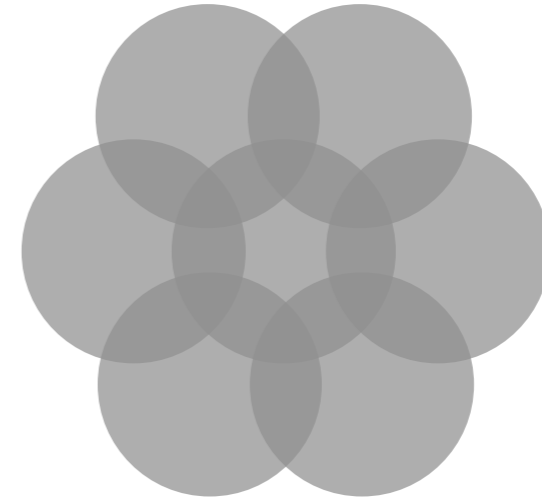
- At a certain  $\phi$ , we construct a subgroup of beads which
- For high  $\phi$ , the subgroup deforms at a later time. the initial positions are inside the shaded region.
- For low  $\phi$ , the subgroup breaks into more pieces.
- Then, we track these beads as a function of time.

# Morphological parameter when $r=2R$

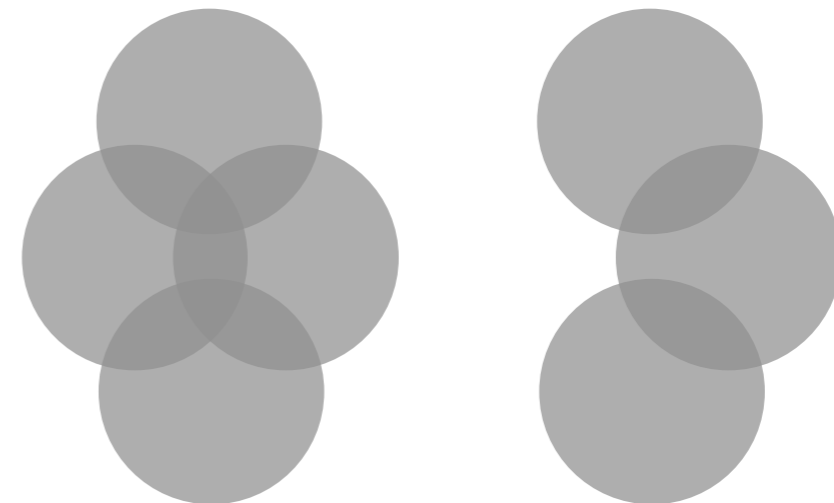
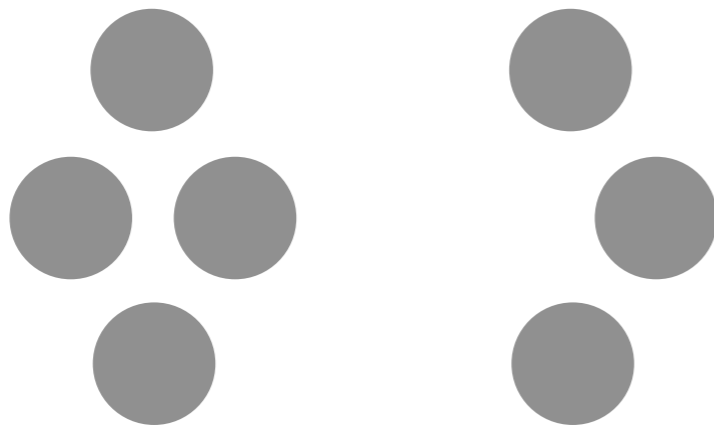
- When  $r=1R$ :



- When  $r=2R$ :



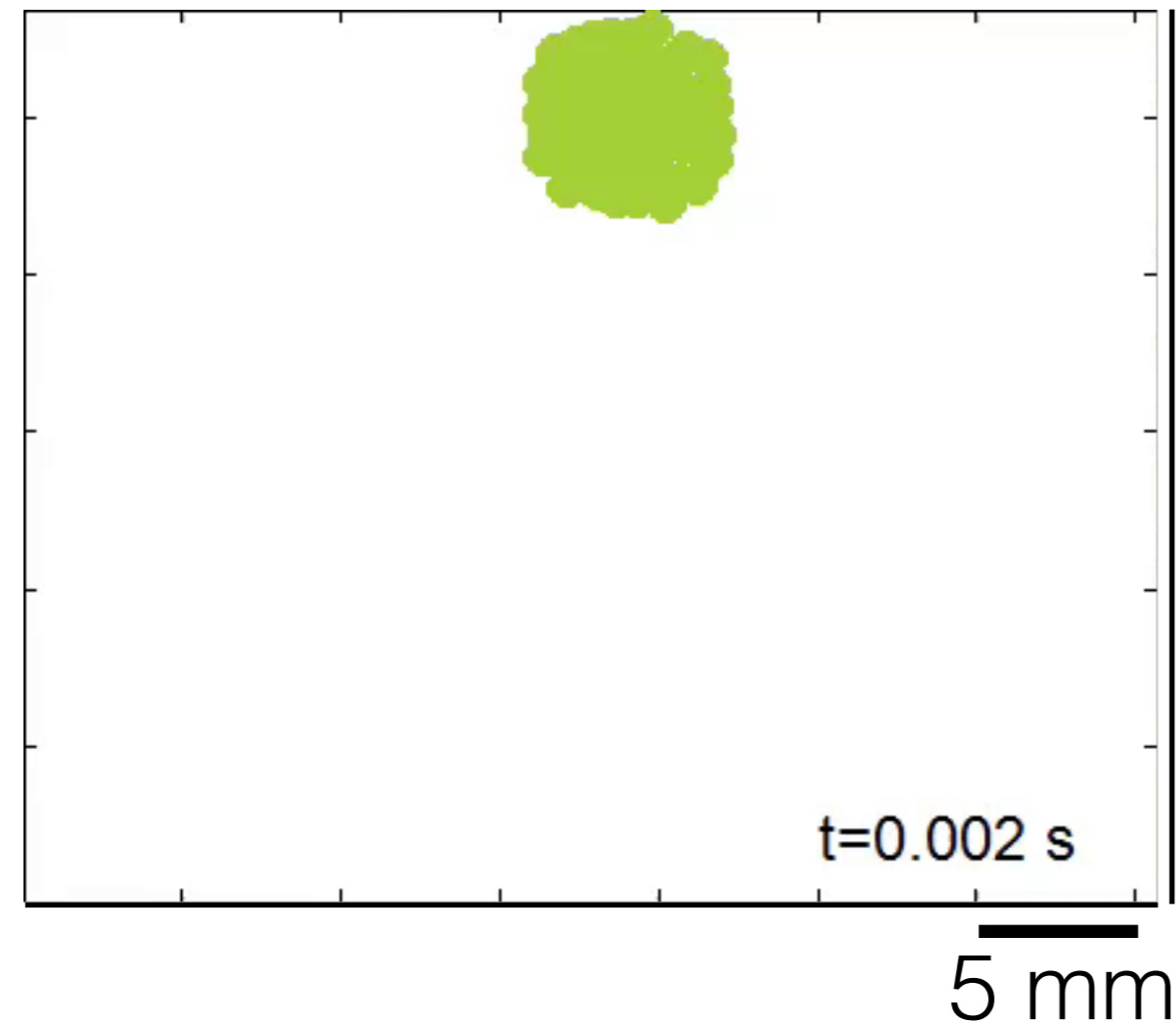
➤ a morphological parameter =  $\frac{\text{total area of the connected region}}{\# \text{ of connected regions}}$



# Morphological analysis of subgroups

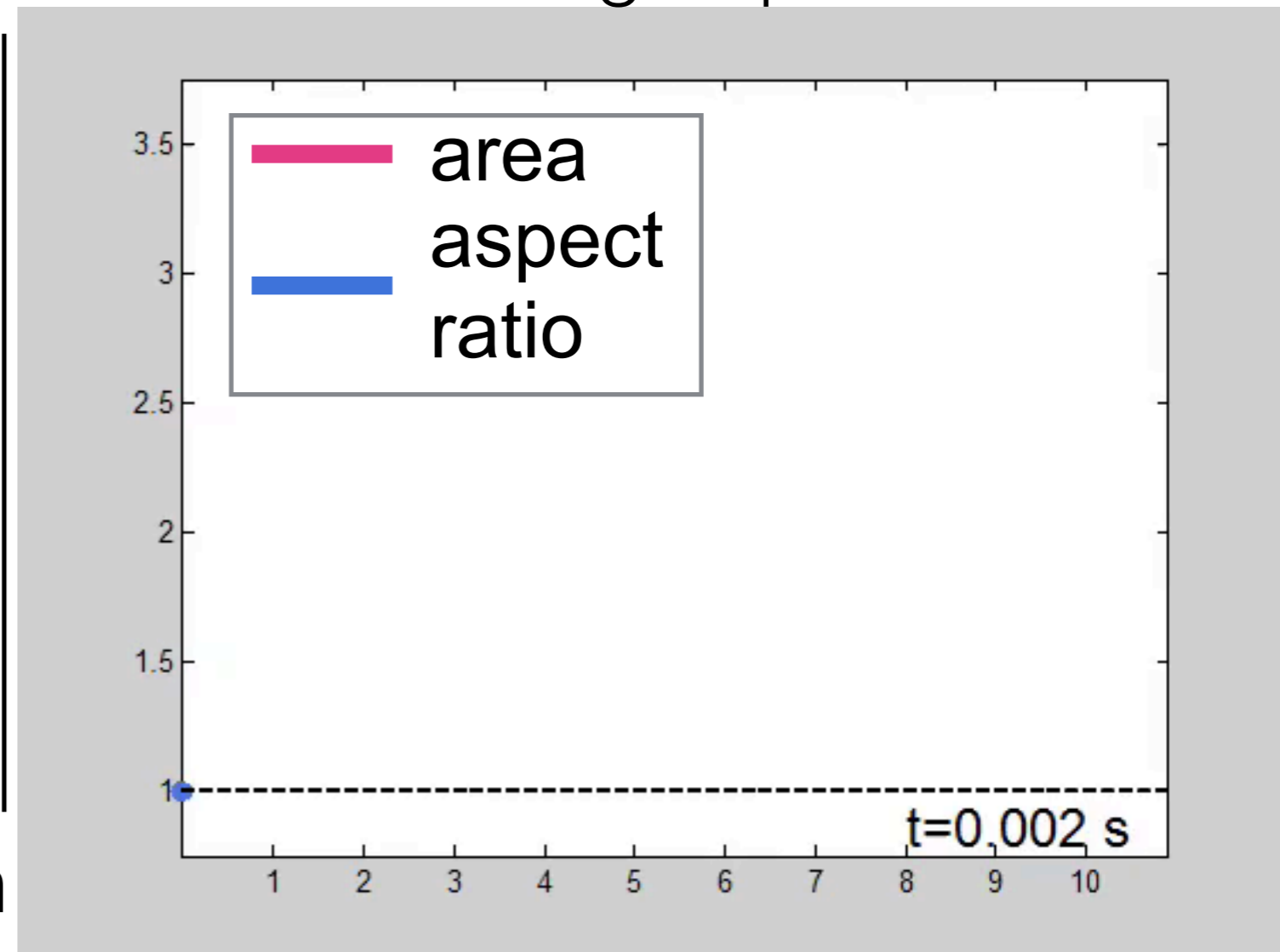
$$\phi = 0.76$$

- time evaluation of a subgroup:



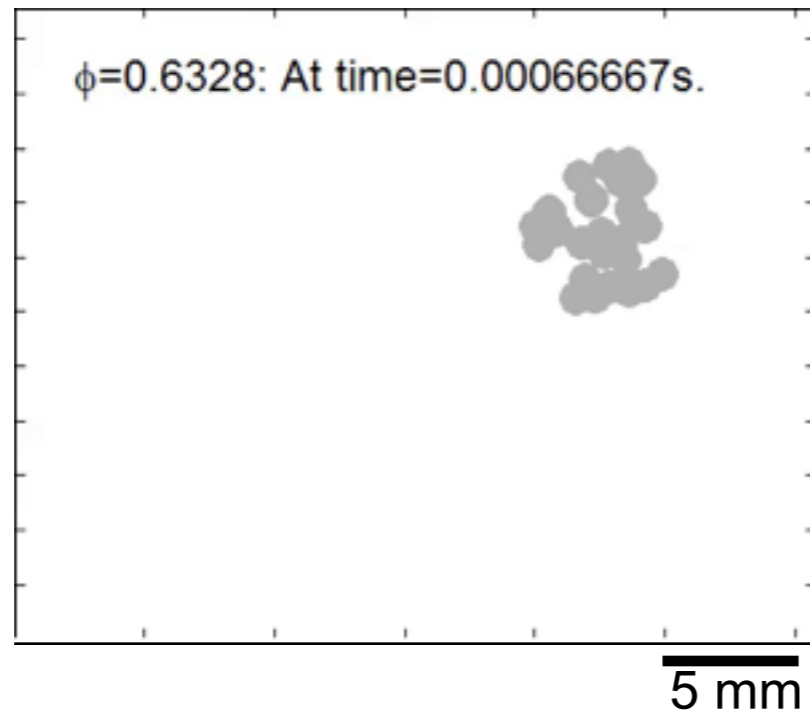
$$\phi = 0.76$$

- morphological deformation of the subgroup:

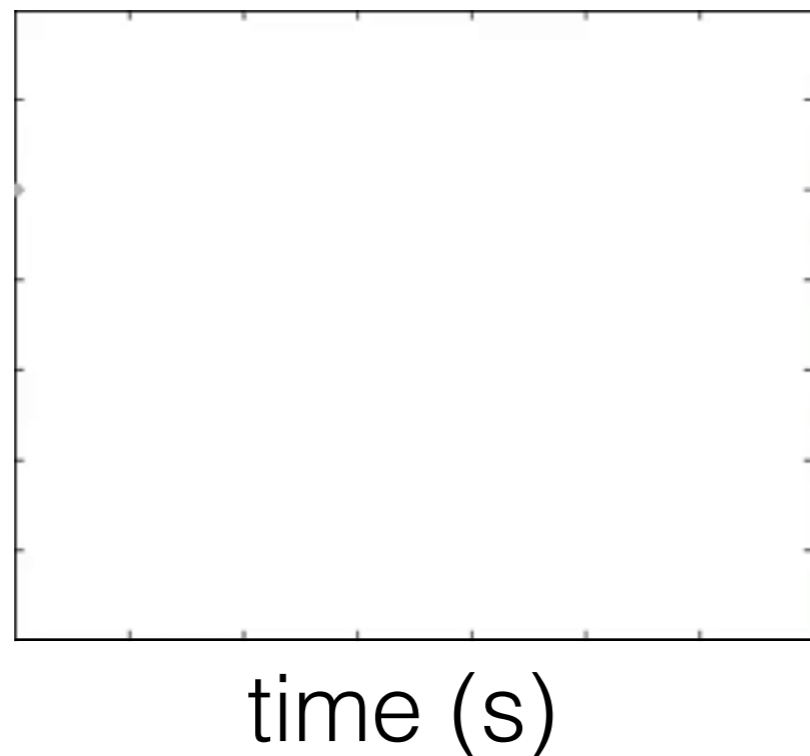


# Comparative analysis of subgroups

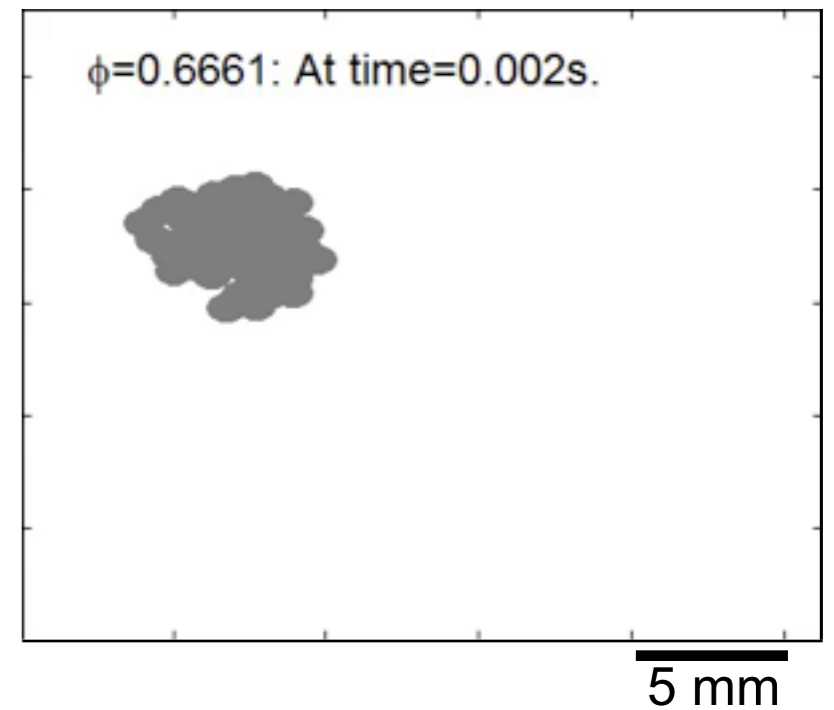
• low  $\phi$ :



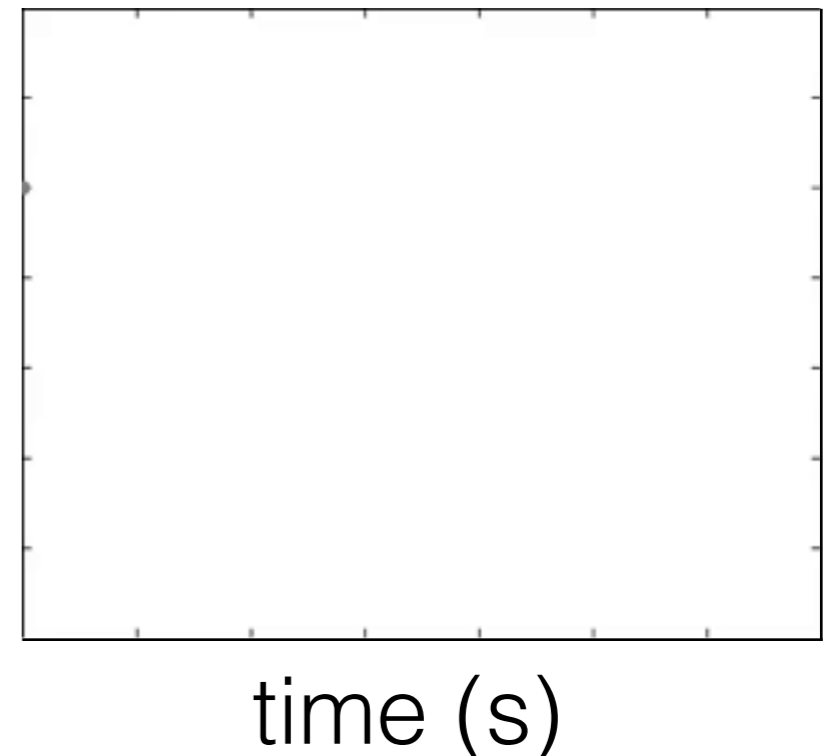
morphological  
parameter (area)



• high  $\phi$ :



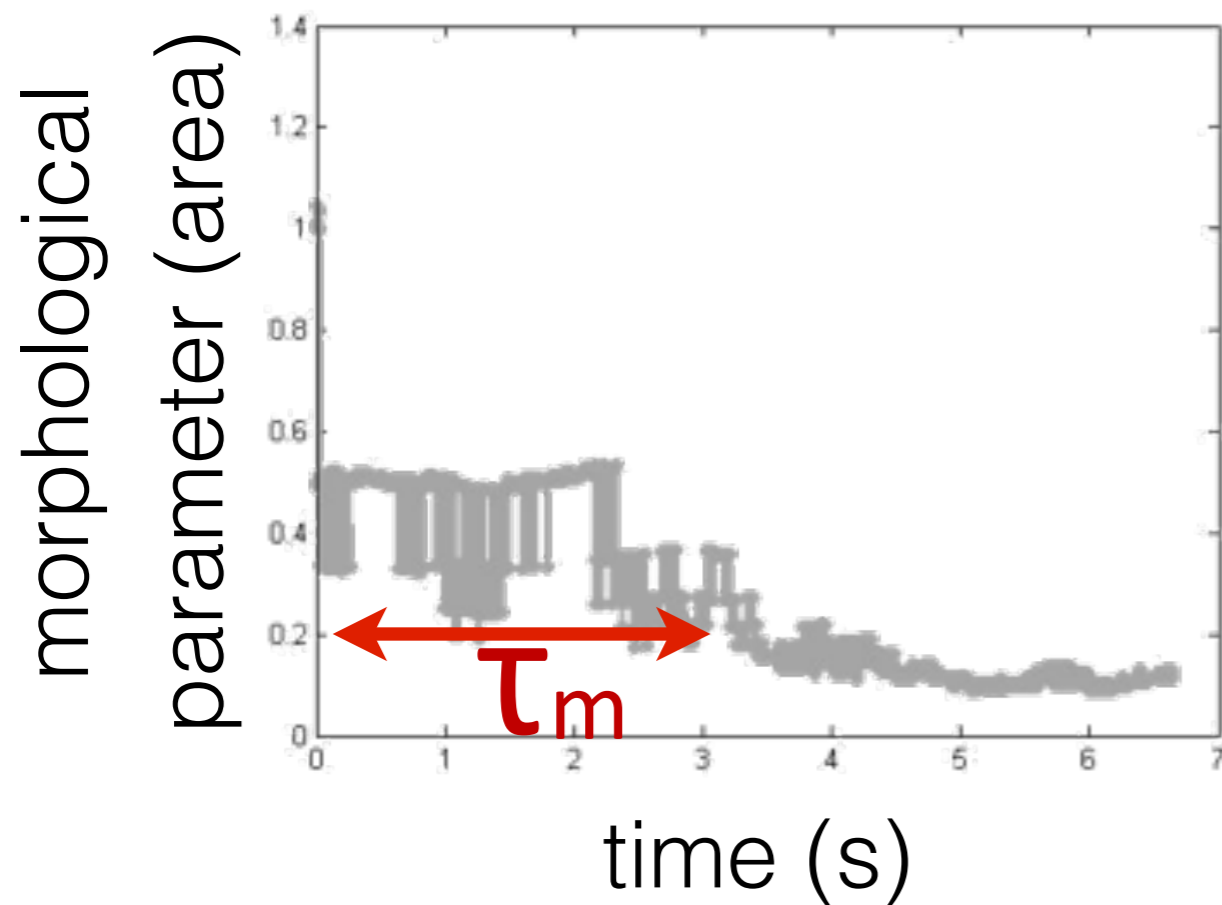
morphological  
parameter (area)



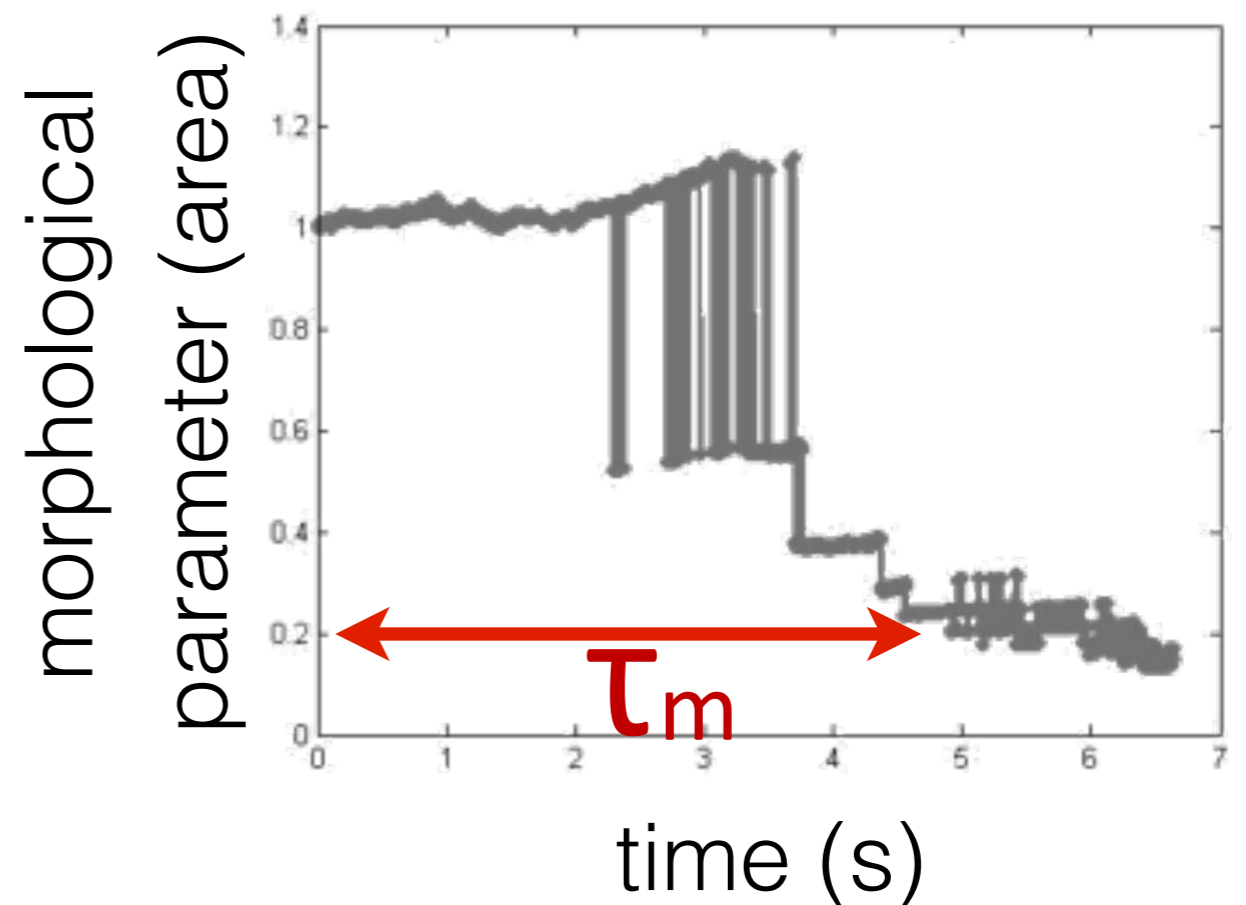


# Quantifying break-up time

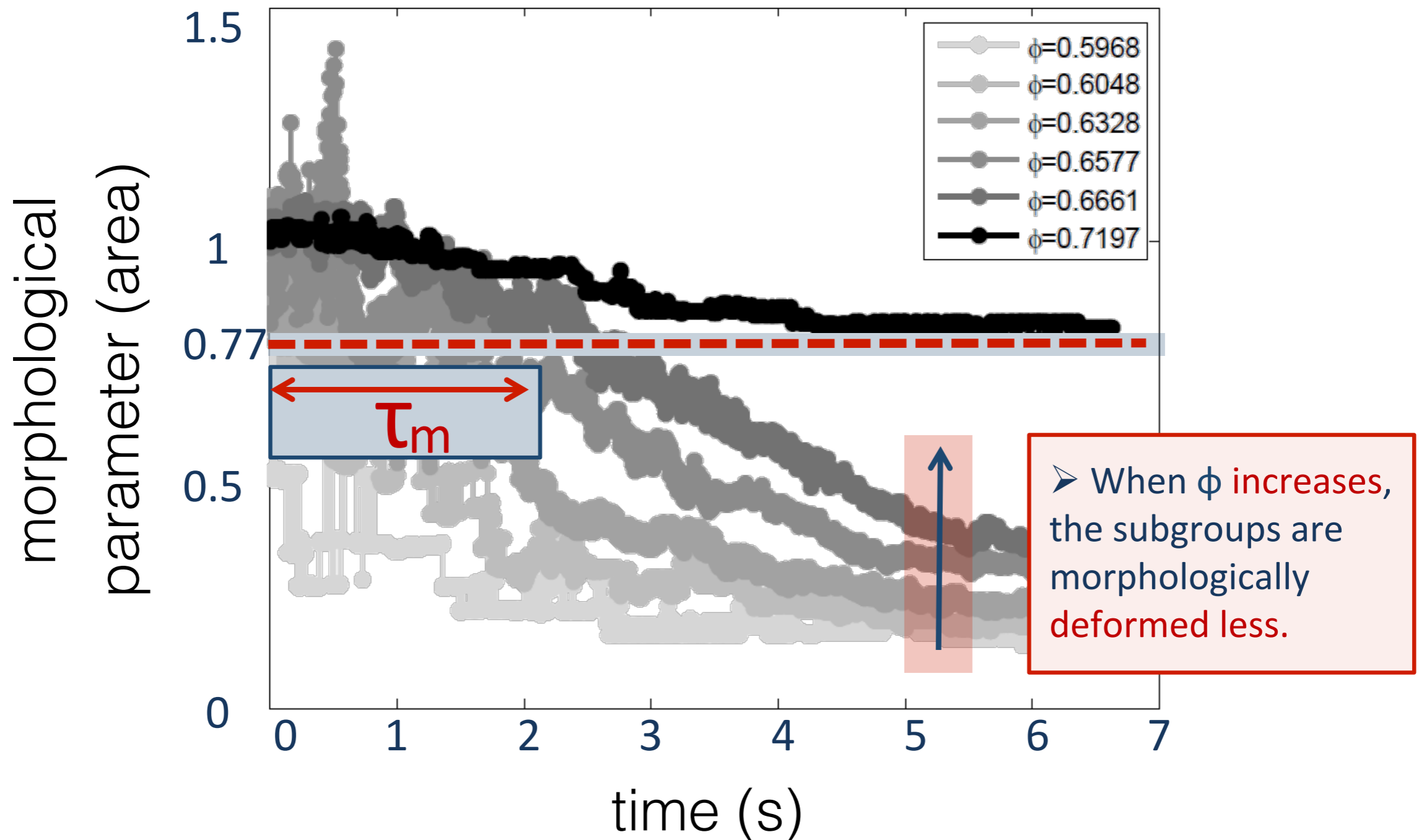
- low  $\phi$ :



- high  $\phi$ :

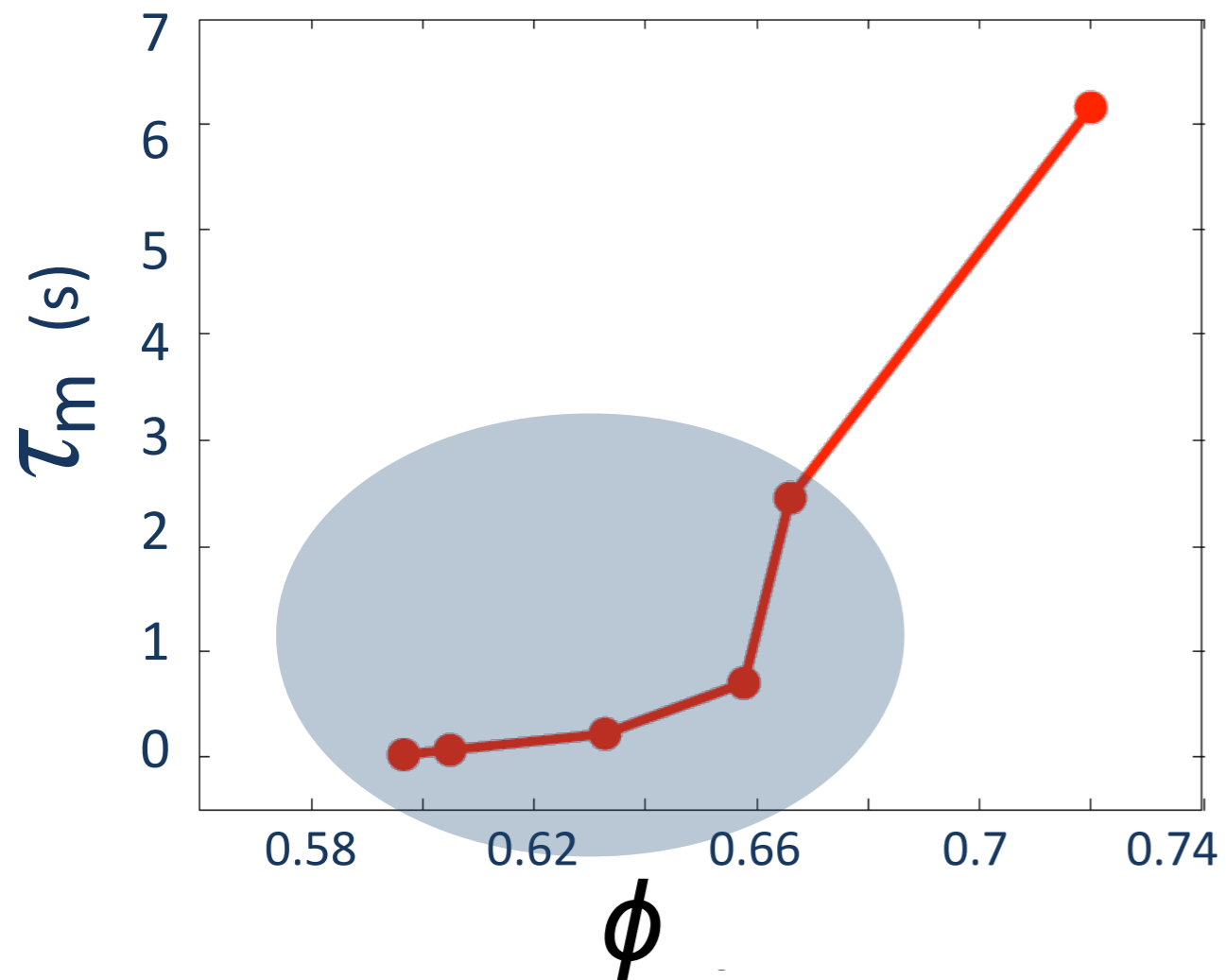


# Morphological parameter

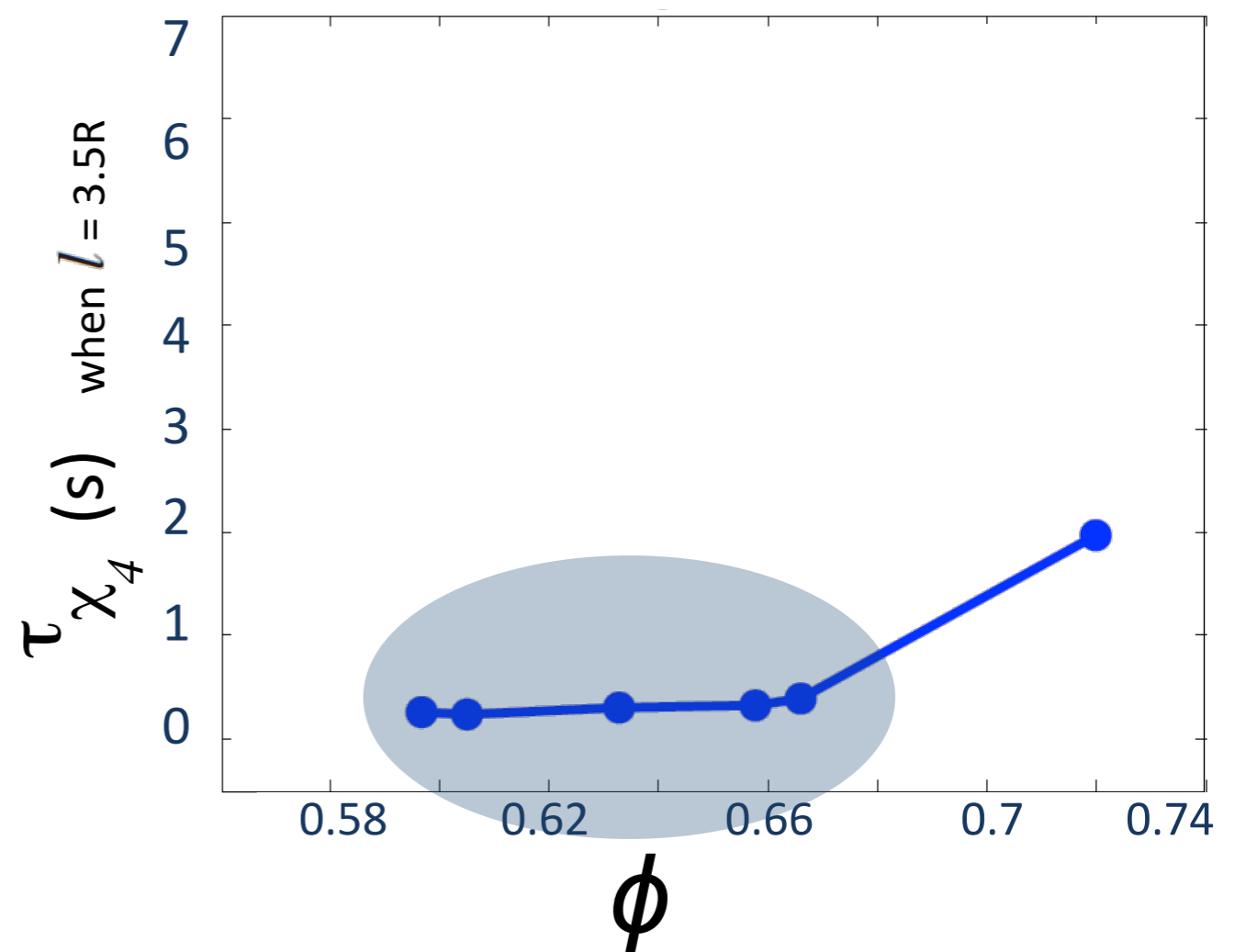


# Comparison: In progress

- Morphological analysis

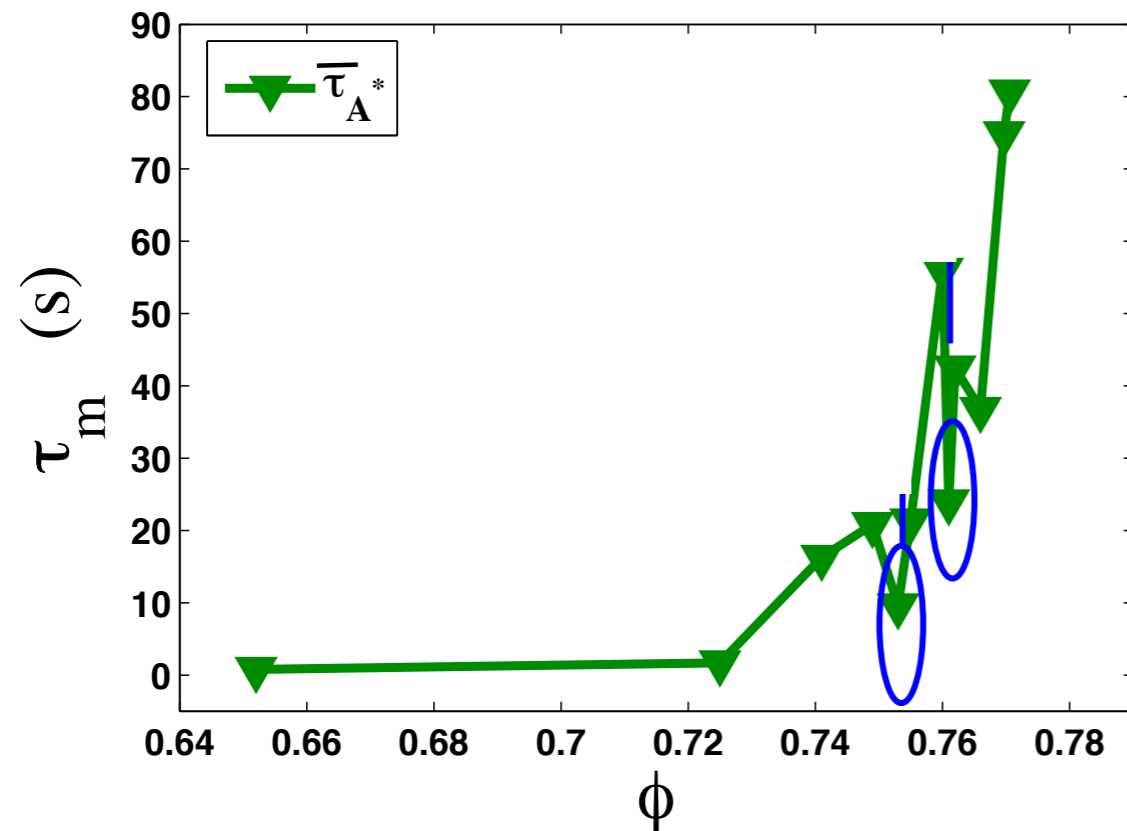


- Four-point dynamic susceptibility

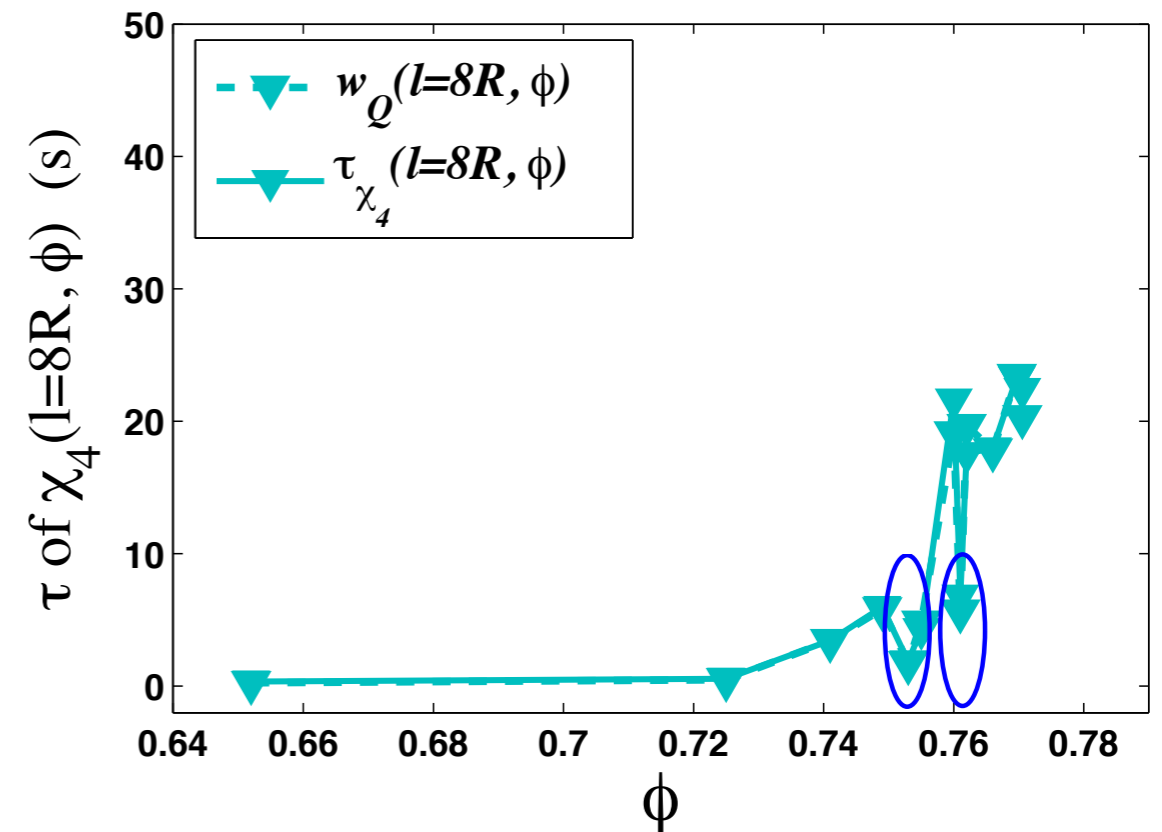


# Comparison: In progress

- Morphological analysis

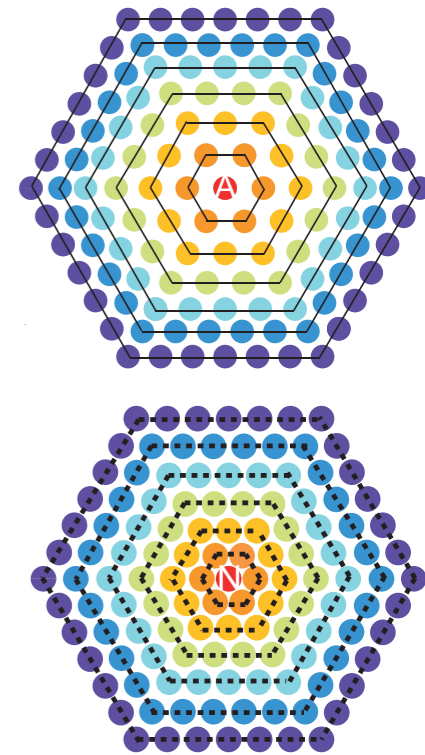
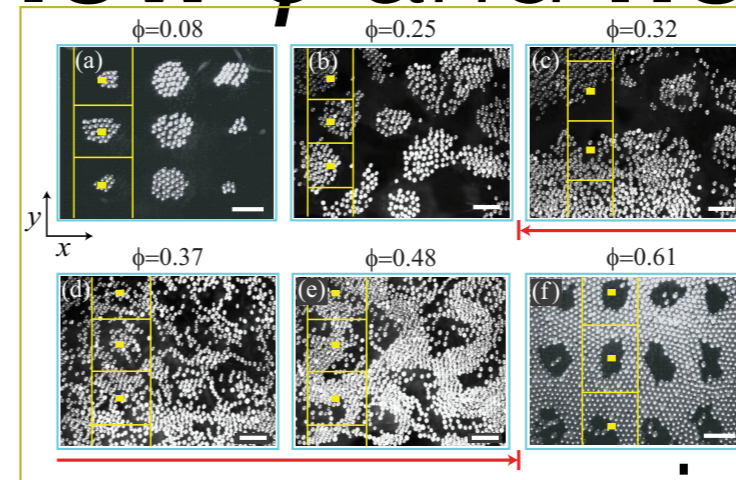


- Four-point dynamic susceptibility

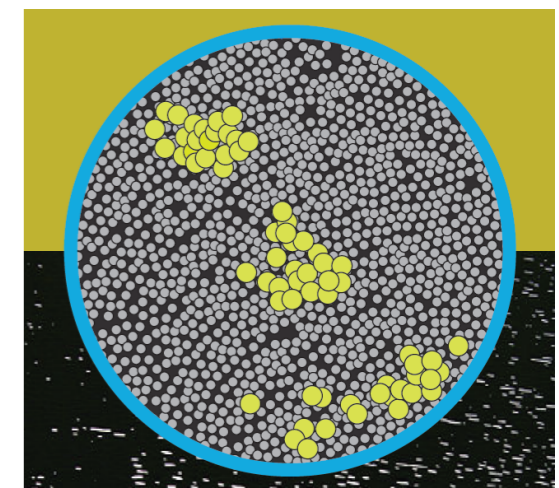


# Summary

- On a **standing** Faraday wave regime. **Antinode** clusters at **low  $\phi$**  and **node** clusters at **high  $\phi$** .

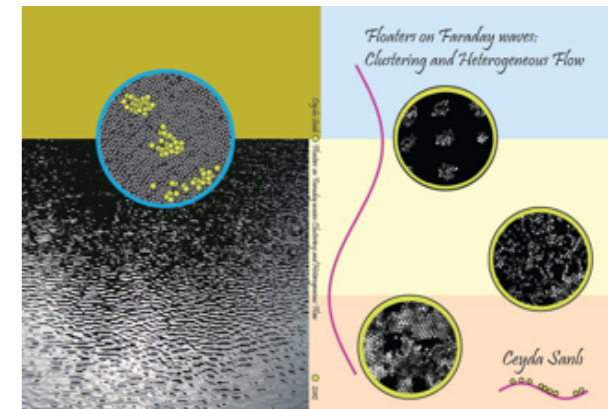


- On a **capillary** Faraday wave regime: We investigate an alternative approach to quantify **heterogeneous** flow by **group deformation**, usually done by the four-point dynamic susceptibility.

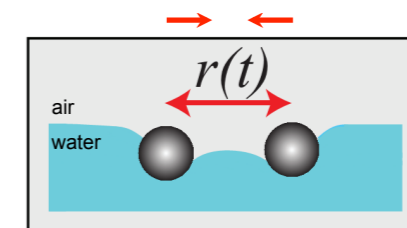


# Take home messages

- Dynamics of floaters is too complicated  
~ further **4** years project is indeed necessary.



- The deep **theory** behind and there is no exact solution, even for just **two floating spheres on a static surface**



- Currently, our **numerical simulation** cannot reach the experimental limits at **high  $\phi$** .